

# ON THE DIFFERENCE OF CONSECUTIVE EIGENVALUES OF UNIFORMLY ELLIPTIC OPERATORS OF HIGHER ORDERS\*\*

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## Abstract

This paper considers the upper bound for the difference of consecutive eigenvalues of a class of uniformly elliptic operators of higher orders. The upper bound of  $\lambda_{n+1} - \lambda_n$  is dependent on the first  $n - 1$  eigenvalues and the coefficients in equations.

**Keywords** Uniformly elliptic operator, Eigenvalue, Eigenfunction.

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## §1. Introduction

The estimates for the bound of  $(n+1)$ th eigenvalue as well as the difference of consecutive eigenvalues of the harmonic operator and its polynomials are well known (see [1-5]). These estimates are that the  $(n+1)$ th eigenvalue is bounded from above by an amount depending on the first  $n$  eigenvalues and being independent of the measure of the domain in which the problem is concerned. In this paper we generalize the same kind of problems to a certain uniformly elliptic operator of higher orders, and obtain the results similar to in form the ones for harmonic operator and its polynomials. The results in [1-5] are all corollaries of the theorems in this paper.

Suppose that  $\Omega$  is a bounded domain in  $R^m$  ( $m \geq 2$ ) with the piecewise smooth boundary  $\partial\Omega$  and  $u(x)$  is a solution of the problem

$$\begin{aligned} & (-1)^l \sum_{i_1, i_2, \dots, i_l=1}^m D_{i_1 \dots i_l} (a_{i_1 \dots i_l} D_{i_1 \dots i_l} u) = \lambda u, \quad x \in \Omega, \\ & u = \frac{\partial u}{\partial \nu} = \dots = \frac{\partial^{l-1} u}{\partial \nu^{l-1}} = 0, \quad l \geq 1, \quad x \in \partial\Omega, \end{aligned} \tag{1.1}$$

where  $\vec{\nu}$  is the unit outward normal to  $\partial\Omega$ .  $a_{i_1 i_2 \dots i_l}$  is a  $C^l(\bar{\Omega})$  function in  $x$  and satisfies the uniformly elliptic condition, i.e., there exists a constant  $\mu > 0$  such that

$$\min_{\bar{\Omega}} a_{i_1 i_2 \dots i_l} \geq \mu, \quad i_1, i_2, \dots, i_l = 1, 2, \dots, m. \tag{1.2}$$

Our main result is as follows.

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**Theorem 1.1.** Let  $\lambda_j$  be the eigenvalues of (1.1). Then we have

$$\begin{aligned} \lambda_{n+1} - \lambda_n &\leq \frac{4l}{\mu m^2 n^2} \left[ (2l+m-2) \max_{\bar{\Omega}} a_{i_1 i_2 \dots i_l} \left( \sum_{j=1}^n \lambda_j^{1-1/l} \right) + \right. \\ &\quad \left. + \mu^{1/2l} (l-1) \max_{\bar{\Omega}} |\nabla a_{i_1 \dots i_l}| \left( \sum_{j=1}^n \lambda_j^{1-3/2l} \right) \right] \left( \sum_{j=1}^n \lambda_j^{1/l} \right), \end{aligned} \quad (1.3)$$

$$\begin{aligned} \lambda_{n+1} &\leq \frac{1}{\mu m^2} \left[ \mu m^2 + 4l(2l+m-2) \max_{\bar{\Omega}} a_{i_1 \dots i_l} \right] \lambda_n + \\ &\quad + \frac{4}{\mu m^2} \mu^{1/2l} l(l-1) \max_{\bar{\Omega}} |\nabla a_{i_1 \dots i_l}| \lambda_n^{1-1/2l}. \end{aligned} \quad (1.4)$$

**Remark 1.1.** If we take  $a_{i_1 \dots i_l} = 1$  in (1.1) and  $\mu = 1$  in (1.2), then  $\max_{\bar{\Omega}} a_{i_1 \dots i_l} = 1$  and  $\max_{\bar{\Omega}} |\nabla a_{i_1 i_2 \dots i_l}| = 0$ . For this case, (1.3) yields

$$\lambda_{n+1} - \lambda_n \leq \frac{4l}{m^2 n^2} (2l+m-2) \left( \sum_{j=1}^n \lambda_j^{1-1/l} \right) \left( \sum_{j=1}^n \lambda_j^{1/l} \right), \quad (1.5)$$

which is just the result in [4] for polyharmonic operators. So the corresponding results in [1-5] are all included.

**Remark 1.2.** To see the sharpness of estimates (1.3) and (1.4), we consider, for simplicity, the following problem:

$$\begin{cases} -\Delta u = \lambda u, & x \in \Omega, \\ u = 0, & x \in \partial\Omega, \end{cases}$$

where  $\Omega \subset R^2$  is a unit square with  $0 < x_1 < 1$  and  $0 < x_2 < 1$ . A straightforward calculation yields  $\lambda_1 = 2\pi^2$ ,  $\lambda_2 = 5\pi^2$ ,  $\lambda_3 = 8\pi^2$  and so on. Therefore,

$$\lambda_2 - \lambda_1 = 3\pi^2, \text{ i.e., } \lambda_2 = 2.5\lambda_1; \quad \lambda_3 - \lambda_2 = 3\pi^2, \text{ i.e., } \lambda_3 = \lambda_1 + 1.2\lambda_2.$$

In this case estimate (1.5) with  $l = 1$ ,  $m = 2$  and  $n = 1, 2$  gives respectively

$$\lambda_2 - \lambda_1 \leq 2\lambda_1 (= 4\pi^2), \text{ i.e., } \lambda_2 \leq 3\lambda_1; \quad \lambda_3 - \lambda_2 \leq \lambda_1 + \lambda_2 (= 7\pi^2), \text{ i.e., } \lambda_3 \leq \lambda_1 + 2\lambda_2.$$

It is usually very difficult to get the exact values of eigenvalues. But it is enough to obtain the bounds of eigenvalues in many cases.

## §2. Proof of Theorem 1.1

In the sequel we will use the standard notations:

$$D_k = \frac{\partial}{\partial x_k}, \quad k = 1, 2, \dots, m, \quad \nabla = (D_1, D_2, \dots, D_m),$$

$$D_{i_1 i_2 \dots i_l} = D_{i_1} D_{i_2} \cdots D_{i_l}, \quad i_1, i_2, \dots, i_l = 1, 2, \dots, m, \text{ and } D_{i_0} = D_{i_{l+1}} = I.$$

From (1.1) we know

$$\lambda_j = \sum_{i_1, i_2, \dots, i_l=1}^m \int_{\Omega} a_{i_1 i_2 \dots i_l} |D_{i_1 i_2 \dots i_l} u_j|^2 dx, \quad j = 1, 2, \dots. \quad (2.1)$$

So we can order the eigenvalues of problem (1.1) so that  $0 < \lambda_1 < \lambda_2 \leq \dots \leq \lambda_n \leq \dots$  with the corresponding eigenfunctions  $u_1, u_2, \dots, u_n, \dots$  satisfying

$$\int_{\Omega} u_i u_j dx = \delta_{ij}, \quad i, j = 1, 2, \dots. \quad (2.2)$$

By virtue of (1.2),

$$\sum_{i_1, i_2, \dots, i_l=1}^m \int_{\Omega} |D_{i_1 \dots i_l} u_j|^2 dx \leq \lambda_j / \mu. \quad (2.3)$$

Let  $\phi_{jk}$  ( $j = 1, 2, \dots, n; k = 1, 2, \dots, m$ ) be the trial functions defined by

$$\phi_{jk} = \phi_{jk}(x) = x_k u_j - \sum_{p=1}^n b_{jp}^k u_p$$

where  $b_{jp}^k = \int_{\Omega} x_k u_j u_p dx$ ,  $j = 1, 2, \dots, n$ ,  $k = 1, 2, \dots, m$ . It is obvious that  $\phi_{jk}$  are orthogonal to  $u_1, u_2, \dots, u_n$  and

$$\phi_{jk} = \frac{\partial \phi_{jk}}{\partial \nu} = \dots = \frac{\partial^{l-1} \phi_{jk}}{\partial \nu^{l-1}} = 0, \quad x \in \partial \Omega. \quad (2.4)$$

Hence, we can use the well-known Rayleigh theorem to obtain

$$\lambda_{n+1} \leq \frac{(-1)^l \sum_{i_1, \dots, i_l=1}^m \int_{\Omega} \phi_{jk} D_{i_1 \dots i_l} (a_{i_1 \dots i_l} D_{i_1 \dots i_l} \phi_{jk}) dx}{\int_{\Omega} \phi_{jk}^2 dx}. \quad (2.5)$$

By a rather complicated calculus it yields

$$\begin{aligned} & (-1)^l \sum_{i_1, \dots, i_l=1}^m D_{i_1 \dots i_l} (a_{i_1 \dots i_l} D_{i_1 \dots i_l} \phi_{jk}) \\ &= \lambda_j x_k u_j - \sum_{p=1}^n b_{jp}^k \lambda_p u_p + \sum_{i_1, \dots, i_l=1}^m (-1)^l \sum_{t=1}^l \delta_{ki_t} D_{i_1 \dots i_{t-1} i_{t+1} \dots i_l} (a_{i_1 \dots i_l} D_{i_1 \dots i_l} u_j) \\ &+ \sum_{i_1, \dots, i_l=1}^m (-1)^l \sum_{t=1}^l \delta_{ki_t} D_{i_1 \dots i_l} (a_{i_1 \dots i_l} D_{i_1 \dots i_{t-1} i_{t+1} \dots i_l} u_j), \end{aligned}$$

where

$$\delta_{ki_t} = \begin{cases} 1, & \text{when } i_t = k, \\ 0, & \text{when } i_t \neq k, \end{cases} \quad i_t = 1, 2, \dots, m.$$

The orthogonality of  $\phi_{jk}$  with  $u_i$  ( $i = 1, 2, \dots, n$ ) gives the following result:

$$\begin{aligned} & \sum_{i_1, \dots, i_l=1}^m (-1)^l \int_{\Omega} \phi_{jk} D_{i_1 \dots i_l} (a_{i_1 \dots i_l} D_{i_1 \dots i_l} \phi_{jk}) dx \\ &= \lambda_j \int_{\Omega} \phi_{jk}^2 dx + \sum_{i_1, \dots, i_l=1}^m (-1)^l \sum_{t=1}^l \delta_{ki_t} \int_{\Omega} \phi_{jk} D_{i_1 \dots i_{t-1} i_{t+1} \dots i_l} (a_{i_1 \dots i_l} D_{i_1 \dots i_l} u_j) dx \\ &+ \sum_{i_1, \dots, i_l=1}^m (-1)^l \sum_{t=1}^l \delta_{ki_t} \int_{\Omega} \phi_{jk} D_{i_1 \dots i_l} (a_{i_1 \dots i_l} D_{i_1 \dots i_{t-1} i_{t+1} \dots i_l} u_j) dx. \end{aligned} \quad (2.6)$$

For simplifying writing we define  $I_{jki_t}, J_{jki_t}, I$  and  $J$  as follows:

$$I_{jki_t} = \sum_{i_1, \dots, i_l=1}^m (-1)^l \delta_{ki_t} \int_{\Omega} \phi_{jk} D_{i_1 \dots i_{t-1} i_{t+1} \dots i_l} (a_{i_1 \dots i_l} D_{i_1 \dots i_l} u_j) dx,$$

$$J_{jki_t} = \sum_{i_1, \dots, i_l=1}^m (-1)^l \delta_{ki_t} \int_{\Omega} \phi_{jk} D_{i_1 \dots i_l} (a_{i_1 \dots i_l} D_{i_1 \dots i_{t-1} i_{t+1} \dots i_l} u_j) dx,$$

$$I = \sum_{j=1}^n \sum_{k=1}^m \sum_{t=1}^l I_{jki_t}, \quad J = \sum_{j=1}^n \sum_{k=1}^m \sum_{t=1}^l J_{jki_t}.$$

Then, it follows from (2.6) that

$$\begin{aligned} & \sum_{j=1}^n \sum_{k=1}^m \sum_{i_1, \dots, i_l=1}^m (-1)^l \int_{\Omega} \phi_{jk} D_{i_1 \dots i_l} (a_{i_1 \dots i_l} D_{i_1 \dots i_l} \phi_{jk}) dx \\ &= \sum_{j=1}^n \sum_{k=1}^m \lambda_j \int_{\Omega} \phi_{jk}^2 dx + I + J. \end{aligned} \quad (2.7)$$

By (2.5),

$$\lambda_{n+1} \left( \sum_{j=1}^n \sum_{k=1}^m \int_{\Omega} \phi_{jk}^2 dx \right) \leq \sum_{j=1}^n \sum_{k=1}^m \lambda_j \int_{\Omega} \phi_{jk}^2 dx + I + J. \quad (2.8)$$

Replacing  $\lambda_j$  in (2.8) by  $\lambda_n$  yields

$$(\lambda_{n+1} - \lambda_n) \left( \sum_{j=1}^n \sum_{k=1}^m \int_{\Omega} \phi_{jk}^2 dx \right) \leq I + J. \quad (2.9)$$

Introducing an operator  $\nabla^p$  such that

$$\nabla^p = \begin{cases} \Delta^{p/2}, & p = 2k, k = 1, 2, \dots, \\ \nabla(\Delta^{(p-1)/2}), & p = 2k-1, k = 1, 2, \dots, \end{cases}$$

we can easily prove, by the inductive method, that for an eigenfunction  $u_j$  of (1.1) it holds that

$$\sum_{i_1, \dots, i_l=1}^m \int_{\Omega} |D_{i_1 \dots i_l} u_j|^2 dx = \int_{\Omega} |\nabla^r u_j|^2 dx, \quad 1 \leq r \leq l \quad (2.10)$$

and

$$\left( \int_{\Omega} |\nabla^s u_j|^2 dx \right)^{1/s} \leq \left( \int_{\Omega} |\nabla^{s+1} u_j|^2 dx \right)^{1/(s+1)}, \quad 1 \leq s \leq l. \quad (2.11)$$

(2.10) and (2.3) yield

$$\int_{\Omega} |\nabla^l u_j|^2 dx \leq \lambda_j / \mu. \quad (2.12)$$

Using (2.11) inductively and (2.12), we find

$$\int_{\Omega} |\nabla^s u_j|^2 dx \leq (\lambda_j / \mu)^{s/l}, \quad 1 \leq s \leq l. \quad (2.13)$$

Now, we have the following essential lemma which plays a key role in the proof of Theorem 1.1.

**Lemma.** For  $I + J$  the following estimate holds

$$\begin{aligned} I + J \leq & l(2l + m - 2) \max_{\bar{\Omega}} a_{i_1 \dots i_l} \sum_{j=1}^n (\lambda_j / \mu)^{1-1/l} + \\ & + l(l-1) \max_{\bar{\Omega}} |\nabla a_{i_1 \dots i_l}| \sum_{j=1}^n (\lambda_j / \mu)^{1-3/2l}. \end{aligned}$$

**Proof.** In the sequel, we denote  $i_1 i_2 \dots i_{t-1} i_{t+1} \dots i_l$  and  $i_1 \dots i_{t-1} i_{t+1} \dots i_{s-1} i_{s+1} \dots i_l$  by  $(i_t)$  and  $(i_t i_s)$  respectively. Integrating by parts for  $I_{jki_t}$ , it yields

$$\begin{aligned} I_{jki_t} = & - \sum_{i(l)=1}^m \delta_{k i_t} \int_{\Omega} x_k a_{i(l)} D_{(i_t)} u_j D_{i(l)} u_j dx - \sum_{i(l)=1}^m \sum_{\substack{s=1 \\ s \neq t}}^l \delta_{k i_t} \delta_{k i_s} \int_{\Omega} a_{i(l)} D_{(i_t i_s)} u_j D_{i(l)} u_j dx \\ & + \sum_{p=1}^n b_{jp}^k \sum_{i(l)=1}^m \delta_{k i_t} \int_{\Omega} a_{i(l)} D_{(i_t)} u_p D_{i(l)} u_j dx, \end{aligned} \quad (2.14)$$

where  $D_{i(l)} = D_{i_1 \dots i_l}$ ,  $a_{i(l)} = a_{i_1 \dots i_l}$ , and  $\sum_{i(l)=1}^m = \sum_{i_1, \dots, i_l=1}^m$ .

Similarly,

$$\begin{aligned} J_{jki_t} = & \sum_{i(l)=1}^m \delta_{k i_t} \int_{\Omega} x_k a_{i(l)} D_{i(l)} u_j D_{(i_t)} u_j dx + \sum_{i(l)=1}^m \sum_{s=1}^l \delta_{k i_t} \delta_{k i_s} \int_{\Omega} a_{i(l)} D_{(i_t)} u_j D_{(i_s)} u_j dx \\ & - \sum_{i(l)=1}^m \sum_{p=1}^n b_{jp}^k \delta_{k i_t} \int_{\Omega} a_{i(l)} D_{i(l)} u_p D_{(i_t)} u_j dx. \end{aligned} \quad (2.15)$$

Combining (2.14) with (2.15) yields

$$\begin{aligned} \sum_{j=1}^n (I_{jki_t} + J_{jki_t}) = & - \sum_{j=1}^n \sum_{i(l)=1}^m \sum_{\substack{s=1 \\ s \neq t}}^l \delta_{k i_t} \delta_{k i_s} \int_{\Omega} a_{i(l)} D_{i(l)} u_j D_{(i_t i_s)} u_j dx \\ & + \sum_{j=1}^n \sum_{i(l)=1}^m \sum_{s=1}^l \delta_{k i_t} \delta_{k i_s} \int_{\Omega} a_{i(l)} D_{(i_t)} u_j D_{(i_s)} u_j dx. \end{aligned}$$

Define

$$\begin{aligned} I_{jki_t i_s}^* = & - \sum_{i(l)=1}^m \delta_{k i_t} \delta_{k i_s} \int_{\Omega} a_{i(l)} D_{i(l)} u_j D_{(i_t i_s)} u_j dx, \\ J_{jki_t i_s}^* = & \sum_{i(l)=1}^m \delta_{k i_t} \delta_{k i_s} \int_{\Omega} a_{i(l)} D_{(i_s)} u_j D_{(i_t)} u_j dx. \end{aligned} \quad (2.16)$$

We then have

$$\sum_{j=1}^n (I_{jki_t} + J_{jki_t}) = \sum_{j=1}^n \left( \sum_{\substack{s=1 \\ s \neq t}}^l I_{jki_t i_s}^* + \sum_{s=1}^l J_{jki_t i_s}^* \right).$$

By (2.16),

$$\begin{aligned} I_{jki_t i_s}^* &= \sum_{i(l)=1}^m \delta_{ki_t} \delta_{ki_s} \int_{\Omega} a_{i(l)} D_{(i_t)} u_j D_{(i_s)} u_j dx + \\ &\quad + \sum_{i(l)=1}^m \delta_{ki_t} \delta_{ki_s} \int_{\Omega} (D_{i_t} a_{i(l)}) D_{(i_t)} u_j D_{(i_t i_s)} u_j dx. \end{aligned}$$

Using Schwartz inequality yields

$$\begin{aligned} \sum_{k=1}^m I_{jki_t i_s}^* &\leq \max_{\bar{\Omega}} a_{i(l)} \sum_{k=1}^m \sum_{i(l)=1}^m \delta_{ki_t} \delta_{ki_s} \int_{\Omega} |D_{(i_t)} u_j| |D_{(i_s)} u_j| dx \\ &\quad + \left( \sum_{k=1}^m \sum_{i(l)=1}^m \delta_{ki_t} \delta_{ki_s} \int_{\Omega} |D_{(i_t)} u_j|^2 dx \right)^{1/2} \left( \sum_{k=1}^m \sum_{i(l)=1}^m \delta_{ki_t} \delta_{ki_s} \int_{\Omega} |D_{i_t} a_{i(l)}|^2 |D_{(i_t i_s)} u_j|^2 dx \right)^{1/2}. \end{aligned} \quad (2.17)$$

For  $t \neq s$ , using (2.10) and (2.13) we find

$$\begin{aligned} \sum_{k=1}^m \sum_{i(l)=1}^m \delta_{ki_t} \delta_{ki_s} \int_{\Omega} |D_{(i_t)} u_j| |D_{(i_s)} u_j| dx &= \sum_{k=1}^m \sum_{(i_t i_s)=1}^m \int_{\Omega} |D_k D_{(i_t i_s)} u_j|^2 dx \\ &= \int_{\Omega} |\nabla^{l-1} u_j|^2 dx \leq (\lambda_j/\mu)^{1-1/l}. \end{aligned} \quad (2.18)$$

Similarly,

$$\sum_{k=1}^m \sum_{i(l)=1}^m \delta_{ki_t} \delta_{ki_s} \int_{\Omega} |D_{(i_t)} u_j|^2 dx = \int_{\Omega} |\nabla^{l-1} u_j|^2 dx \leq (\lambda_j/\mu)^{1-1/l}, \quad (2.19)$$

$$\sum_{k=1}^m \sum_{i(l)=1}^m \delta_{ki_t} \delta_{ki_s} \int_{\Omega} |D_{i_t} a_{i(l)}|^2 |D_{(i_t i_s)} u_j|^2 dx \leq \max_{\bar{\Omega}} |\nabla a_{i(l)}|^2 (\lambda_j/\mu)^{1-2/l}. \quad (2.20)$$

Substituting (2.18), (2.19) and (2.20) into (2.17), we have

$$\sum_{k=1}^m I_{jki_t i_s}^* \leq \max_{\bar{\Omega}} a_{i(l)} (\lambda_j/\mu)^{1-1/l} + \max_{\bar{\Omega}} |\nabla a_{i(l)}| (\lambda_j/\mu)^{1-3/2l}.$$

Hence,

$$\sum_{\substack{s,t=1 \\ s \neq t}}^l \sum_{k=1}^m I_{jki_t i_s}^* \leq l(l-1) \left[ \max_{\bar{\Omega}} a_{i(l)} (\lambda_j/\mu)^{1-1/l} + \max_{\bar{\Omega}} |\nabla a_{i(l)}| (\lambda_j/\mu)^{1-3/2l} \right]. \quad (2.21)$$

Similarly, for  $J_{jki_t i_s}$  we have

$$\sum_{\substack{s,t=1 \\ s \neq t}}^l \sum_{k=1}^m J_{jki_t i_s}^* \leq l(l-1) \max_{\bar{\Omega}} a_{i(l)} (\lambda_j/\mu)^{1-1/l}, \quad (2.22)$$

$$\sum_{s=t=1}^l \sum_{k=1}^m J_{jki_t i_s}^* \leq ml \max_{\bar{\Omega}} a_{i(l)} (\lambda_j/\mu)^{1-1/l}. \quad (2.23)$$

Combining (2.21), (2.22) and (2.23) the lemma then follows.

To finish the proof of Theorem 1.1, we need the estimate that

$$\sum_{j=1}^n \sum_{k=1}^m \int_{\Omega} \phi_{jk}^2 dx \geq \mu^{1/l} m^2 n^2 \left( \sum_{j=1}^n \lambda_j^{1/l} \right)^{-1} / 4. \quad (2.24)$$

This is true. In fact, from the definition of  $\phi_{jk}$  we have

$$\sum_{j=1}^n \int_{\Omega} \phi_{jk} D_k u_j dx = \sum_{j=1}^n \int_{\Omega} x_k u_j D_k u_j dx - \sum_{p,j=1}^n b_{jp}^k \int_{\Omega} u_p D_k u_j dx. \quad (2.25)$$

Because of  $b_{jp}^k = b_{pj}^k$  the second term on the right hand side of (2.25) is zero. From the identity of  $\int_{\Omega} x_k u_j D_k u_j dx = -\frac{1}{2}$ , we get

$$\sum_{k=1}^m \sum_{j=1}^n \int_{\Omega} \phi_{jk} D_k u_j dx = -mn/2.$$

By Schwartz inequality,

$$m^2 n^2 / 4 \leq \left( \sum_{k=1}^m \sum_{j=1}^n \int_{\Omega} \phi_{jk}^2 dx \right) \sum_{j=1}^n \int_{\Omega} |\nabla u_j|^2 dx. \quad (2.26)$$

From (2.13) we know that  $\int_{\Omega} |\nabla u_j|^2 dx \leq (\lambda_j / \mu)^{1/l}$ , which with (2.26) gives (2.24). Substituting (2.24) into (2.9), and using our lemma, assertion (1.3) then follows. Replacing  $\lambda_j$  in (1.3) by  $\lambda_n$  yields (1.4). This completes the proof.

**Remark 2.1.** Similar to the method in [3], we can get for problem (1.1) the following estimate

$$\begin{aligned} \sum_{j=1}^n \frac{\lambda_j^{1/l}}{\lambda_{n+1} - \lambda_j} &\geq \frac{\mu m^2 n^2}{4} \left[ l(2l+m-2) \max_{\bar{\Omega}} a_{i_1 \dots i_l} \left( \sum_{j=1}^n \lambda_j^{1-1/l} \right) \right. \\ &\quad \left. + l(l-1) \mu^{1/2l} \max_{\bar{\Omega}} |\nabla a_{i_1 \dots i_l}| \left( \sum_{j=1}^n \lambda_j^{1-3/2l} \right) \right]^{-1}, \end{aligned} \quad (2.27)$$

where  $m \geq 2$ ,  $n \geq 1$  and  $l \geq 1$ . If  $a_{i_1 \dots i_l}$  is as in Remark 1.1 and  $\mu = 1$ , then (2.27) takes the form of

$$\sum_{j=1}^n \frac{\lambda_j^{1/l}}{\lambda_{n+1} - \lambda_j} \geq \frac{m^2 n^2}{4l(2l+m-2)} \left( \sum_{j=1}^n \lambda_j^{1-1/l} \right)^{-1},$$

which is the result obtained in [4] for polyharmonic operator. Moreover, let  $l = 1, 2$  respectively. Then the results in [1] and [3] follow.

**Remark 2.2.** By the same method as above we can deduce the upper bounds of eigenvalues for the more general uniformly elliptic operators as follows

$$\begin{cases} \sum_{r=1}^l \sum_{i_1, \dots, i_r=1}^m (-1)^r D_{i_1 \dots i_r} (a_{i_1 \dots i_r} D_{i_1 \dots i_r} u) = \lambda u, & x \in \Omega, \\ u = \frac{\partial u}{\partial \nu} = \dots = \frac{\partial^{l-1} u}{\partial \nu^{l-1}} = 0, & x \in \partial \Omega, \end{cases} \quad (2.28)$$

where  $a_{i_1 \dots i_l} \in C^l(\bar{\Omega})$  and  $\min_{\bar{\Omega}} a_{i_1 \dots i_l} \geq \mu > 0$ ,  $a_{i_1 \dots i_r} \geq 0$ ,  $a_{i_1 \dots i_r} \in C^r(\bar{\Omega})$ ,  $r = 1, 2, \dots, l-1$ ,

$i_1, i_2, \dots, i_l = 1, 2, \dots, m$ . The results are

$$\sum_{j=1}^n \frac{\lambda_j^{1/l}}{\lambda_{n+1} - \lambda_j} \geq \frac{\mu^{1/l} m^2 n^2}{4} \left[ \sum_{r=1}^l \sum_{j=1}^n r(2r+m-2) \max_{\bar{\Omega}} a_{i_1 \dots i_r} (\lambda_j/\mu)^{(r-1)/l} \right. \\ \left. + \sum_{r=1}^l \sum_{j=1}^n r(r-1) \max_{\bar{\Omega}} |\nabla a_{i_1 \dots i_r}| (\lambda_j/\mu)^{(2r-3)/2l} \right]^{-1}, \quad (2.29)$$

$$\lambda_{n+1} - \lambda_n \leq \frac{4}{m^2 n^2} \left[ \sum_{r=1}^l \sum_{j=1}^n r(2r+m-2) \max_{\bar{\Omega}} a_{i_1 \dots i_r} (\lambda_j/\mu)^{(r-1)/l} \right. \\ \left. + \sum_{r=1}^l \sum_{j=1}^n r(r-1) \max_{\bar{\Omega}} |\nabla a_{i_1 \dots i_r}| (\lambda_j/\mu)^{(2r-3)/2l} \right] \sum_{j=1}^n (\lambda_j/\mu)^{1/l}. \quad (2.30)$$

These are the generalization of the corresponding results in [4]. Therefore, [1], [2], [3] and [5] are special cases of this paper since [4] is a generalization of them.

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