

“MEANED” EXACT CONTROLLABILITY FOR WAVE EQUATIONS

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Abstract

With a control of Dirichlet type the author wants to bring the wave system to a “meaned” state in the time interval $[T-h, T]$. With the HUM Method the optimal control is got and its limit behavior when h goes to zero is studied.

Keywords Wave equation, Control of Dirichlet type, Optimal control.

1991 MR Subject Classification 35L05, 35B37.

§1. Introduction and the Principal Results

Let $\Omega \subset \mathbb{R}^n$ be a bounded domain with smooth boundary $\partial\Omega = \Gamma$. Following the notations of J. L. Lions^[1,2], we set

$$\Gamma(x^0) = \{x \in \Gamma; \quad \nu(x) \cdot m(x) > 0\}$$

for some point $x^0 \in \mathbb{R}^n$, $\nu(x)$ is the outunit normal to the boundary at the point x and $m(x) = x - x^0$. Then we denote:

$$\Sigma = \Gamma \times (0, T), \quad \Sigma(x^0) = \Gamma(x^0) \times (0, T), \quad R(x^0) = \max\{|x - x^0|; x \in \bar{\Omega}\}, \quad T(x^0) = 2R(x^0).$$

We consider the following problem:

$$\begin{cases} y'' - \Delta y = 0 & \text{in } Q = \Omega \times (0, T), \\ y = \begin{cases} v \text{ (control)} & \text{on } \Sigma(x^0), \\ 0 & \text{on } \Sigma \setminus \Sigma(x^0), \end{cases} \\ y(0) = y^0, y'(0) = y^1 & \text{on } \Omega \times \{0\}. \end{cases} \quad (1.1)$$

For any given initial data $\{y^0, y^1\} \in L^2(\Omega) \times H^{-1}(\Omega)$, we will denote for simplicity:

$$y := y(v) := y(t; v) := y(x, t; v),$$

the weak solution of (1.1) associated with the control v .

By the results of J. L. Lions^[1,2], we know that if $T > T(x^0)$, then for any $\{z^0, z^1\} \in L^2(\Omega) \times H^{-1}(\Omega)$, there exists a v_0 such that

$$\begin{aligned} v_0 \in \mathcal{U}_{ad} &= \{v \in L^2(\Sigma(x^0)); \quad \text{s.t. } y(T; v) = z^0, y'(T; v) = z^1\}, \\ J(v_0) &= \min_{v \in \mathcal{U}_{ad}} J(v), \quad \text{where } J(v) = \frac{1}{2} \int_{\Sigma(x^0)} v^2 d\Sigma, \\ v_0 &= \frac{\partial \Phi}{\partial \nu} \quad \text{on } \Sigma(x^0), \end{aligned} \quad (1.2)$$

Manuscript received March 5, 1991. Revised May 6, 1992.

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where Φ is the solution of the following problem for a certain $\{\Phi^0, \Phi^1\} \in H^1(\Omega) \times L^2(\Omega)$:

$$\begin{cases} \Phi'' - \Delta \Phi = 0 & \text{in } Q, \\ \Phi = 0 & \text{on } \Sigma, \\ \Phi(T) = \Phi^0, \Phi'(T) = \Phi^1 & \text{on } \Omega \times \{T\}. \end{cases} \quad (1.3)$$

Now we are going to search $\{z_h^0, z_h^1\}$ and v_h for any $h > 0$, such that

$$v_h \in \mathcal{U}_{ad,h} = \left\{ v \in L^2(\Sigma(x^0)); \text{ s.t. } \begin{cases} \frac{1}{h} \int_{T-h}^T y(t;v) dt = z_h^0, & \frac{1}{h} \int_{T-h}^T y'(t;v) dt = z_h^1 \end{cases} \right\}, \quad (1.4)$$

$$J(v_h) = \min_{v \in \mathcal{U}_{ad,h}} J(v).$$

And when $h \rightarrow 0$, if we have

$$\{z_h^0, z_h^1\} \rightarrow \{z^0, z^1\} \quad (1.5)$$

by a convenient topology, we would like to have

$$v_h \rightarrow v_0 \quad (1.6)$$

by a certain topology.

Remark 1.1. We can see that for any $\{y^0, y^1\} \in L^2(\Omega) \times H^{-1}(\Omega)$ and any $v \in L^2(\Sigma(x^0))$, we have

$$\Delta\left(\frac{1}{h} \int_{T-h}^T y(t;v) dt\right) \in H^{-1}(\Omega), \quad \frac{1}{h} \int_{T-h}^T y'(t;v) dt \in L^2(\Omega).$$

So, it is necessary to introduce a process like (1.4) and (1.5), since we only have $\{z^0, z^1\} \in L^2(\Omega) \times H^{-1}(\Omega)$.

In what follows, we will suppose that there is a certain $h_0 > 0$, such that

$$0 < h \leq h_0, \quad T(x^0) < T - h_0 \quad (1.7)$$

and observing the linearity of the system (1.1), we can suppose without restricting the generality that

$$y^0 = y^1 = 0 \quad \text{in } \Omega. \quad (1.8)$$

Theorem 1.1. For any $\{z_h^0, z_h^1\} \in H_0^1(\Omega) \times L^2(\Omega)$, there exists a unique v_h which verifies (1.4), and we have

$$\|v_h\|_{L^2(\Sigma(x^0))} \leq C \left(\|z_h^0\|_{H_0^1(\Omega)}^2 + \|z_h^1\|_{L^2(\Omega)}^2 \right)^{1/2}, \quad (1.9)$$

where the constant C is independent of h .

Lemma 1.1. When $h \rightarrow 0$, if $\{z_h^0, z_h^1\} \in H_0^1(\Omega) \times L^2(\Omega)$ satisfies

$$\{z_h^0, z_h^1\} \rightharpoonup \{z^0, z^1\} \quad \text{weakly in } L^2(\Omega) \times H^{-1}(\Omega), \quad (1.10)$$

and if $\{v_h\}_{h>0}$ is bounded in $L^2(\Sigma(x^0))$, then

$$\begin{aligned} v_h &\rightharpoonup v_0 && \text{weakly in } L^2(\Sigma(x^0)), \\ y_h &\overset{*}{\rightharpoonup} y_0 && \text{weak } * \text{ in } L^\infty(0, T; L^2(\Omega)) \cap W^{1,\infty}(0, T; H^{-1}(\Omega)), \end{aligned} \quad (1.11)$$

where $y_h = y(v_h)$, $y_0 = y(v_0)$.

For the strong convergence of $\{v_h\}_{h>0}$, here, we give a construction of the series

$$\{z_h^0, z_h^1\}_{h>0} \subset H_0^1(\Omega) \times L^2(\Omega)$$

for any $\{z^0, z^1\} \in L^2(\Omega) \times H^{-1}(\Omega)$. In fact, we can choose

$$\bar{z}_h^0 = \frac{1}{h} \int_{T-h}^T z_h(t) dt, \quad \bar{z}_h^1 = \frac{1}{h} \int_{T-h}^T z_h'(t) dt, \quad (1.12)$$

where $z_h = z_h(x, t)$ satisfies

$$\begin{cases} z_h'' - \Delta z_h = 0 & \text{in } \Omega \times (0, h), \\ z_h = 0 & \text{on } (\Gamma \setminus \Gamma(x^0)) \times (0, h), \\ \int_0^h z_h dt = 0 & \text{on } \Gamma(x^0), \\ \int_{\Gamma(x^0) \times (0, h)} z_h^2 d\Gamma dt \rightarrow 0 & \text{when } h \rightarrow 0, \\ \{z_h(0), z_h'(0)\} \rightarrow \{z^0, z^1\} & \text{strongly in } L^2(\Omega) \times H^{-1}(\Omega), \text{ as } h \rightarrow 0. \end{cases} \quad (1.13)$$

We have the following results.

Theorem 1.2. *If $\{z_h^0, z_h^1\}$ satisfies*

$$\{z_h^0 - \bar{z}_h^0, z_h^1 - \bar{z}_h^1\} \rightarrow \{0, 0\} \quad \text{strongly in } H_0^1(\Omega) \times L^2(\Omega) \quad (1.14)$$

for a certain $\{z_h\}_{h>0}$, which satisfies (1.13), then we have

$$\|v_h\|_{L^2(\Sigma(x^0))} \leq \|v_0\|_{L^2(\Sigma(x^0))} + o(1), \quad (1.15)$$

$$\{z_h^0, z_h^1\} \rightarrow \{z^0, z^1\} \quad \text{strongly in } L^2(\Omega) \times H^{-1}(\Omega). \quad (1.16)$$

Hence

$$\begin{aligned} v_h &\rightarrow v_0 \quad \text{strongly in } L^2(\Sigma(x^0)), \\ y_h &\rightarrow y_0 \quad \text{strongly in } L^\infty(0, T; L^2(\Omega)) \cap W^{1,\infty}(0, T; H^{-1}(\Omega)). \end{aligned} \quad (1.17)$$

Remark 1.2. Particularly, we can choose

$$z_h \equiv z \quad \text{in } \Omega \times (0, h), \quad (1.18)$$

where z satisfies

$$\begin{cases} z'' - \Delta z = 0 & \text{in } \Omega \times (0, h_0), \\ z = 0 & \text{on } \Gamma \times (0, h_0), \\ z(0) = z^0, z'(0) = z^1 & \text{on } \Omega \times \{0\}, \end{cases} \quad (1.19)$$

and we take

$$\{z_h^0, z_h^1\} = \{\bar{z}_h^0, \bar{z}_h^1\}, \quad (1.20)$$

or, when $\{z^0, z^1\} \in H_0^1(\Omega) \times L^2(\Omega)$, we can simply take for any $h > 0$,

$$\{z_h^0, z_h^1\} \equiv \{z^0, z^1\}. \quad (1.21)$$

§2 Proof of Theorem 1.1

Following the HUM method of J. L. Lions^[1,2], we consider the optimal control problem:

$$\min_{v \in \mathcal{U}_{ad,h}} J(v), \quad J(v) = \frac{1}{2} \int_{\Sigma(x^0)} v^2 d\Sigma. \quad (2.1)$$

If

$$\mathcal{U}_{ad,h} \neq \emptyset, \quad (2.2)$$

then (2.1) admits a unique solution $v_h \in \mathcal{U}_{ad,h}$:

$$J(v_h) = \min_{v \in \mathcal{U}_{ad,h}} J(v). \quad (2.3)$$

Then in order to characterize v_h by an optimal system, we use the penalization method

$$J_\epsilon(v, z) = \frac{1}{2} \int_{\Sigma(x^0)} v^2 d\Sigma + \frac{1}{2\epsilon} |z'' - \Delta z|_{L^2(Q)}^2, \quad (2.4)$$

where v and z satisfy

$$\begin{aligned} v &\in L^2(\Sigma(x^0)), z'' - \Delta z \in L^2(Q), \\ z &= \begin{cases} v & \text{on } \Sigma(x^0), \\ 0 & \text{on } \Sigma \setminus \Sigma(x^0), \end{cases} \\ \frac{1}{h} \int_{T-h}^T z(t) dt &= z_h^0, \quad \frac{1}{h} \int_{T-h}^T z'(t) dt = z_h^1, \\ z(0) = z'(0) &= 0 \quad \text{on } \Omega \times \{0\}. \end{aligned} \quad (2.5)$$

We consider the following problem:

$$\min J_\epsilon(v, z), \quad (2.6)$$

where v and z satisfy (2.5).

According to (2.2), this problem admits a unique solution:

$$v_{h\epsilon}, \quad y_{h\epsilon}. \quad (2.7)$$

When $\epsilon \rightarrow 0$, we can prove that (cf. [1])

$$\begin{aligned} v_{h\epsilon} &\rightarrow v_h \quad \text{strongly in } L^2(\Sigma(x^0)), \\ y_{h\epsilon} &\rightarrow y_h \quad \text{strongly in } L^\infty(0, T; L^2(\Omega)) \cap W^{1,\infty}(0, T; H^{-1}(\Omega)). \end{aligned} \quad (2.8)$$

We define

$$P_{h\epsilon} = \begin{cases} P_{1h\epsilon} = -\frac{1}{\epsilon}(y_{h\epsilon}'' - \Delta y_{h\epsilon}) & \text{in } \Omega \times (T-h, T), \\ P_{2h\epsilon} = -\frac{1}{\epsilon}(y_{h\epsilon}'' - \Delta y_{h\epsilon}) & \text{in } \Omega \times (0, T-h). \end{cases} \quad (2.9)$$

Then $P_{h\epsilon} \in L^2(Q)$, and for any ζ, z which satisfy

$$\begin{aligned} \zeta &\in L^2(\Sigma(x^0)), z'' - \Delta z \in L^2(Q), \\ z &= \begin{cases} \zeta & \text{on } \Sigma(x^0), \\ 0 & \text{on } \Sigma \setminus \Sigma(x^0), \end{cases} \\ \frac{1}{h} \int_{T-h}^T z(t) dt &= 0, \quad \frac{1}{h} \int_{T-h}^T z'(t) dt = 0, \\ z(0) = z'(0) &= 0 \quad \text{on } \Omega \times \{0\}, \end{aligned} \quad (2.10)$$

we have

$$\int_Q P_{h\epsilon}(z'' - \Delta z) dx dt - \int_{\Sigma(x^0)} v_{h\epsilon} \zeta d\Sigma = 0. \quad (2.11)$$

If $P_{1h\epsilon}, P_{2h\epsilon}$ are regular enough, then $P_{1h\epsilon}, P_{2h\epsilon}$ will satisfy the following equation:

$$\begin{cases} \rho_1'' - \Delta \rho_1 = -\phi^1(x) & \text{in } \Omega \times (T-h, T), \\ \rho_1 = 0 & \text{on } \Gamma \times (T-h, T), \\ \rho_1(T) = 0, \rho_1'(T) = \phi^0(x) & \text{on } \Omega \times \{T\}, \end{cases} \quad (2.12)$$

$$\begin{cases} \rho_2'' - \Delta \rho_2 = 0 & \text{in } \Omega \times (0, T-h), \\ \rho_2 = 0 & \text{on } \Gamma \times (0, T-h), \\ \rho_2(T-h) = \rho_1(T-h) & \text{on } \Omega \times \{T-h\}, \\ \rho_2'(T-h) = \rho_1'(T-h) - \phi^0(x) & \text{on } \Omega \times \{T-h\}. \end{cases} \quad (2.13)$$

with $\{\phi^0, \phi^1\} = \{\phi_{h\epsilon}^0, \phi_{h\epsilon}^1\}$ for a certain $\{\phi_{h\epsilon}^0, \phi_{h\epsilon}^1\} \in L^2(\Omega) \times H^{-1}(\Omega)$.

Remark 2.1. We will denote the solutions ρ_1, ρ_2 of (2.12) and (2.13) by $\rho_1 = \rho_1(\phi^0, \phi^1)$, $\rho_2 = \rho_2(\phi^0, \phi^1)$, when we want to indicate their relations with $\{\phi^0, \phi^1\}$.

We have the following result, which says that $P_{1h\epsilon}$, $P_{2h\epsilon}$ are, in fact, regular enough, and which will be proved in §5.

Lemma 2.1. For any $w \in L^2(\Sigma(x^0))$, $P \in L^2(Q)$, if w, P satisfy

$$\begin{aligned} \int_{\Omega \times (0, T)} P(z'' - \Delta z) dx dt - \int_{\Sigma(x^0)} w \zeta d\Sigma &= 0, \\ \int_{T-h}^T w(t) dt &= 0 \quad \text{on } \Gamma(x^0) \end{aligned} \quad (2.14)$$

for any ζ, z which satisfy (2.10), then there exists a $\{\phi^0, \phi^1\} \in L^2(\Omega) \times H^{-1}(\Omega)$, such that

$$P = \begin{cases} \rho_1(\phi^0, \phi^1) & \text{in } \Omega \times (T-h, T), \\ \rho_2(\phi^0, \phi^1) & \text{in } \Omega \times (0, T-h), \end{cases} \quad (2.15)$$

$$w = \bar{v}(\phi^0, \phi^1) \stackrel{\text{def}}{=} \begin{cases} \frac{\partial \rho_1}{\partial \nu} - \frac{1}{h} \int_{T-h}^T \frac{\partial \rho_1}{\partial \nu} dt & \text{on } \Gamma(x^0) \times (T-h, T), \\ \frac{\partial \rho_2}{\partial \nu} & \text{on } \Gamma(x^0) \times (0, T-h). \end{cases} \quad (2.16)$$

So, there really exists a $\{\phi_{h\epsilon}^0, \phi_{h\epsilon}^1\} \in L^2(\Omega) \times H^{-1}(\Omega)$, such that

$$P_{1h\epsilon} = \rho_1(\phi_{h\epsilon}^0, \phi_{h\epsilon}^1), P_{2h\epsilon} = \rho_2(\phi_{h\epsilon}^0, \phi_{h\epsilon}^1), v_{h\epsilon} = \bar{v}(\phi_{h\epsilon}^0, \phi_{h\epsilon}^1). \quad (2.17)$$

By the results of §4, we have

$$\begin{aligned} C_h \|\bar{v}(\phi^0, \phi^1)\|_{L^2(\Sigma(x^0))}^2 &\leq |\phi^0|_{L^2(\Omega)}^2 + \|\phi\|_{H^{-1}(\Omega)}^2 \\ &\leq C\left(\frac{1}{h} + \frac{1}{h^2}\right) \|\bar{v}(\phi^0, \phi^1)\|_{L^2(\Sigma(x^0))}^2, \end{aligned} \quad (2.18)$$

where C_h is a constant dependent of h .

Combining (2.8), (2.17) and (2.18), we can deduce that there exists a $\{\bar{\phi}_h^0, \bar{\phi}_h^1\} \in L^2(\Omega) \times H^{-1}(\Omega)$, such that

$$\begin{aligned} \phi_{h\epsilon}^0 &\rightarrow \bar{\phi}_h^0 & \text{strongly in } L^2(\Omega), \\ \phi_{h\epsilon}^1 &\rightarrow \bar{\phi}_h^1 & \text{strongly in } H^{-1}(\Omega), \\ v_h &= \bar{v}(\bar{\phi}_h^0, \bar{\phi}_h^1). \end{aligned} \quad (2.19)$$

We have proved that if v_h exists, that is to say, if (2.2) is true, then v_h will verifies (2.19). Now we are going to verify (2.2) by the HUM Method.

Firstly, for any $\{\phi^0, \phi^1\} \in L^2(\Omega) \times H^{-1}(\Omega)$, we solve

$$\begin{aligned} \psi'' - \Delta \psi &= 0 & \text{in } Q, \\ \psi &= \bar{v}(\phi^0, \phi^1) = \begin{cases} \frac{\partial \rho_1}{\partial \nu} - \frac{1}{h} \int_{T-h}^T \frac{\partial \rho_1}{\partial \nu} dt & \text{on } \Gamma(x^0) \times (T-h, T), \\ \frac{\partial \rho_2}{\partial \nu} & \text{on } \Gamma(x^0) \times (0, T-h), \end{cases} \\ \psi &= 0 & \text{on } \Sigma \setminus \Sigma(x^0), \\ \psi(0) &= \psi'(0) = 0 & \text{on } \Omega \times \{0\}, \end{aligned} \quad (2.20)$$

where $\rho_1 = \rho_1(\phi^0, \phi^1)$, $\rho_2 = \rho_2(\phi^0, \phi^1)$. We set

$$\begin{aligned} u^0(x) &= -\Delta^{-1} \phi^1(x) \in H_0^1(\Omega), & u^1(x) &= \phi^0(x) \in L^2(\Omega), \\ u(x, t) &= \rho_1(x, t) + u^0(x), & \text{in } \Omega \times (T-h, T). \end{aligned} \quad (2.21)$$

u will verify

$$\begin{cases} u'' - \Delta u = 0 & \text{in } \Omega \times (T-h, T), \\ u = 0 & \text{on } \Gamma \times (T-h, T), \\ u(T) = u^0, u'(T) = u^1 & \text{on } \Omega \times \{T\}. \end{cases} \quad (2.22)$$

By the results of J. L. Lions^[1,2], we know that

$$\frac{\partial u}{\partial \nu} \in L^2(\Gamma \times (T-h, T)), \quad \frac{\partial \rho_2}{\partial \nu} \in L^2(\Gamma \times (0, T-h)).$$

Hence

$$\frac{\partial \rho_1}{\partial \nu} - \frac{1}{h} \int_{T-h}^T \frac{\partial \rho_1}{\partial \nu} dt = \frac{\partial u}{\partial \nu} - \frac{1}{h} \int_{T-h}^T \frac{\partial u}{\partial \nu} dt \in L^2(\Gamma \times (T-h, T)). \quad (2.23)$$

Then, we can define

$$\begin{aligned} \Lambda_h : L^2(\Omega) \times H^{-1}(\Omega) &\rightarrow L^2(\Omega) \times H_0^1(\Omega), \\ \{\phi^0, \phi^1\} &\mapsto \left\{ \frac{1}{h} \int_{T-h}^T \psi' dt, \frac{1}{h} \int_{T-h}^T \psi dt \right\}, \end{aligned} \quad (2.24)$$

where $\psi = \psi(\phi^0, \phi^1)$ is the solution of (2.20) associated with the initial data $\{\phi^0, \phi^1\}$.

Now, we consider another couple $\{\bar{\phi}^0, \bar{\phi}^1\} \in L^2(\Omega) \times H^{-1}(\Omega)$, and $\bar{\rho}_1 = \rho_1(\bar{\phi}^0, \bar{\phi}^1)$, $\bar{\rho}_2 = \rho_2(\bar{\phi}^0, \bar{\phi}^1)$; we multiply (2.20) by $\bar{\rho}_1, \bar{\rho}_2$; at the end, we will get

$$\begin{aligned} &\left(\frac{\psi(T) - \psi(T-h)}{h}, \bar{\phi}^0 \right)_{L^2(\Omega), L^2(\Omega)} + \left\langle \frac{1}{h} \int_{T-h}^T \psi dt, \bar{\phi}^1 \right\rangle_{H_0^1(\Omega), H^{-1}(\Omega)} \\ &= \frac{1}{h} \int_{\Gamma(x^0) \times (T-h, T)} \left(\frac{\partial \rho_1}{\partial \nu} - \frac{1}{h} \int_{T-h}^T \frac{\partial \rho_1}{\partial \nu} dt \right) \frac{\partial \bar{\rho}_1}{\partial \nu} d\Sigma + \frac{1}{h} \int_{\Gamma(x^0) \times (0, T-h)} \frac{\partial \rho_2}{\partial \nu} \frac{\partial \bar{\rho}_2}{\partial \nu} d\Sigma \\ &= \frac{1}{h} \int_{\Gamma(x^0) \times (T-h, T)} \left(\frac{\partial \rho_1}{\partial \nu} - \frac{1}{h} \int_{T-h}^T \frac{\partial \rho_1}{\partial \nu} dt \right) \left(\frac{\partial \bar{\rho}_1}{\partial \nu} - \frac{1}{h} \int_{T-h}^T \frac{\partial \bar{\rho}_1}{\partial \nu} dt \right) d\Sigma \\ &\quad + \frac{1}{h} \int_{\Gamma(x^0) \times (0, T-h)} \frac{\partial \rho_2}{\partial \nu} \frac{\partial \bar{\rho}_2}{\partial \nu} d\Sigma. \end{aligned} \quad (2.25)$$

This means

$$\begin{aligned} &\langle \Lambda_h \{\phi^0, \phi^1\}, \{\bar{\phi}^0, \bar{\phi}^1\} \rangle_{L^2(\Omega) \times H_0^1(\Omega), L^2(\Omega) \times H^{-1}(\Omega)} \\ &= \frac{1}{h} \int_{\Sigma(x^0)} \bar{v}(\phi^0, \phi^1) \bar{v}(\bar{\phi}^0, \bar{\phi}^1) d\Sigma, \end{aligned} \quad (2.26)$$

hence

$$\Lambda_h = \Lambda_h^*. \quad (2.27)$$

By observing (2.18), we can deduce

$$\begin{aligned} &\langle \Lambda_h \{\phi^0, \phi^1\}, \{\phi^0, \phi^1\} \rangle_{L^2(\Omega) \times H_0^1(\Omega), L^2(\Omega) \times H^{-1}(\Omega)} \\ &= \frac{1}{h} \int_{\Sigma(x^0)} \bar{v}(\phi^0, \phi^1) \bar{v}(\phi^0, \phi^1) d\Sigma \geq C_h \|\{\phi^0, \phi^1\}\|_{L^2(\Omega) \times H^{-1}(\Omega)}^2. \end{aligned} \quad (2.28)$$

Hence, Λ_h is an isomorphism.

So for any $\{z_h^1, z_h^0\} \in L^2(\Omega) \times H_0^1(\Omega)$ we can find a $\{\phi_h^0, \phi_h^1\} \in L^2(\Omega) \times H^{-1}(\Omega)$, such that

$$\Lambda_h \{\phi_h^0, \phi_h^1\} = \{z_h^1, z_h^0\}. \quad (2.29)$$

we can deduce that

$$\bar{v}(\phi_h^0, \phi_h^1) \in \mathcal{U}_{ad,h}. \quad (2.30)$$

So, (2.2) is verified.

Since Λ_h is an isomorphism, we can see from (2.19), (2.30) that

$$\begin{aligned} \{\bar{\phi}_h^0, \bar{\phi}_h^1\} &= \{\phi_h^0, \phi_h^1\}, \\ v_h &= \bar{v}(\phi_h^0, \phi_h^1). \end{aligned} \quad (2.31)$$

From (2.18), we can deduce

$$\begin{aligned}
 \left(\|z_h^0\|_{H_0^1(\Omega)}^2 + |z_h^1|_{L^2(\Omega)}^2 \right)^{\frac{1}{2}} &= \|\Lambda_h \{\phi_h^0, \phi_h^1\}\|_{L^2(\Omega) \times H_0^1(\Omega)} \\
 &\geq \frac{\langle \Lambda_h \{\phi_h^0, \phi_h^1\}, \{\phi_h^0, \phi_h^1\} \rangle_{L^2(\Omega) \times H_0^1(\Omega), L^2(\Omega) \times H^{-1}(\Omega)}}{\|\{\phi_h^0, \phi_h^1\}\|_{L^2(\Omega) \times H^{-1}(\Omega)}} \\
 &\geq C \frac{\frac{1}{h} \|\bar{v}(\phi_h^0, \phi_h^1)\|_{L^2(\Sigma(x^0))}^2}{\left(\frac{1}{h} + \frac{1}{h^2}\right)^{\frac{1}{2}} \|\bar{v}(\phi_h^0, \phi_h^1)\|_{L^2(\Sigma(x^0))}}.
 \end{aligned} \tag{2.32}$$

So we get (1.9).

§3. Proofs of Lemma 1.1 and Theorem 1.2

Proof of Lemma 1.1. If for a certain $\bar{v}_0 \in L^2(\Sigma(x^0))$ we have a subseries such that

$$v_h \rightharpoonup \bar{v}_0 \quad \text{weakly in } L^2(\Sigma(x^0)), \tag{3.1}$$

by the results of J. L. Lions^[1,2] we know that the series $\{\rho_{2h}(T-h), \rho'_{2h}(T-h)\}$ remains bounded in $H_0^1(\Omega) \times L^2(\Omega)$, where $\rho_{2h} = \rho_2(\phi_h^0, \phi_h^1)$. So we can choose a subseries (still denoted by h) such that

$$\begin{aligned}
 \rho_{2h}(T-h) &\rightharpoonup \bar{\Phi}^0 && \text{weakly in } H_0^1(\Omega), \\
 \rho'_{2h}(T-h) &\rightharpoonup \bar{\Phi}^1 && \text{weakly in } L^2(\Omega)
 \end{aligned} \tag{3.2}$$

for a certain $\{\bar{\Phi}^0, \bar{\Phi}^1\} \in H_0^1(\Omega) \times L^2(\Omega)$, and

$$\begin{aligned}
 y_h &\xrightarrow{*} \bar{y}_0 && \text{weak } * \text{ in } L^\infty(0, T; L^2(\Omega)) \cap W^{1,\infty}(0, T; H^{-1}(\Omega)), \\
 \rho_{2h} &\xrightarrow{*} \bar{\Phi}_0 && \text{weak } * \text{ in } L^\infty(0, T; H_0^1(\Omega)) \cap W^{1,\infty}(0, T; L^2(\Omega)),
 \end{aligned} \tag{3.3}$$

where ρ_{2h} is extended into $\Omega \times (T-h, T)$ by

$$\begin{aligned}
 \rho_{2h}'' - \Delta \rho_{2h} &= 0 && \text{in } \Omega \times (T-h, T), \\
 \rho_{2h} &= 0 && \text{on } \Gamma \times (T-h, T),
 \end{aligned} \tag{3.4}$$

$\bar{y}_0 = y(\bar{v}_0)$, $\bar{\Phi}_0$ satisfies the following equation:

$$\begin{cases} \Phi'' - \Delta \Phi = 0 & \text{in } Q, \\ \Phi = 0 & \text{on } \Sigma, \\ \Phi(T) = \bar{\Phi}^0, \Phi'(T) = \bar{\Phi}^1 & \text{on } \Omega \times \{T\}, \end{cases} \tag{3.5}$$

with the initial data $\{\bar{\Phi}^0, \bar{\Phi}^1\} = \{\bar{\Phi}^0, \bar{\Phi}^1\}$, and still more, we have

$$\frac{\partial \bar{\Phi}_0}{\partial \nu} = \bar{v}_0 \quad \text{on } \Sigma(x^0). \tag{3.6}$$

It still remains to prove that

$$\bar{y}_0(T) = z^0, \quad \bar{y}'_0(T) = z^1. \tag{3.7}$$

For any $\{\bar{\Phi}^0, \bar{\Phi}^1\} \in H_0^1(\Omega) \times L^2(\Omega)$, and $\Phi = \Phi(\bar{\Phi}^0, \bar{\Phi}^1)$ the associated solution of (3.5), and for any $s \in (T-h, T)$, we have

$$\begin{aligned}
 &(y_h(s), \bar{\Phi}^1)_{L^2(\Omega), L^2(\Omega)} - \langle y'_h(s), \bar{\Phi}^0 \rangle_{H^{-1}(\Omega), H_0^1(\Omega)} \\
 &= \int_{\Gamma(x^0) \times (0, s)} v_h(t) \frac{\partial \Phi}{\partial \nu}(t+T-s) d\Sigma.
 \end{aligned} \tag{3.8}$$

Integrating (3.8) in $(T-h, T)$, we can get

$$\begin{aligned} & (z_h^0, \Phi^1)_{L^2(\Omega), L^2(\Omega)} - \langle z_h^1, \Phi^0 \rangle_{H^{-1}(\Omega), H_0^1(\Omega)} \\ &= \frac{1}{h} \int_{T-h}^T \int_{\Gamma(x^0) \times (0, s)} v_h(t) \frac{\partial \Phi}{\partial \nu}(t+T-s) d\Sigma ds. \end{aligned} \quad (3.9)$$

In (3.9) letting $h \rightarrow 0$, we will get

$$\begin{aligned} & (z^0, \Phi^1)_{L^2(\Omega), L^2(\Omega)} - \langle z^1, \Phi^0 \rangle_{H^{-1}(\Omega), H_0^1(\Omega)} \\ &= \int_{\Sigma(x^0) \times (0, T)} \bar{v}_0(t) \frac{\partial \Phi}{\partial \nu} d\Sigma = \int_{\Sigma(x^0) \times (0, T)} \frac{\partial \bar{\Phi}_0}{\partial \nu} \frac{\partial \Phi}{\partial \nu} d\Sigma. \end{aligned} \quad (3.10)$$

So we have (3.7). By the results of J. L. Lions^[1,2] we can deduce from (1.2), (3.6) and (3.7) that

$$\begin{aligned} \bar{y}_0 &= y_0 \quad \text{in } Q, \\ \bar{v}_0 &= v_0 \quad \text{on } \Sigma(x^0). \end{aligned} \quad (3.11)$$

Proof of Theorem 1.2. It suffices to prove (1.15). If $\bar{v}_h$ is such that

$$\frac{1}{h} \int_{T-h}^T \bar{y}_h(t) dt = \bar{z}_h^0, \quad \frac{1}{h} \int_{T-h}^T \bar{y}_h'(t) dt = \bar{z}_h^1, \quad (3.12)$$

where $\bar{y}_h = y(\bar{v}_h)$, and $\bar{v}_h$ is optimal in the following sense:

$$\begin{aligned} & \|\bar{v}_h\|_{L^2(\Sigma(x^0))} = \min \|v\|_{L^2(\Sigma(x^0))}, \\ & \text{for any } v \in L^2(\Sigma(x^0)) \quad \text{which makes (3.12) true with} \\ & y = y(v) \quad \text{being in the place of } \bar{y}_h, \end{aligned} \quad (3.13)$$

According to (1.9) we have

$$\|v_h - \bar{v}_h\|_{L^2(\Sigma(x^0))} \leq C \left(\|z_h^0 - \bar{z}_h^0\|_{H_0^1(\Omega)}^2 + \|z_h^1 - \bar{z}_h^1\|_{L^2(\Omega)}^2 \right)^{\frac{1}{2}} = o(1). \quad (3.14)$$

If $\bar{v}_0 \in L^2(\Sigma(x^0) \times (0, T-h))$ is such that:

$$\bar{y}_0(T-h) = z_h(0), \quad \bar{y}_0'(T-h) = z_h'(0), \quad (3.15)$$

where $\bar{y}_0 = y(\bar{v}_0)$ is the solution of (1.1) in $\Omega \times (0, T-h)$ with $v = \bar{v}_0$ on $\Gamma(x^0) \times (0, T-h)$, and $\bar{v}_0$ is optimal in the following sense:

$$\begin{aligned} & \|\bar{v}_0\|_{L^2(\Gamma(x^0) \times (0, T-h))} = \min \|v\|_{L^2(\Gamma(x^0) \times (0, T-h))}, \\ & \text{for any } v \in L^2(\Gamma(x^0) \times (0, T-h)) \quad \text{which makes (3.15) true with} \\ & y = y(v) \quad \text{being in the place of } \bar{y}_0, \end{aligned} \quad (3.16)$$

we extend $\bar{v}_0$ into $\Gamma(x^0) \times (T-h, T)$ by:

$$\bar{v}_0(t) = z_h(t-T+h) \quad \text{on } \Gamma(x^0) \times (T-h, T). \quad (3.17)$$

Then, by the results of J. L. Lions^[1,2], we have

$$\begin{aligned} \|v_0 - \bar{v}_0\|_{L^2(\Sigma(x^0))} &\leq \|v_0 - \bar{v}_0\|_{L^2(\Gamma(x^0) \times (0, T-h))} + \|v_0 - \bar{v}_0\|_{L^2(\Gamma(x^0) \times (T-h, T))} \\ &\leq C \left(\|y(T-h; v_0) - z_h(0)\|_{L^2(\Omega)}^2 + \|y'(T-h; v_0) - z_h'(0)\|_{H^{-1}(\Omega)}^2 \right)^{\frac{1}{2}} \\ &\quad + \|v_0\|_{L^2(\Gamma(x^0) \times (T-h, T))} + \|z_h\|_{L^2(\Gamma(x^0) \times (0, h))} \\ &= o(1). \end{aligned} \quad (3.18)$$

Since $\bar{v}_0$ verifies (3.12), we have from (3.13)

$$\|\bar{v}_h\|_{L^2(\Sigma(x^0))} \leq \|\bar{v}_0\|_{L^2(\Sigma(x^0))}. \quad (3.19)$$

So we have

$$\begin{aligned} \|v_h\|_{L^2(\Sigma(x^0))} &\leq \|v_h - \bar{v}_h\|_{L^2(\Sigma(x^0))} + \|\bar{v}_h\|_{L^2(\Sigma(x^0))} \leq o(1) + \|\bar{v}_0\|_{L^2(\Sigma(x^0))} \\ &\leq \|v_0\|_{L^2(\Sigma(x^0))} + \|v_0 - \bar{v}_0\|_{L^2(\Sigma(x^0))} + o(1) \\ &\leq \|v_0\|_{L^2(\Sigma(x^0))} + o(1). \end{aligned} \quad (3.20)$$

§4. The Inverse Inequality

In this section we are going to verify (2.18). Since

$$\begin{aligned} \|\bar{v}(\phi^0, \phi^1)\|_{L^2(\Sigma(x^0))}^2 &= \int_{\Gamma(x^0) \times (T-h, T)} \left(\frac{\partial \rho_1}{\partial \nu} - \frac{1}{h} \int_{T-h}^T \frac{\partial \rho_1}{\partial \nu} dt \right)^2 d\Sigma \\ &\quad + \int_{\Gamma(x^0) \times (0, T-h)} \left(\frac{\partial \rho_2}{\partial \nu} \right)^2 d\Sigma \end{aligned} \quad (4.1)$$

by the results of J. L. Lions^[1,2], we know that (cf.(2.21))

$$\begin{aligned} \int_{\Gamma(x^0) \times (0, T-h)} \left(\frac{\partial \rho_2}{\partial \nu} \right)^2 d\Sigma &\geq C \left(\|\rho_2(T-h)\|_{H_0^1(\Omega)}^2 + |\rho_2'(T-h)|_{L^2(\Omega)}^2 \right) \\ &= C \left(\|\rho_1(T-h)\|_{H_0^1(\Omega)}^2 + |\rho_1'(T-h) - \phi^0|_{L^2(\Omega)}^2 \right) \\ &= C \left(\|u(T) - u(T-h)\|_{H_0^1(\Omega)}^2 + |u'(T) - u'(T-h)|_{L^2(\Omega)}^2 \right). \end{aligned} \quad (4.2)$$

Hence (cf. (2.23))

$$\begin{aligned} \|\bar{v}(\phi^0, \phi^1)\|_{L^2(\Sigma(x^0))}^2 &\geq \int_{\Gamma(x^0) \times (T-h, T)} \left(\frac{\partial u}{\partial \nu} - \frac{1}{h} \int_{T-h}^T \frac{\partial u}{\partial \nu} dt \right)^2 d\Sigma \\ &\quad + C \left(\|u(T) - u(T-h)\|_{H_0^1(\Omega)}^2 + |u'(T) - u'(T-h)|_{L^2(\Omega)}^2 \right). \end{aligned} \quad (4.3)$$

So, it suffices to prove the following results.

Theorem 4.1. For any $\{u^0, u^1\} \in H_0^1(\Omega) \times L^2(\Omega)$, we have

$$\begin{aligned} \|u^0\|_{H_0^1(\Omega)}^2 + |u^1|_{L^2(\Omega)}^2 &\leq C \left(\frac{1}{h} \int_{\Gamma(x^0) \times (0, h)} \left(\frac{\partial u}{\partial \nu} - \frac{1}{h} \int_0^h \frac{\partial u}{\partial \nu} dt \right)^2 d\Sigma \right. \\ &\quad \left. + \frac{1}{h^2} \|u(h) - u(0)\|_{H_0^1(\Omega)}^2 + \frac{1}{h^2} |u'(h) - u'(0)|_{L^2(\Omega)}^2 \right), \end{aligned} \quad (4.4)$$

where $u = u(x, t)$ satisfies

$$\begin{cases} u'' - \Delta u = 0 & \text{in } \Omega \times (0, h), \\ u = 0 & \text{on } \Gamma \times (0, h), \\ u(0) = u^0, u'(0) = u^1 & \text{on } \Omega \times \{0\}. \end{cases} \quad (4.5)$$

Proof. By the results of J. L. Lions^[1,2], we have

$$\frac{\partial u}{\partial \nu} \in L^2(\Gamma \times (0, h)), \quad (4.6)$$

$$Z(t) \Big|_0^h + hE(0) \leq \frac{R(x^0)}{2} \int_{\Gamma(x^0) \times (0, h)} \left| \frac{\partial u}{\partial \nu} \right|^2 d\Sigma, \quad (4.7)$$

where

$$\begin{aligned} E(0) &= \frac{1}{2} \int_{\Omega} (|\nabla u^0|^2 + |u^1|^2) dx, \\ Z(t) &= (u'(t), m \cdot \nabla u(t) + \frac{n-1}{2} u(t))_{L^2(\Omega), L^2(\Omega)}. \end{aligned} \quad (4.8)$$

For any $0 < \epsilon \ll 1$ we will denote by C_ϵ some corresponding sufficiently small constants.

$$\begin{aligned} |Z(t)|_0^h &= |(u'(h) - u(0), m \cdot \nabla u(h) + \frac{n-1}{2} u(h))_{L^2(\Omega), L^2(\Omega)} \\ &\quad + (u'(0), m \cdot \nabla u(h) + \frac{n-1}{2} u(h) - m \cdot \nabla u(0) - \frac{n-1}{2} u(0))_{L^2(\Omega), L^2(\Omega)}| \quad (4.9) \\ &\leq \epsilon h E(0) + C_\epsilon \frac{1}{h} \left(\|u(h) - u(0)\|_{H_0^1(\Omega)}^2 + |u'(h) - u'(0)|_{L^2(\Omega)}^2 \right). \end{aligned}$$

Hence

$$\begin{aligned} E(0) &\leq C \left(\frac{1}{h} \int_{\Gamma(x^0) \times (0, h)} \left(\frac{\partial u}{\partial \nu} \right)^2 d\Sigma + \frac{1}{h^2} \|u(h) - u(0)\|_{H_0^1(\Omega)}^2 \right. \\ &\quad \left. + \frac{1}{h^2} |u'(h) - u'(0)|_{L^2(\Omega)}^2 \right). \end{aligned} \quad (4.10)$$

Since we have

$$\begin{aligned} &\frac{1}{h} \int_{\Gamma(x^0) \times (0, h)} \left(\frac{\partial u}{\partial \nu} \right)^2 d\Sigma \\ &= \frac{1}{h} \int_{\Gamma(x^0) \times (0, h)} \left(\frac{\partial u}{\partial \nu} - \frac{1}{h} \int_0^h \frac{\partial u}{\partial \nu} dt \right)^2 d\Sigma + \frac{1}{h^2} \int_{\Gamma(x^0)} \left(\int_0^h \frac{\partial u}{\partial \nu} dt \right)^2 d\Gamma, \end{aligned} \quad (4.11)$$

it suffices to verify

$$\int_{\Gamma(x^0)} \left(\int_0^h \frac{\partial u}{\partial \nu} dt \right)^2 d\Gamma \leq \epsilon h^2 E(0) + C_\epsilon (\|u(h) - u(0)\|_{H_0^1(\Omega)}^2 + |u'(h) - u'(0)|_{L^2(\Omega)}^2). \quad (4.12)$$

There exists (cf. Lions [1,2]) a vector field $q(x) = (q_1(x), q_2(x), \dots, q_n(x)) \in (C^2(\bar{\Omega}))^n$, such that

$$q(x) = \nu(x) \quad \text{on } \Gamma. \quad (4.13)$$

We multiply (4.5) by $q_k \int_0^h \frac{\partial u}{\partial x_k} dt$ ($:= \sum_{k=1}^n q_k \int_0^h \frac{\partial u}{\partial x_k} dt$, (we will always omit the $\sum$ notation for simplicity)

$$\begin{aligned} 0 &= \int_{\Omega \times (0, h)} q_k \left(\int_0^h \frac{\partial u}{\partial x_k} dt \right) (u'' - \Delta u) dx dt \\ &= \int_{\Omega} q_k \left(\int_0^h \frac{\partial u}{\partial x_k} dt \right) u'(t) dx \Big|_0^h - \int_{\Gamma \times (0, h)} q_k \left(\int_0^h \frac{\partial u}{\partial x_k} dt \right) \frac{\partial u}{\partial \nu} d\Gamma dt \\ &\quad + \int_{\Omega \times (0, h)} \frac{\partial q_k}{\partial x_j} \left(\int_0^h \frac{\partial u}{\partial x_k} dt \right) \frac{\partial u}{\partial x_j} dx dt + \int_{\Omega \times (0, h)} q_k \left(\int_0^h \frac{\partial^2 u}{\partial x_j \partial x_k} dt \right) \frac{\partial u}{\partial x_j} dx dt. \end{aligned} \quad (4.14)$$

For the last term, we have

$$\begin{aligned} \int_{\Omega \times (0, h)} q_k \left(\int_0^h \frac{\partial^2 u}{\partial x_j \partial x_k} dt \right) \frac{\partial u}{\partial x_j} dx dt &= \int_{\Omega} \frac{1}{2} q_k \frac{\partial}{\partial x_k} \left(\int_0^h \frac{\partial u}{\partial x_j} dt \right)^2 dx \\ &= \int_{\Omega} \frac{1}{2} q_k \nu_k \left(\int_0^h \frac{\partial u}{\partial x_j} dt \right)^2 dx - \int_{\Omega} \frac{1}{2} \frac{\partial q_k}{\partial x_k} \left(\int_0^h \frac{\partial u}{\partial x_j} dt \right)^2 dx. \end{aligned} \quad (4.15)$$

Since $u = 0$, on $\Gamma \times (0, h)$, we have

$$\nabla u = \frac{\partial u}{\partial \nu} \nu \quad \text{on } \Gamma \times (0, h). \quad (4.16)$$

By combining (4.14), (4.15) and (4.16), we can see that

$$\begin{aligned} &\frac{1}{2} \int_{\Gamma} \left(\int_0^h \frac{\partial u}{\partial \nu} dt \right)^2 d\Gamma \\ &= \int_{\Omega} q_k \left(\int_0^h \frac{\partial u}{\partial x} dt \right) u'(t) dx \Big|_0^h + \int_{\Omega} \frac{\partial q_k}{\partial x_j} \left(\int_0^h \frac{\partial u}{\partial x_k} dt \right) \left(\int_0^h \frac{\partial u}{\partial x_j} dt \right) dx \\ &\quad - \int_{\Omega} \frac{1}{2} \frac{\partial q_k}{\partial x_k} \left(\int_0^h \frac{\partial u}{\partial x_j} dt \right)^2 dx \\ &\leq C(|u'(h) - u'(0)|_{L^2(\Omega)}^2 + \int_{\Omega} \left(\int_0^h \frac{\partial u}{\partial x_k} dt \right)^2 dx). \end{aligned} \quad (4.17)$$

But, we have

$$\Delta \left(\int_0^h u dt \right) = u'(h) - u'(0). \quad (4.18)$$

Hence

$$\left\| \int_0^h u dt \right\|_{H_0^1(\Omega)}^2 \leq C|u'(h) - u'(0)|_{L^2(\Omega)}^2. \quad (4.19)$$

Then, by combining (4.17), (4.19), we can obtain (4.12).

Corollary 4.1 (Uniqueness Theorem). *If u is a solution of (4.5), which satisfies*

$$u(h) = u(0), \quad u'(h) = u'(0) \text{ in } \Omega, \quad (4.20)$$

$$\frac{\partial u}{\partial \nu} \quad \text{is independent of } t \text{ on } \Gamma(x^0) \times (0, h), \quad (4.21)$$

then

$$u \equiv 0 \quad \text{in } \Omega \times (0, h). \quad (4.22)$$

§5. Proof of Lemma 2.1

Firstly, we can observe that the propriety that Λ_h is an isomorphism is independent of this proof, so we can use it here.

Using this $w \in L^2(\Sigma(x^0))$ in (1.1), we can get $y_w = y(w)$, the solution of (1.1) associated with w , and

$$\left\{ \frac{1}{h} \int_{T-h}^T y'_w(t) dt, \quad \frac{1}{h} \int_{T-h}^T y_w(t) dt \right\} \in L^2(\Omega) \times H_0^1(\Omega). \quad (5.1)$$

So, there exists a unique $\{\phi^0, \phi^1\} \in L^2(\Omega) \times H^{-1}(\Omega)$, such that

$$\Lambda_h\{\phi^0, \phi^1\} = \left\{ \frac{1}{h} \int_{T-h}^T y'_w(t) dt, \quad \frac{1}{h} \int_{T-h}^T y_w(t) dt \right\}. \quad (5.2)$$

We set

$$\begin{aligned} \rho &= \rho(x, t) = \rho(x, t; \phi^0, \phi^1) \\ &= \rho(\phi^0, \phi^1) = \begin{cases} \rho_1(\phi^0, \phi^1) & \text{on } \Omega \times (T-h, T), \\ \rho_2(\phi^0, \phi^1) & \text{on } \Omega \times (0, T-h). \end{cases} \end{aligned} \quad (5.3)$$

Then, we have

$$\int_Q (P - \rho)(z'' - \Delta z) dx dt - \int_{\Sigma(x^0)} (w - \bar{v}(\phi^0, \phi^1)) \zeta d\Sigma = 0, \quad (5.4)$$

for any ζ, z which satisfy (2.10).

We can take in (5.4)

$$\begin{aligned} z &= y_w - \psi(\phi^0, \phi^1), \\ \zeta &= w - \bar{v}(\phi^0, \phi^1). \end{aligned} \quad (5.5)$$

We can deduce that

$$w = \bar{v}(\phi^0, \phi^1) \quad \text{on } \Sigma(x^0). \quad (5.6)$$

So we have

$$\int_Q (P - \rho)(z'' - \Delta z) dx dt = 0, \quad (5.7)$$

for any z which satisfies (2.10). For any $f \in L^2(Q)$, we solve the following problem:

$$\begin{cases} z_1'' - \Delta z_1 = f & \text{in } Q, \\ z_1 = 0 & \text{on } \Sigma, \\ z_1(0) = 0, z_1'(0) = 0 & \text{on } \Omega \times \{0\}. \end{cases} \quad (5.8)$$

There exists a $\{\phi_f^0, \phi_f^1\} \in L^2(\Omega) \times H^{-1}(\Omega)$, such that

$$\Lambda_h\{\phi_f^0, \phi_f^1\} = \left\{ \frac{1}{h} \int_{T-h}^T z_1'(t) dt, \quad \frac{1}{h} \int_{T-h}^T z_1(t) dt \right\}. \quad (5.9)$$

We can see that

$$z_f = z_1 - \psi(\phi_f^0, \phi_f^1) \quad (5.10)$$

verifies (2.10). By submitting z_f into (5.7), we get

$$\int_Q (P - \rho) f dx dt = 0, \quad (5.11)$$

for any $f \in L^2(Q)$. So, we obtain

$$P \equiv \rho \quad \text{in } Q. \quad (5.12)$$

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