

Bochner-Martinelli Formula for Higher Spin Operators of Several \mathbb{R}^6 Variables*

Guangzhen REN¹ Qianqian KANG¹

Abstract The higher spin operator of several \mathbb{R}^6 variables is an analogue of the $\bar{\partial}$ -operator in theory of several complex variables. The higher spin representation of $\mathfrak{so}_6(\mathbb{C})$ is $\odot^k \mathbb{C}^4$ and the higher spin operator \mathcal{D}_k acts on $\odot^k \mathbb{C}^4$ -valued functions. In this paper, the authors establish the Bochner-Martinelli formula for higher spin operator \mathcal{D}_k of several \mathbb{R}^6 variables. The embedding of \mathbb{R}^{6n} into the space of complex $4n \times 4$ matrices allows them to use two-component notation, which makes the spinor calculus on \mathbb{R}^{6n} more concrete and explicit. A function annihilated by \mathcal{D}_k is called k -monogenic. They give the Penrose integral formula over \mathbb{R}^{6n} and construct many k -monogenic polynomials.

Keywords Higher spin operator, k -Monogenic, Bochner-Martinelli formula, Penrose integral formula

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1 Introduction

After Pertici [17] proved the Hartogs' phenomenon for several quaternionic variables, people began to study the regular function theory of several quaternionic variables, generalizing the theory of several complex variables. So far the k -Cauchy-Fueter complex, the quaternionic counterpart of Dolbeault complex is known explicitly (see [1–2, 5–6] for $k = 1$, [23] and references there in).

More generally, people are interested in constructing the corresponding differential complex for several \mathbb{R}^n variables. Such complexes exist in the stable case, i.e., the number of variables is less than or equal to $\frac{n}{2}$, see [7, 15, 20]. It is interesting to develop the function of regular functions of several \mathbb{R}^n variables. Ren and Wang [19] investigated several octonionic variables.

On the other hands, physicists are interested in 6-dimensional physics, in particular, the superconformal field theory over space-time $\mathbb{R}^{5,1}$ and the 6-dimensional theory of self-dual three-forms, the reduction of which gives us the self-dual string equations in 4-dimensions (see e.g. [14, 16, 21] and references there in). It is well known that the orthogonal Lie algebra is isomorphic to the special linear algebra $\mathfrak{sl}_4(\mathbb{C})$, i.e.,

$$\mathfrak{so}_6(\mathbb{C}) \cong \mathfrak{sl}_4(\mathbb{C}), \quad (1.1)$$

see [11, P. 327]. Then the basic spin representation of $\mathfrak{so}_6(\mathbb{C})$ is isomorphic to \mathbb{C}^4 as a $\mathfrak{sl}_4(\mathbb{C})$ -module and $\odot^k \mathbb{C}^4$ is the higher spin representation of $\mathfrak{so}_6(\mathbb{C})$. So Kang and Wang [13] defined

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¹Department of Mathematics, Zhejiang International Studies University, Hangzhou 310023, China.

E-mail: gzren@zisu.edu.cn qqkang@zisu.edu.cn

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the elliptic version of the massless field operators on \mathbb{R}^6 by the imbedding

$$\tau : \mathbb{R}^6 \hookrightarrow \wedge^2 \mathbb{C}^4 \subseteq \mathbb{C}^{4 \times 4},$$

$$\mathbf{x} = (x^0, x^1, \dots, x^5) \mapsto \tau_0(\mathbf{x}) := \begin{pmatrix} 0 & \mathbf{i}x^1 + x^6 & x^4 + \mathbf{i}x^5 & x^2 + \mathbf{i}x^3 \\ -\mathbf{i}x^1 - x^6 & 0 & x^2 - \mathbf{i}x^3 & -x^4 + \mathbf{i}x^5 \\ -x^4 - \mathbf{i}x^5 & -x^2 + \mathbf{i}x^3 & 0 & -\mathbf{i}x^1 + x^6 \\ -x^2 - \mathbf{i}x^3 & x^4 - \mathbf{i}x^5 & \mathbf{i}x^1 - x^6 & 0 \end{pmatrix}. \quad (1.2)$$

This embedding is the generalization of the embedding of quaternionic space $\mathbb{H} \hookrightarrow 2 \times 2$, which is based on

$$\mathfrak{so}_4(\mathbb{C}) \cong \mathfrak{sl}_2(\mathbb{C}) \oplus \mathfrak{sl}_2(\mathbb{C}). \quad (1.3)$$

Sämann and Wolf [21] and Mason et al. [16] defined the massless field operators on Lorentzian space $\mathbb{R}^{5,1}$ by an embedding $\mathbb{R}^{5,1}$ into $\mathbb{C}^{4 \times 4}$, which is the generalization of the embedding of Minkowski space $\mathbb{R}^{3,1}$ into 2×2 -Hermitian matrix space. The massless field operators are higher spin operators.

By embedding \mathbb{H}^n into $\mathbb{C}^{2n \times 2}$, Wang defined the k -Cauchy-Fueter operator over \mathbb{H}^n and gave the k -Cauchy-Fueter complex and twistor transform over \mathbb{H}^n by using the twistor method with the advantage of two-component station, see [23]. Motivated by the embedding (1.2), we give an embedding $\tau : \mathbb{R}^{6n} \hookrightarrow \mathbb{C}^{4n \times 4}$:

$$(z^{\alpha B}) := \tau(\mathbf{x}) := \begin{pmatrix} \tau_0(\mathbf{x}_0) \\ \tau_0(\mathbf{x}_1) \\ \dots \\ \tau_0(\mathbf{x}_{n-1}) \end{pmatrix}, \quad (1.4)$$

$\alpha = 0, 1, \dots, 4n-1$, $B = 0, 1, 2, 3$, where we denote the element of \mathbb{R}^{6n} by $\mathbf{x} = (\mathbf{x}_0, \mathbf{x}_1, \dots, \mathbf{x}_{n-1})$ with $\mathbf{x}_l = (x_{6l+1}, x_{6l+2}, \dots, x_{6l+6})$. In this paper, we use the notation $A, B, C, D \in \{0, 1, 2, 3\}$, $\alpha, \beta \in \{0, 1, \dots, 4n-1\}$, $j, l, m \in \{0, 1, \dots, n-1\}$. We denote by $\odot^k \mathbb{C}^4$ the k -th symmetric power of \mathbb{C}^4 , whose element is 4^k -tuple $(\phi_{A_1 \dots A_k})$ with $\phi_{A_1 \dots A_k} \in \mathbb{C}$, which are invariant under the permutations of subscripts $A_1, \dots, A_k = 0, 1, 2, 3$.

k -Cauchy-Fueter operator is valued in the higher spin representations of $\mathfrak{so}_4(\mathbb{C})$. Recently, there are many works to construct higher spin operators on higher dimensional Euclidean space and to investigate its function theory (see [8–10] and references there in).

Motivated by the higher spin operators on \mathbb{R}^6 (see [13]), we introduce higher spin operators \mathcal{D}_k of several \mathbb{R}^6 variables as

$$\mathcal{D}_k : C^\infty(\mathbb{R}^{6n}, \odot^k \mathbb{C}^4) \rightarrow C^\infty(\mathbb{R}^{6n}, \mathbb{C}^4 \otimes \odot^{k-1} \mathbb{C}^4), \quad (1.5)$$

with

$$(\mathcal{D}_k \phi)_{A_2 \dots A_k}^\alpha := \sum_{B=0}^3 Z^{\alpha B} \phi_{A_2 \dots A_k B} \quad (1.6)$$

for $k = 1, 2, \dots$, where $Z^{\alpha B}$ are complex vector fields and the matrix $(Z^{\alpha B})$ is just the embedding matrix (1.2)–(1.4) with the coordinate x^j replaced by ∂_{x^j} , see (2.2) for details. A $\odot^k \mathbb{C}^4$ -valued function ϕ on a domain U of \mathbb{R}^{6n} is called k -monogenic on U if

$$\mathcal{D}_k \phi = 0.$$

In this way, we could begin to study regular functions of several \mathbb{R}^6 variables.

Integral representation formula, i.e., Bochner-Martinelli integral formula is a fundamental and powerful tool to study holomorphic functions of several complex variables, see [18–19, 24–26] and references there in. For quaternion Siegel upper half-space, we know a better kernel, Cauchy-Szegö kernel reproduces regular functions, see [3–4]. In [23], Wang gave the Bochner-Martinelli type formula for k -Cauchy-Fueter operator. Wang and Ren [22] established the explicit Bochner-Martinelli kernel by using technique of Clifford analysis. They also established Bochner-Martinelli type formula for several octonionic variables, see [19]. Kang and Wang [13] gave Bochner-Martinelli type formula for higher spin massless field operators on \mathbb{R}^6 . In this paper, we get the Bochner-Martinelli type formula for higher spin operator \mathcal{D}_k on \mathbb{R}^{6n} .

Theorem 1.1 (Bochner-Martinelli type formula) *Let U be an open bounded set of \mathbb{R}^{6n} with C^1 boundary. For $\phi \in C(\overline{U}, \odot^k \mathbb{C}^4) \cap C^1(U, \odot^k \mathbb{C}^4)$, $k = 1, 2, \dots$ and any fixed $\mathbf{x} \in U$, we have*

$$\begin{aligned} \phi_{A_1 \dots A_k}(\mathbf{x}) &= \sum_{\alpha} \sum_B \int_{\partial U} \mathcal{E}_{\alpha(A_1}(\mathbf{x} - \mathbf{y}) \phi_{A_2 \dots A_k)B}(\mathbf{y}) d\sigma^{\alpha B}(\mathbf{y}) \\ &\quad - \sum_{\alpha} \int_U \mathcal{E}_{\alpha(A_1}(\mathbf{x} - \mathbf{y}) (\mathcal{D}_k \phi(\mathbf{y}))^{\alpha}_{A_2 \dots A_k)B} dv(\mathbf{y}). \end{aligned} \quad (1.7)$$

Wang and Kang introduced the Penrose integral formula and series expansion of k -regular functions over \mathbb{H}^n (see [12]). Sämann and Wolf [21] constructed the Penrose integral formula over the twistor space \mathbb{CP}^3 . We construct lots of k -monogenic functions by writing down the Penrose integral formula explicitly in terms of nonhomogeneous coordinates.

Let (ω, p) be the coordinates of \mathbb{C}^{4n+3} , $\omega \in \mathbb{C}^3, p \in \mathbb{C}^{4n}$. The Penrose type contour integral transformation is

$$\begin{aligned} &(\mathcal{P}f)_{A_1 \dots A_k}(\mathbf{x}) \\ &:= \int_{|\omega_1|=1} d\omega_1 \int_{|\omega_2|=1} d\omega_2 \int_{|\omega_3|=1} d\omega_3 \cdot \omega_{A_1} \dots \omega_{A_k} f\left(\omega_2, \omega_3, \omega_4, \dots, \sum_B z^{\alpha B} \omega_B, \dots\right) \end{aligned} \quad (1.8)$$

for f holomorphic on $\mathbb{C}^{4n+3} \setminus \{\omega_1 \omega_2 \omega_3 = 0\}$. We prove that the Penrose-type integral formula produces k -monogenic functions and find many nontrivial k -monogenic polynomials by choosing suitable holomorphic functions in the formula. Abundance of k -monogenic functions show the theory of several \mathbb{R}^6 variables is interesting.

Based on isomorphism (1.1) and (1.3), we can use two-component notation, which makes the spinor calculus on \mathbb{R}^{4n} and \mathbb{R}^{6n} more concrete and explicit. There is no such isomorphism for $\mathfrak{so}_n(\mathbb{C})$ with $n > 6$. That is why we only consider several \mathbb{R}^6 variables.

2 The Higher Spin Massless Field Operator \mathcal{D}_k Over \mathbb{R}^{6n}

In \mathbb{R}^6 (see [13]), we use complex vector fields

$$(\nabla^{AB}) := \begin{pmatrix} 0 & \mathbf{i}\partial_{x^0} + \partial_{x^5} & \partial_{x^3} + \mathbf{i}\partial_{x^4} & \partial_{x^1} + \mathbf{i}\partial_{x^2} \\ -\mathbf{i}\partial_{x^0} - \partial_{x^5} & 0 & \partial_{x^1} - \mathbf{i}\partial_{x^2} & -\partial_{x^3} + \mathbf{i}\partial_{x^4} \\ -\partial_{x^3} - \mathbf{i}\partial_{x^4} & -\partial_{x^1} + \mathbf{i}\partial_{x^2} & 0 & -\mathbf{i}\partial_{x^0} + \partial_{x^5} \\ -\partial_{x^1} - \mathbf{i}\partial_{x^2} & \partial_{x^3} - \mathbf{i}\partial_{x^4} & \mathbf{i}\partial_{x^0} - \partial_{x^5} & 0 \end{pmatrix}, \quad (2.1)$$

to define higher spin operators. Here we consider complex vector fields

$$(Z^{\alpha B})_{4n \times 4} := \begin{pmatrix} \vdots \\ Z^{(4l+A)B} \\ \vdots \end{pmatrix}, \quad (2.2)$$

where $l = 0, 1, \dots, n-1$, $A, B = 0, 1, 2, 3$ and

$$(Z^{(4l+A)B})_{4 \times 4} := \begin{pmatrix} 0 & \mathbf{i}\partial_{x_{6l+1}} + \partial_{x_{6l+6}} & \partial_{x_{6l+4}} + \mathbf{i}\partial_{x_{6l+5}} & \partial_{x_{6l+2}} + \mathbf{i}\partial_{x_{6l+3}} \\ -\mathbf{i}\partial_{x_{6l+1}} - \partial_{x_{6l+6}} & 0 & \partial_{x_{6l+2}} - \mathbf{i}\partial_{x_{6l+3}} & -\partial_{x_{6l+4}} + \mathbf{i}\partial_{x_{6l+5}} \\ -\partial_{x_{6l+4}} - \mathbf{i}\partial_{x_{6l+5}} & -\partial_{x_{6l+2}} + \mathbf{i}\partial_{x_{6l+3}} & 0 & -\mathbf{i}\partial_{x_{6l+1}} + \partial_{x_{6l+6}} \\ -\partial_{x_{6l+2}} - \mathbf{i}\partial_{x_{6l+3}} & \partial_{x_{6l+4}} - \mathbf{i}\partial_{x_{6l+5}} & \mathbf{i}\partial_{x_{6l+1}} - \partial_{x_{6l+6}} & 0 \end{pmatrix}. \quad (2.3)$$

Let $\varepsilon_{ABCD} = \varepsilon^{ABCD}$ be the sign of the permutation from $(0, 1, 2, 3)$ to (A, B, C, D) . Obviously, we have $\sum_{C,D} \varepsilon_{ABCD} \varepsilon^{CDEF} = 2(\delta_A^E \delta_B^F - \delta_A^F \delta_B^E)$. We will use $\varepsilon_{ABCD} = \varepsilon^{ABCD}$ to lower and raise indices:

$$Z_{\alpha B} := \frac{1}{2} \sum_{C,D} \varepsilon_{ABCD} Z^{(4l+C)D}, \quad Z^{\alpha B} := \frac{1}{2} \sum_{C,D} \varepsilon^{ABCD} Z_{(4l+C)D}, \quad (2.4)$$

when $\alpha = 4l + A$. By (2.2) and (2.4), we have

$$(Z_{\alpha B}) = (\overline{Z^{\alpha B}}). \quad (2.5)$$

The following three propositions generalize the corresponding results in 6 dimensional case in [13].

Proposition 2.1 For $z^{\alpha B}$ defined by (1.4) and the operator $Z_{\beta D}$ defined by (2.5), we have

$$Z_{(4m+C)D} z^{(4l+A)B} = 2\delta_m^l (\delta_C^A \delta_D^B - \delta_C^B \delta_D^A). \quad (2.6)$$

Proof When $m \neq l$, $Z_{(4m+C)D} z^{(4l+A)B} = 0$ for any $A, B, C, D \in \{0, 1, 2, 3\}$, by using (1.2)–(1.4) and (2.3)–(2.5). Now we consider the case $m = l$. It is direct to check that

$$\begin{aligned} Z_{(4l)1} z^{(4l)1} &= (-\mathbf{i}\partial_{x_{6l+1}} + \partial_{x_{6l+6}})(\mathbf{i}x^{6l+1} + x^{6l+6}) = 2, \\ Z_{(4l)1} z^{(4l+2)3} &= (-\mathbf{i}\partial_{x_{6l+1}} + \partial_{x_{6l+6}})(-\mathbf{i}x^{6l+1} + x^{6l+6}) = 0. \end{aligned}$$

So we have $Z_{(4l)1} z^{(4l+1)0} = -2$ and $Z_{(4l)1} z^{(4l+3)2} = 0$ by $z^{(4l+A)B} = -z^{(4l+B)A}$. And we have $Z_{(4l)1} z^{(4l+A)B} = 0$ for other case, since z^{AB} is independent of x^{6l+1} and x^{6l+6} . Hence we have

$$Z_{(4l)1} z^{(4l+A)B} = 2(\delta_0^A \delta_1^B - \delta_0^B \delta_1^A),$$

i.e., (2.6) holds for $(C, D) = (0, 1)$. Similarly, when $(C, D) \neq (0, 1)$, (2.6) also holds. The proposition is proved.

Proposition 2.2 For operator $Z^{\alpha B}$ and $Z_{\alpha B}$ defined by (2.2)–(2.3) and (2.5), we have

$$\sum_{\alpha} Z_{\alpha B_1} Z^{\alpha B_2} = \delta_{B_1}^{B_2} \Delta, \quad (2.7)$$

where $\Delta = \sum_{i=1}^{6n} \partial_{x_i}^2$ is the Laplace operator on \mathbb{R}^{6n} .

Proof When $B_1 \neq B_2$, we have

$$\sum_{\alpha} Z_{\alpha B_1} Z^{\alpha B_2} = \frac{1}{2} \sum_l \sum_A \sum_{C,D} \varepsilon_{AB_1 CD} Z^{(4l+C)D} Z^{(4l+A)B_2}$$

by definition (2.4). Note that $Z^{(4l+A)B_2} = 0$ if $A = B_2$. Thus $\varepsilon_{AB_1 CD} \nabla^{(4l+A)B_2} \neq 0$ if and only if C or $D = B_2$. So we have

$$\begin{aligned} \sum_{\alpha} Z_{\alpha B_1} Z^{\alpha B_2} &= \frac{1}{2} \sum_l \sum_A \left(\sum_C \varepsilon_{AB_1 C B_2} Z^{(4l+C)B_2} + \sum_D \varepsilon_{AB_1 B_2 D} Z^{(4l+B_2)D} \right) Z^{(4l+A)B_2} \\ &= \frac{1}{2} \sum_l \sum_A \left(\sum_C \varepsilon_{AB_1 C B_2} Z^{(4l+C)B_2} + \sum_D \varepsilon_{AB_1 D B_2} Z^{(4l+D)B_2} \right) Z^{(4l+A)B_2} \\ &= \sum_l \sum_A \sum_C \varepsilon_{AB_1 C B_2} Z^{(4l+C)B_2} Z^{(4l+A)B_2} = 0 \end{aligned}$$

by $\varepsilon_{AB_1 B_2 D} = -\varepsilon_{AB_1 D B_2}$ and $Z^{B_2 D} = -Z^{D B_2}$. When $B_1 = B_2$, we have

$$\sum_{\alpha} Z_{\alpha B} Z^{\alpha B} = \sum_{\alpha} \overline{Z^{\alpha B}} Z^{\alpha B}$$

by using (2.5). It is directed to check that

$$\sum_A \overline{Z^{(4l+A)B}} Z^{(4l+A)B} = \partial_{x^{6l+1}}^2 + \partial_{x^{6l+6}}^2 + \partial_{x^{6l+4}}^2 + \partial_{x^{6l+5}}^2 + \partial_{x^{6l+2}}^2 + \partial_{x^{6l+3}}^2$$

for fixed $l \in \{0, 1, \dots, n-1\}$ and $B \in \{0, 1, 2, 3\}$. So we have

$$\sum_{\alpha} \overline{Z^{\alpha B}} Z^{\alpha B} = \Delta.$$

The proposition is proved.

Define the inner products

$$(\phi, \psi) := \sum_{A_1, A_2, \dots, A_k} \int_{\mathbb{R}^{6n}} \phi_{A_1 A_2 \dots A_k}(\mathbf{y}) \overline{\psi_{A_1 A_2 \dots A_k}(\mathbf{y})} dv(\mathbf{y}), \quad \phi, \psi \in C_0^\infty(\mathbb{R}^{6n}, \mathbb{C}^{4n} \otimes \odot^{k-1} \mathbb{C}^4)$$

and

$$(\psi, \Psi) := \sum_{\alpha} \sum_{A_2, \dots, A_k} \int_{\mathbb{R}^{6n}} \psi_{A_2 \dots A_k}^{\alpha}(\mathbf{y}) \overline{\Psi_{A_2 \dots A_k}^{\alpha}(\mathbf{y})} dv(\mathbf{y}), \quad \psi, \Psi \in C_0^\infty(\mathbb{R}^{6n}, \mathbb{C}^{4n} \otimes \odot^{k-1} \mathbb{C}^4).$$

Let $L^2(\mathbb{R}^{6n}, \odot^k \mathbb{C}^4)$ and $L^2(\mathbb{R}^{6n}, \mathbb{C}^{4n} \otimes \odot^{k-1} \mathbb{C}^4)$ be the completion of $C_0^\infty(\mathbb{R}^{6n}, \odot^k \mathbb{C}^4)$ and $C_0^\infty(\mathbb{R}^{6n}, \mathbb{C}^{4n} \otimes \odot^{k-1} \mathbb{C}^4)$, respectively. Let \mathcal{D}_k^* be the adjoint of the operator \mathcal{D}_k in (1.5) as

$$(\mathcal{D}_k \phi, \psi) = (\phi, \mathcal{D}_k^* \psi)$$

for any $\phi \in C_0^\infty(\mathbb{R}^{6n}, \odot^k \mathbb{C}^4)$, $\psi \in C_0^\infty(\mathbb{R}^{6n}, (\mathbb{C}^4)^* \otimes \odot^{k-1} \mathbb{C}^4)$.

We will use primed symmetrization of indices

$$\phi_{\dots(A_1 \dots A_k)\dots} := \frac{1}{k!} \sum_{\sigma \in S_k} \phi_{\dots A_{\sigma(1)} \dots A_{\sigma(k)} \dots}, \quad (2.8)$$

where S_k is the group of permutations of k letters.

For any $\phi, \psi \in L^2(\mathbb{R}^{6n}, \odot^k \mathbb{C}^4)$, by their symmetrization, we have

$$\sum_{A_1, \dots, A_k} \phi_{A_1 A_2 \dots A_k} \overline{\psi_{A_1 A_2 \dots A_k}} = \sum_{A_1, \dots, A_k} \phi_{A_1 A_2 \dots A_k} \overline{\psi_{(A_1 A_2 \dots A_k)}}, \quad \text{a.e..} \quad (2.9)$$

Proposition 2.3 For any $\phi \in C_0^\infty(\mathbb{R}^{6n}, \odot^k \mathbb{C}^4)$, $k = 1, 2, \dots$, we have

$$\mathcal{D}_k^* \mathcal{D}_k \phi = -\Delta \phi,$$

where Δ is the Laplace operator on \mathbb{R}^{6n} .

Proof For $\phi \in C_0^\infty(\mathbb{R}^{6n}, \odot^k \mathbb{C}^4)$ and $\psi \in C_0^\infty(\mathbb{R}^{6n}, \mathbb{C}^{4n} \otimes \odot^{k-1} \mathbb{C}^4)$, we have

$$\begin{aligned} (\mathcal{D}_k \phi, \psi) &= \sum_{\alpha} \sum_{A_2, \dots, A_k} \int_{\mathbb{R}^{6n}} \sum_{A_1} Z^{\alpha A_1} \phi_{A_2 \dots A_k A_1} \cdot \overline{\psi_{A_2 \dots A_k}^{\alpha}} dv(\mathbf{y}) \\ &= - \sum_{A_1, \dots, A_k} \int_{\mathbb{R}^{6n}} \phi_{A_1 A_2 \dots A_k} \cdot \overline{\sum_{\alpha} Z_{\alpha A_1} \psi_{A_2 \dots A_k}^{\alpha}} dv(\mathbf{y}) \\ &= - \sum_{A_1, \dots, A_k} \int_{\mathbb{R}^{6n}} \phi_{A_1 A_2 \dots A_k} \cdot \overline{\sum_{\alpha} Z_{\alpha(A_1} \psi_{A_2 \dots A_k)}^{\alpha}} dv(\mathbf{y}) = (\phi, \mathcal{D}_k^* \psi), \end{aligned}$$

which follows from Stokes' formula. Here we use (2.9) in the third identity. Note that

$$(\mathcal{D}_k^* \psi)_{A_1 \dots A_k} = - \sum_{\alpha} Z_{\alpha(A_1} \psi_{A_2 \dots A_k)}^{\alpha}.$$

We get

$$\begin{aligned} (\mathcal{D}_k^* \mathcal{D}_k \phi)_{A_1 \dots A_k} &= -\frac{1}{k!} \sum_{\alpha} \sum_{\sigma \in S_k} Z_{\alpha A_{\sigma(1)}} \left(\sum_B Z^{\alpha B} \phi_{\alpha_{\sigma(2)} \dots A_{\sigma(k)} B} \right) \\ &= -\frac{1}{k!} \sum_{\sigma \in S_k} \sum_B \delta_{A_{\sigma(1)}}^B \Delta \phi_{A_{\sigma(2)} \dots A_{\sigma(k)} B} = -\frac{1}{k!} \sum_{\sigma \in S_k} \Delta \phi_{A_{\sigma(2)} \dots A_{\sigma(k)} A_{\sigma(1)}} \\ &= -\Delta \phi_{A_1 A_2 \dots A_k} \end{aligned}$$

by using (1.6) and (2.7). Then $(\mathcal{D}_k^* \mathcal{D}_k) \phi = -\Delta \phi$. The proposition is proved.

3 The Bochner-Martinelli Integral Formula

It is directed to check that $K(\mathbf{y}) := \frac{C}{|\mathbf{y}|^{6n-2}}$ is the fundamental solution of the Laplacian operator Δ on \mathbb{R}^{6n} , where $C = \frac{(3n)!}{12n(1-3n)\pi^{3n}}$ and $|\mathbf{y}| = (y_1^2 + y_2^2 + \dots + y_{6n}^2)^{\frac{1}{2}}$. Define the Bochner-Martinelli kernel

$$\mathcal{E}_{\alpha B}(\mathbf{y}) := \overline{Z^{\alpha B}} K(\mathbf{y}) = Z_{\alpha B} K(\mathbf{y}),$$

which is a $\mathbb{C}^{4n} \otimes \odot^k \mathbb{C}^4$ valued function. Define the $(6n-1)$ -form

$$d\widehat{y}^{\alpha} := i_{\partial_{y^{\alpha}}} dv(\mathbf{y}), \quad d\sigma^{\alpha B}(\mathbf{y}) := i_{Z^{\alpha B}} dv(\mathbf{y}), \quad (3.1)$$

where i is the interior product and $dv(\mathbf{y}) := dy^1 \wedge dy^2 \wedge \dots \wedge dy^{6n}$ is the standard volume form on \mathbb{R}^{6n} . Obviously, we have

$$i_{(c_1 \partial_{y^{\beta_1}} + c_2 \partial_{y^{\beta_2}})} dv(\mathbf{y}) = c_1 d\widehat{y}^{\beta_1} + c_2 d\widehat{y}^{\beta_2}, \quad \beta_1, \beta_2 \in \{1, \dots, 6n\}.$$

Then $d\sigma^{\alpha B}(\mathbf{y})$ in (3.1) can be rewritten as

$$(d\sigma^{\alpha B}(\mathbf{y})) = \begin{pmatrix} \vdots \\ d\sigma^{(4l+A)B}(\mathbf{y}_l) \\ \vdots \end{pmatrix},$$

where $\mathbf{y}_l \in \mathbb{R}^6$, $l = 0, 1, \dots, n-1$, $A, B = 0, 1, 2, 3$ and

$$(d\sigma^{(4l+A)B}(\mathbf{y}_l)) = \begin{pmatrix} 0 & \widehat{\text{id}y^{6l+1} + dy^{6l+6}} & \widehat{dy^{6l+4} + \text{id}y^{6l+5}} & \widehat{dy^{6l+2} + \text{id}y^{6l+3}} \\ -\widehat{\text{id}y^{6l+1} - dy^{6l+6}} & 0 & \widehat{dy^{6l+2} - \text{id}y^{6l+3}} & -\widehat{dy^{6l+4} + \text{id}y^{6l+5}} \\ -\widehat{dy^{6l+4} - \text{id}y^{6l+5}} & -\widehat{dy^{6l+2} + \text{id}y^{6l+3}} & 0 & -\widehat{\text{id}y^{6l+1} + dy^{6l+6}} \\ -\widehat{dy^{6l+2} - \text{id}y^{6l+3}} & \widehat{dy^{6l+4} - \text{id}y^{6l+5}} & \widehat{\text{id}y^{6l+1} - dy^{6l+6}} & 0 \end{pmatrix}.$$

This matrix is just (1.2) with x^i replaced by $\widehat{dy^{6l+i}}$.

Lemma 3.1 For a domain U in \mathbb{R}^{6n} and $u, w \in C^1(U, \mathbb{C})$, we have

$$d(u \cdot w d\sigma^{\alpha B}(\mathbf{y})) = (Z^{\alpha B} u \cdot w + u \cdot Z^{\alpha B} w) dv(\mathbf{y}). \quad (3.2)$$

Proof For fixed $j_1, j_2 \in \{1, \dots, 6n\}$, we have

$$\sum_{i=1}^{6n} \partial_{y^i} h \cdot dy^i \wedge (i_{c_1 \partial_{y^{j_1}} + c_2 \partial_{y^{j_2}}} dv(\mathbf{y})) = (c_1 \partial_{y^{j_1}} + c_2 \partial_{y^{j_2}}) h \cdot dv(\mathbf{y}). \quad (3.3)$$

Thus for $h(\mathbf{y}) \in C^1(U)$, we have

$$d(h \cdot d\sigma^{\alpha B}) = \sum_{i=1}^{6n} \partial_{y^i} h \cdot dy^i \wedge d\sigma^{\alpha B} = \sum_{i=1}^{6n} \partial_{y^i} h \cdot dy^i \wedge i_{Z^{\alpha B}} dv(\mathbf{y}) = Z^{\alpha B} h \cdot dv(\mathbf{y})$$

by (3.3). Then (3.2) follows.

Lemma 3.2 Let U be an bounded domain of \mathbb{R}^{6n} with C^1 boundary ∂U and $\Phi \in C(\overline{U}, \odot^k \mathbb{C}^4) \cap C^1(U, \odot^k \mathbb{C}^4)$. For a fixed point $\mathbf{x} \notin U$ and fixed $A_1, \dots, A_k \in \{0, 1, 2, 3\}$, we have

$$\begin{aligned} & \int_{\partial U} \sum_{\alpha} \sum_B \mathcal{E}_{\alpha A_1}(\mathbf{x} - \mathbf{y}) \Phi_{A_2 \dots A_k B}(\mathbf{y}) d\sigma^{\alpha B}(\mathbf{y}) \\ &= \sum_{\alpha} \int_U \mathcal{E}_{\alpha A_1}(\mathbf{x} - \mathbf{y}) (\mathcal{D}_k \Phi(\mathbf{y}))_{A_2 \dots A_k}^{\alpha} dv(\mathbf{y}). \end{aligned} \quad (3.4)$$

Proof Apply Lemma 3.1 to $u(\cdot) = Z_{AA_1} K(\mathbf{x} - \cdot)$, $w(\cdot) = \Phi_{A_2 \dots A_k B}(\cdot)$ to get

$$\begin{aligned} & \int_{\partial U} dy \left[\sum_{\alpha} \sum_B \mathcal{E}_{\alpha A_1}(\mathbf{x} - \mathbf{y}) \Phi_{A_2 \dots A_k B}(\mathbf{y}) d\sigma^{\alpha B}(\mathbf{y}) \right] \\ &= \sum_{\alpha} \sum_B \int_U [-Z^{\alpha B} Z_{\alpha A_1} K(\mathbf{x} - \mathbf{y}) \cdot \Phi_{A_2 \dots A_k B}(\mathbf{y}) + Z_{\alpha A_1} K(\mathbf{x} - \mathbf{y}) \cdot Z^{\alpha B} \Phi_{A_2 \dots A_k B}(\mathbf{y})] dv(\mathbf{y}) \\ &= \sum_B \int_U -\delta_{A_1}^B \Delta K(\mathbf{x} - \mathbf{y}) \cdot \Phi_{A_2 \dots A_k B}(\mathbf{y}) dv(\mathbf{y}) + \sum_{\alpha} \int_U \mathcal{E}_{\alpha A_1}(\mathbf{x} - \mathbf{y}) (\mathcal{D}_k \Phi(\mathbf{y}))_{A_2 \dots A_k}^{\alpha} dv(\mathbf{y}) \\ &= \sum_{\alpha} \int_U \mathcal{E}_{\alpha A_1}(\mathbf{x} - \mathbf{y}) (\mathcal{D}_k \Phi(\mathbf{y}))_{A_2 \dots A_k}^{\alpha} dv(\mathbf{y}) \end{aligned}$$

by $\sum_{\alpha} Z^{\alpha B} Z_{\alpha A_1} = \sum_{\alpha} Z_{\alpha A_1} Z^{\alpha B} = \delta_{A_1}^B \Delta$ and $\Delta K(\mathbf{x} - \mathbf{y}) = 0$ for $\mathbf{x} \notin U$.

We also need the identity corresponding to (3.4) for $x \in U$.

Lemma 3.3 *Let U be an bounded domain in \mathbb{R}^{6n} with C^1 boundary ∂U and $\phi \in C(\overline{U}, \odot^k \mathbb{C}^4) \cap C^1(U, \odot^k \mathbb{C}^4)$. Then we have*

$$\int_{\partial U} \sum_{\alpha} \sum_B \mathcal{E}_{\alpha(A_1}(\mathbf{x} - \mathbf{y}) \phi_{A_2 \dots A_k B}(\mathbf{x}) d\sigma^{\alpha B}(\mathbf{y}) = \begin{cases} \phi_{A_1 \dots A_k}(\mathbf{x}), & \text{if } \mathbf{x} \in U, \\ 0, & \text{if } \mathbf{x} \notin \overline{U}. \end{cases} \quad (3.5)$$

Proof For a fixed point $\mathbf{x} \notin \overline{U}$, applying Lemma 3.2 to $\Phi(\mathbf{y}) \equiv \phi(\mathbf{x})$ to get

$$\int_{\partial U} \sum_{\alpha} \sum_B \mathcal{E}_{\alpha A_1}(\mathbf{x} - \mathbf{y}) \phi_{A_2 \dots A_k B}(\mathbf{x}) d\sigma^{\alpha B}(\mathbf{y}) = 0$$

by $\mathcal{D}_k \Phi(\mathbf{y}) \equiv 0$. It implies (3.5) holds for $\mathbf{x} \notin \overline{U}$ by symmetrization.

When $\mathbf{x} \in U$, let $B(\mathbf{x}, \varepsilon) = \{\mathbf{y} \in \mathbb{R}^{6n} : |\mathbf{y} - \mathbf{x}| \leq \varepsilon\} \subseteq U$ with ε sufficiently small. Applying Lemma 3.2 to $\Phi(\mathbf{y}) \equiv \phi(\mathbf{x})$ and the domain $U \setminus B(\mathbf{x}, \varepsilon)$, we have

$$\int_{\partial U \setminus \partial B(\mathbf{x}, \varepsilon)} \sum_{\alpha} \sum_B \mathcal{E}_{\alpha A_1}(\mathbf{x} - \mathbf{y}) \phi_{A_2 \dots A_k B}(\mathbf{x}) d\sigma^{\alpha B}(\mathbf{y}) = 0,$$

i.e.,

$$\begin{aligned} & \int_{\partial U} \sum_{\alpha} \sum_B \mathcal{E}_{\alpha A_1}(\mathbf{x} - \mathbf{y}) \phi_{A_2 \dots A_k B}(\mathbf{x}) d\sigma^{\alpha B}(\mathbf{y}) \\ &= \int_{\partial B(\mathbf{x}, \varepsilon)} \sum_{\alpha} \sum_B \mathcal{E}_{\alpha A_1}(\mathbf{x} - \mathbf{y}) \phi_{A_2 \dots A_k B}(\mathbf{x}) d\sigma^{\alpha B}(\mathbf{y}). \end{aligned} \quad (3.6)$$

Obviously, we have $Z_{\alpha B} |\mathbf{x} - \mathbf{y}|^2 = -2\overline{\iota(\mathbf{x} - \mathbf{y})}_{\alpha B}$. Then the right hand side of (3.6) equals to

$$\begin{aligned} & \sum_{\alpha} \sum_B \phi_{A_2 \dots A_k B}(\mathbf{x}) \int_{\partial B(\mathbf{x}, \varepsilon)} \frac{(-3n+1)C}{\varepsilon^{6n}} Z_{\alpha A_1} |\mathbf{x} - \mathbf{y}|^2 d\sigma^{\alpha B}(\mathbf{y}) \\ &= \sum_{\alpha} \sum_B \phi_{A_2 \dots A_k B}(\mathbf{x}) \int_{\partial B(\mathbf{x}, \varepsilon)} \frac{(3n)!}{12n\pi^{3n}\varepsilon^{6n}} Z_{\alpha A_1} |\mathbf{x} - \mathbf{y}|^2 d\sigma^{\alpha B}(\mathbf{y}) \\ &= \sum_{\alpha} \sum_B \phi_{A_2 \dots A_k B}(\mathbf{x}) \int_{B(\mathbf{x}, \varepsilon)} \frac{(3n)!}{12n\pi^{3n}\varepsilon^{6n}} Z^{\alpha B} Z_{\alpha A_1} |\mathbf{x} - \mathbf{y}|^2 dv(\mathbf{y}) \\ &= \sum_B \phi_{A_2 \dots A_k B}(\mathbf{x}) \int_{B(\mathbf{x}, \varepsilon)} \frac{(3n)!}{12n\pi^{3n}\varepsilon^{6n}} \delta_{A_1}^B \Delta |\mathbf{x} - \mathbf{y}|^2 dv(\mathbf{y}) \\ &= \frac{(3n)!}{\pi^{3n}\varepsilon^{6n}} |B(\mathbf{x}, \varepsilon)| \phi_{A_2 \dots A_k A_1}(\mathbf{x}) = \phi_{A_1 \dots A_k}(\mathbf{x}) \end{aligned} \quad (3.7)$$

by using Stokes' formula and (3.2) in the first identity, where $|B(\mathbf{x}, \varepsilon)| = \frac{\pi^{3n}\varepsilon^{6n}}{(3n)!}$ is the volume of the ball $B(\mathbf{x}, \varepsilon)$ in \mathbb{R}^{6n} . The lemma is proved.

Proof of Theorem 1.1 For any fixed $x \in U$, applying Lemma 3.2 to $\Phi(\mathbf{y}) = \phi(\mathbf{y})$ and the domain $U \setminus B(\mathbf{x}, \varepsilon)$ to get

$$\begin{aligned} & \left\{ \int_{\partial U} - \int_{\partial B(\mathbf{x}, \varepsilon)} \right\} \sum_{\alpha} \sum_B \mathcal{E}_{\alpha A_1}(\mathbf{x} - \mathbf{y}) \phi_{A_2 \dots A_k B}(\mathbf{y}) d\sigma^{\alpha B} \\ &= \sum_{\alpha} \int_{U \setminus B(\mathbf{x}, \varepsilon)} \mathcal{E}_{\alpha A_1}(\mathbf{x} - \mathbf{y}) (\mathcal{D}_k \phi(\mathbf{y}))_{A_2 \dots A_k}^{\alpha} dv. \end{aligned}$$

By symmetrizing indices A_1, \dots, A_k , we have

$$\begin{aligned} & \left\{ \int_{\partial U} - \int_{\partial B(\mathbf{x}, \varepsilon)} \right\} \sum_{\alpha} \sum_B \mathcal{E}_{\alpha(A_1}(\mathbf{x} - \mathbf{y}) \phi_{A_2 \dots A_k)B}(\mathbf{y}) d\sigma^{\alpha B} \\ &= \sum_{\alpha} \int_{U \setminus B(\mathbf{x}, \varepsilon)} \mathcal{E}_{\alpha(A_1}(\mathbf{x} - \mathbf{y}) (\mathcal{D}_k \phi(\mathbf{y}))_{A_2 \dots A_k}^{\alpha} dv. \end{aligned} \quad (3.8)$$

Since $Z^{AB} \phi(\mathbf{y})_{A_2 \dots A_k B}$ is locally bounded and $|\mathcal{E}_{AB}(\mathbf{x} - \mathbf{y})| = O(|\mathbf{x} - \mathbf{y}|^{-6n+1})$, we find that

$$\lim_{\varepsilon \rightarrow 0} \sum_{\alpha} \int_{B(\mathbf{x}, \varepsilon)} \mathcal{E}_{\alpha A_1}(\mathbf{x} - \mathbf{y}) (\mathcal{D}_k \phi(\mathbf{y}))_{A_2 \dots A_k}^{\alpha} dv(\mathbf{y}) = 0.$$

On the other hand,

$$\lim_{\varepsilon \rightarrow 0} \sum_{\alpha} \sum_B \int_{\partial B(\mathbf{x}, \varepsilon)} \mathcal{E}_{\alpha A_1}(\mathbf{x} - \mathbf{y}) [\phi_{A_2 \dots A_k B}(\mathbf{y}) - \phi_{A_2 \dots A_k B}(\mathbf{x})] d\sigma^{\alpha B}(\mathbf{y}) = 0,$$

since $|f(\mathbf{y}) - f(\mathbf{x})| \leq \|f\|_{C^1} |\mathbf{x} - \mathbf{y}|$ for any $\mathbf{y} \in \partial B(\mathbf{x}, \varepsilon)$ and $f \in C^1(\overline{B(\mathbf{x}, \varepsilon)})$.

This together with Lemma 3.3 and symmetrizing the indices A_1, \dots, A_k imply that

$$\lim_{\varepsilon \rightarrow 0} \int_{\partial B(\mathbf{x}, \varepsilon)} \sum_{\alpha} \sum_B \mathcal{E}_{\alpha(A_1}(\mathbf{x} - \mathbf{y}) \phi_{A_2 \dots A_k)B}(\mathbf{y}) d\sigma^{\alpha B}(\mathbf{y}) = \phi_{A_1 \dots A_k}(\mathbf{x}).$$

The theorem is proved by letting $\varepsilon \rightarrow 0$ in (3.8).

4 The k -Monogenic Functions and Penrose Integral Formula

Sämman and Wolf [21] have established the Penrose integral formula in terms of homogeneous coordinates of \mathbb{CP}^7 . We can get the contour integral formula (1.8) over \mathbb{R}^{6n} by taking nonhomogeneous coordinates of \mathbb{CP}^3 . We check that $\mathcal{P}f$ satisfies the equation $\mathcal{D}_k(\mathcal{P}f) = 0$ by direct differentiation.

Theorem 4.1 *For any holomorphic function f on $\mathbb{C}^{4n+3} \setminus \{\omega_1 \omega_2 \omega_3 = 0\}$, the $\odot^k \mathbb{C}^4$ -valued function $\mathcal{P}f$ on \mathbb{R}^{6n} given by (1.8) is k -monogenic.*

Proof Take $p^{\alpha} := \sum_B z^{\alpha B} \omega_B$ with $\omega_0 = 1$. When $\alpha = 4l + A$,

$$\begin{aligned} Z_{(4m+C)D} p^{\alpha} &= \sum_B Z_{(4m+C)D} z^{(4l+A)B} \omega_B \\ &= 2 \sum_B \delta_m^l (\delta_C^A \delta_D^B - \delta_C^B \delta_D^A) \omega_B = 2 \delta_m^l (\delta_C^A \omega_D - \delta_D^A \omega_C) \end{aligned} \quad (4.1)$$

by (2.6). For a fixed A_2, \dots, A_k, B , differentiate the contour integral formula (1.8) to get

$$\begin{aligned} & Z_{(4m+C)D} (\mathcal{P}f)_{A_2 \dots A_k B}(\mathbf{x}) \\ &= \int_{|\omega_1|=1} d\omega_1 \int_{|\omega_2|=1} d\omega_2 \int_{|\omega_3|=1} d\omega_3 \cdot \omega_{A_2} \dots \omega_{A_k} \omega_B \sum_{\alpha} \frac{\partial f}{\partial p^{\alpha}} Z_{(4m+C)D} p^{\alpha} \\ &= 2 \int_{|\omega_1|=1} d\omega_1 \int_{|\omega_2|=1} d\omega_2 \int_{|\omega_3|=1} d\omega_3 \cdot \omega_{A_2} \dots \omega_{A_k} \omega_B \left(\frac{\partial f}{\partial p^{4m+C}} \omega_D - \frac{\partial f}{\partial p^{4m+D}} \omega_C \right) \end{aligned} \quad (4.2)$$

by (4.1). For fixed A_2, \dots, A_k , set

$$H_{AB(4m+C)}(\mathbf{x}) := \int_{|\omega_1|=1} d\omega_1 \int_{|\omega_2|=1} d\omega_2 \int_{|\omega_3|=1} d\omega_3 \cdot \omega_{A_2} \cdots \omega_{A_k} \omega_A \omega_B \frac{\partial f}{\partial p^{4m+C}}.$$

Obviously, we have

$$H_{AB(4m+C)}(\mathbf{x}) = H_{BA(4m+C)}(\mathbf{x}).$$

We can rewrite (4.2) as

$$\nabla_{(4m+C)D}(\mathcal{P}f)_{A_2 \cdots A_k B}(\mathbf{x}) = 2(H_{BD(4m+C)}(\mathbf{x}) - H_{BC(4m+D)}(\mathbf{x})). \quad (4.3)$$

Then we get

$$\begin{aligned} (\mathcal{D}_k(\mathcal{P}f))_{A_2 \cdots A_k}^{4m+A} &= \sum_B Z^{(4m+A)B}(\mathcal{P}f)_{A_2 \cdots A_k B}(\mathbf{x}) = \sum_{B,C,D} \frac{1}{2} \varepsilon^{ABCD} Z_{(4m+C)D}(\mathcal{P}f)_{A_2 \cdots A_k B}(\mathbf{x}) \\ &= \sum_{B,C,D} \varepsilon^{ABCD} (H_{BD(4m+C)}(\mathbf{x}) - H_{BC(4m+D)}(\mathbf{x})) \end{aligned}$$

by raising indices and using (2.4), (4.3) in the second and third identity, respectively. Since ε^{ABCD} is antisymmetric in indices and $H_{AB(4m+C)}$ is symmetric in A and B , we have

$$\sum_{B,D} \varepsilon^{ABCD} H_{BD(4m+C)} = 0 \quad \text{and} \quad \sum_{B,C} \varepsilon^{ABCD} H_{BC(4m+D)} = 0.$$

So $(\mathcal{D}_k(\mathcal{P}f))_{A_2 \cdots A_k}^\alpha = 0$, i.e., $(\mathcal{P}f)(\mathbf{x})$ is k -monogenic on \mathbb{R}^{6n} .

Now we give some concrete k -monogenic functions.

Corollary 4.1 *For fixed nonnegative integers q_1, q_2, q_3 and $m_j, j = 0, 1, \dots, n-1$, the $\odot^k \mathbb{C}^4$ -valued polynomial ϕ with*

$$\begin{aligned} &\phi_{A_1 \cdots A_k}(\mathbf{x}) \\ &= \begin{cases} \sum_{\substack{\alpha_0 + \alpha_1 + \cdots + \alpha_{n-1} + \alpha_1 = q_1 \\ \beta_0 + \beta_1 + \cdots + \beta_{n-1} + \alpha_2 = q_2 \\ \gamma_0 + \gamma_1 + \cdots + \gamma_{n-1} + \alpha_3 = q_3 \\ \alpha_j + \beta_j + \gamma_j = m_j, 0 \leq \alpha_j, \beta_j, \gamma_j \leq m_j}} \prod_{j=0}^{n-1} \frac{m_j!}{\alpha_j! \beta_j! \gamma_j!} (z^{(4j)1})^{\alpha_j} (z^{(4j)2})^{\beta_j} (z^{(4j)3})^{\gamma_j}, & \text{when } |\mathbf{a}| = |\mathbf{q}| - |\mathbf{m}| \\ 0, & \text{else} \end{cases} \end{aligned} \quad (4.4)$$

is k -monogenic, where a_h is the number of h in $\{A_1, \dots, A_k\}$, $h = 1, 2, 3$, $|\mathbf{a}| = a_1 + a_2 + a_3$, $|\mathbf{q}| = q_1 + q_2 + q_3$, $|\mathbf{m}| = \sum_{j=0}^{n-1} m_j$.

Proof Choose

$$f(\omega_1, \omega_2, \omega_3, \dots, p^\alpha, \dots) := \frac{\prod_{j=0}^{n-1} (p^{4j})^{m_j}}{\omega_1^{q_1+1} \omega_2^{q_2+1} \omega_3^{q_3+1}} \quad (4.5)$$

holomorphic on $\mathbb{C}^{4n+3} \setminus \{\omega_1 \omega_2 \omega_3 = 0\}$ for fixed nonnegative integers q_1, q_2, q_3 and $m_j, j = 0, 1, \dots, n-1$. Note that

$$\int_{|\omega_h|=1} \omega_h^s d\omega_h = \begin{cases} 2\pi i, & \text{if } s = -1, \\ 0, & \text{otherwise} \end{cases}$$

for $h = 1, 2, 3$. Recall that $\omega_0 \equiv 1$. Substitute f in (4.5) into the Penrose type contour integral formula (1.8) to get

$$\begin{aligned}
 (\mathcal{P}f)_{A_1 \dots A_k}(\mathbf{x}) &= \int_{|\omega_1|=1} d\omega_1 \int_{|\omega_2|=1} d\omega_2 \int_{|\omega_3|=1} d\omega_3 \cdot \omega_{A_1} \cdots \omega_{A_k} \frac{\prod_{j=0}^{n-1} (p^{4j})^{m_j}}{\omega_1^{q_1+1} \omega_2^{q_2+1} \omega_3^{q_3+1}} \\
 &= \int_{|\omega_1|=1} d\omega_1 \int_{|\omega_2|=1} d\omega_2 \int_{|\omega_3|=1} d\omega_3 \cdot \omega_1^{a_1} \omega_2^{a_2} \omega_3^{a_3} \prod_{j=0}^{n-1} \left(\sum_{\substack{\alpha_j + \beta_j + \gamma_j = m_j \\ 0 \leq \alpha_j, \beta_j, \gamma_j \leq m_j}} \frac{m_j!}{\alpha_j! \beta_j! \gamma_j!} \right. \\
 &\quad \cdot \left. \frac{(z^{(4j)1} \omega_1)^{\alpha_j} (z^{(4j)2} \omega_2)^{\beta_j} (z^{(4j)3} \omega_3)^{\gamma_j}}{\omega_1^{q_1+1} \omega_2^{q_2+1} \omega_3^{q_3+1}} \right) \\
 &= (2\pi i)^3 \sum_{\substack{\alpha_0 + \alpha_1 + \dots + \alpha_{n-1} + a_1 = q_1 \\ \beta_0 + \beta_1 + \dots + \beta_{n-1} + a_2 = q_2 \\ \gamma_0 + \gamma_1 + \dots + \gamma_{n-1} + a_3 = q_3 \\ \alpha_j + \beta_j + \gamma_j = m_j, 0 \leq \alpha_j, \beta_j, \gamma_j \leq m_j}} \prod_{j=0}^{n-1} \frac{m_j!}{\alpha_j! \beta_j! \gamma_j!} (z^{(4j)1})^{\alpha_j} (z^{(4j)2})^{\beta_j} (z^{(4j)3})^{\gamma_j}
 \end{aligned}$$

only if $|\mathbf{a}| = |\mathbf{q}| - |\mathbf{m}|$ and $q_j - a_j \geq 0$ ($j = 1, 2, 3$). Otherwise, it vanishes. By Theorem 4.1, we know $(\mathcal{P}f)(\mathbf{x})$ is k -monogenic on \mathbb{R}^{6n} . The corollary is proved.

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