

## On Hardy's Theorem on $SU(1, 1)^{***}$

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**Abstract** The classical Hardy theorem asserts that  $f$  and its Fourier transform  $\hat{f}$  can not both be very rapidly decreasing. This theorem was generalized on Lie groups and also for the Fourier-Jacobi transform. However, on  $SU(1, 1)$  there are infinitely many “good” functions in the sense that  $f$  and its spherical Fourier transform  $\tilde{f}$  both have good decay. In this paper, we shall characterize such functions on  $SU(1, 1)$ .

**Keywords** Heat kernel, Jacobi transform, Plancherel formula

**2000 MR Subject Classification** 22E30, 43A80, 43A90, 33C45

### 1 Introduction

The classical Hardy theorem (see [2]) asserts that  $f$  and its Fourier transform  $\tilde{f}$  can not both be “very rapidly decreasing”. More precisely, suppose that a measurable function  $f$  on  $\mathbb{R}$  and its Fourier transform  $\tilde{f}$  on  $\mathbb{R}$  satisfy

$$|f(x)| \leq Ae^{-ax^2} \quad \text{and} \quad |\tilde{f}(\lambda)| \leq Be^{-b\lambda^2} \quad (1.1)$$

for some positive constants  $A$ ,  $B$ ,  $a$  and  $b$ . If  $ab > \frac{1}{4}$ , then  $f = 0$ , and if  $ab = \frac{1}{4}$ , then  $f$  is a constant multiple of  $e^{-ax^2}$ . Recently, an analogue of Hardy's theorem was established on Lie groups by various people, where the heat kernel on Lie groups plays an essential role in controlling the decay of  $f$  and, in the case of  $ab = \frac{1}{4}$ , in expressing a unique function up to a constant multiplication. We refer to [9] and the references therein for more information. Moreover, Hardy's theorem was generalized for the Fourier-Jacobi transform (see [1, 3]) and, as an application, Andersen pointed out that Hardy's theorem on  $SU(1, 1)$  does not hold unless the  $K$ -type of  $f$  is fixed: Let  $G = SU(1, 1)$ , and for  $g \in G$  let  $g = k_\phi a_x k_\psi$ ,  $0 \leq x$ ,  $0 \leq \phi, \psi \leq 4\pi$ , denote the Cartan decomposition of  $g$ . Let  $h_t$  denote the heat kernel on  $G$  and for integrable functions  $f$  on  $G$  let  $\tilde{f}_{n,m}$ ,  $n, m \in \frac{1}{2}\mathbb{Z}$ , the spherical Fourier transform of  $f$  corresponding to the  $K$ -type  $(n, m)$  (see (2.10) below). We suppose that a measurable function  $f$  on  $G$  and its spherical Fourier transform  $\tilde{f}_{n,m}$  on  $\mathbb{R}$  satisfy

$$|f_{n,m}(g)| \leq Ah_{1/(4a)}(g) \quad \text{and} \quad |\tilde{f}_{n,m}(\lambda)| \leq Be^{-b\lambda^2} \quad \text{for all } n, m \in \frac{1}{2}\mathbb{Z} \quad (1.2)$$

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for some positive constants  $A, B, a$  and  $b$ . Then,  $f = 0$  if  $ab > \frac{1}{4}$ . However, there are infinitely many linearly independent functions on  $G$  satisfying the above condition if  $ab = \frac{1}{4}$  (see Corollary 4.1).

In this paper, we restrict our attention to functions on  $G$  with  $K$ -types  $(n, m)$ ,  $n, m = 0, 1, 2, \dots$ , and we show that the condition (1.2) under  $ab = \frac{1}{4}$  determines a function on  $G$  uniquely in the following sense: In the classical case the condition (1.1) under  $ab = \frac{1}{4}$  guarantees the limit

$$\lim_{x \rightarrow \infty} e^{ax^2} f(x) = c$$

and then  $f$  is uniquely determined as  $f(x) = ce^{-ax^2}$ . On  $\mathrm{SU}(1, 1)$ , similarly, the condition (1.2) under  $ab = \frac{1}{4}$  guarantees the limit

$$\lim_{x \rightarrow \infty} h_{1/(4a)}(x)^{-1} f(k_\phi a_x) = F(\phi)$$

and then  $f$  is uniquely determined by using the Fourier coefficients of  $F$ . Here  $F \in H^2(\mathbb{T})$  is real analytic. Moreover, the  $L^2$ -norm of  $F$  on  $\mathbb{T}$  coincides with the  $L^2$ -norm of the principal part of  $f$  on  $G$  and the Fourier coefficients  $\{d_n; n = 0, 1, 2, \dots\}$  of  $F$  satisfy

$$\sum_{n=0}^{\infty} |d_n|^2 \left( 1 + \sum_{k=0}^{n-1} k e^{2b(2k+1)^2} \right) < \infty.$$

In Theorem 5.1 we shall give a characterization of  $F$ .

## 2 Notations

Let  $G = \mathrm{SU}(1, 1)$  and  $A, K$  the subgroups of  $G$  of the matrices

$$a_x = \begin{pmatrix} \cosh \frac{x}{2} & \sinh \frac{x}{2} \\ \sinh \frac{x}{2} & \cosh \frac{x}{2} \end{pmatrix}, \quad x \in \mathbb{R} \quad \text{and} \quad k_\phi = \begin{pmatrix} e^{i\phi/2} & 0 \\ 0 & e^{-i\phi/2} \end{pmatrix}, \quad 0 \leq \phi \leq 4\pi$$

respectively. According to the Cartan decomposition of  $G$ , each  $g \in G$  can be written uniquely as  $g = k_\phi a_x k_\psi$  where  $0 \leq x$ ,  $0 \leq \phi, \psi \leq 4\pi$ . Let  $\pi_{j,\lambda}$  ( $j = 0, \frac{1}{2}, \lambda \in \mathbb{R}$ ) denote the principal series representation of  $G$ . Then the (vector-valued) spherical Fourier transform  $\pi_{j,\lambda}(f)$  of  $f$  on  $G$  is defined as  $\pi_{j,\lambda}(f) = \int_G f(g) \pi_{j,\lambda}(g) dg$ , where  $dg$  a Haar measure on  $G$ . In the following, we shall consider functions  $f$  on  $G$  satisfying

$$f(a_x) = f(a_{-x}), \quad x \in \mathbb{R}$$

and we identify  $f$  with an even function on  $\mathbb{R}$ , which is denoted by the same symbol  $f$ . Under this restriction, we may suppose that  $\pi_{j,\lambda}(f)$  is supported on  $j = 0$  and  $\lambda > 0$  and the  $K$ -types  $(m, n)$  of  $f$  is supported on  $m, n \in \mathbb{Z}$  (cf. [6] and [8, §8]).

Before introducing the explicit form of the spherical Fourier transform of  $f$  on  $G$ , we shall recall the theory of the Jacobi transform on  $\mathbb{R}_+$  (see [4, 5]). Let  $\alpha, \beta, \lambda \in \mathbb{C}$  and  $x \in \mathbb{R}$  and consider the differential equation

$$(L_{\alpha,\beta} + \lambda^2 + \rho^2)f(x) = 0, \tag{2.1}$$

where  $\rho = \alpha + \beta + 1$  and

$$L_{\alpha, \beta} = \frac{d^2}{dx^2} + ((2\alpha + 1) \coth x + (2\beta + 1) \tanh x) \frac{d}{dx}.$$

Then, for  $\alpha \notin -\mathbb{N}$ , the Jacobi function of the first kind with order  $(\alpha, \beta)$

$$\phi_{\lambda}^{\alpha, \beta}(x) = F\left(\frac{\rho + i\lambda}{2}, \frac{\rho - i\lambda}{2}; \alpha + 1; -\sinh^2 x\right) \quad (2.2)$$

is a unique solution of (2.1) satisfying  $\phi_{\lambda}^{\alpha, \beta}(0) = 1$  and  $\frac{d\phi_{\lambda}^{\alpha, \beta}}{dx}(0) = 0$ . For  $\lambda \notin -i\mathbb{N}$ , the Jacobi function of the second kind with order  $(\alpha, \beta)$

$$\Phi_{\lambda}^{\alpha, \beta}(t) = (e^t - e^{-t})^{i\lambda - \rho} F\left(\frac{\rho - 2\alpha - i\lambda}{2}, \frac{\rho - i\lambda}{2}; 1 - i\lambda; -\sinh^{-2} t\right) \quad (2.3)$$

is another solution of (2.1). Then  $\Gamma(\alpha + 1)^{-1} \phi_{\lambda}^{\alpha, \beta}$  is entire of  $\alpha, \beta$ , and for  $\lambda \notin i\mathbb{Z}$ , we have the identity

$$\frac{\sqrt{\pi}}{\Gamma(\alpha + 1)} \phi_{\lambda}^{\alpha, \beta}(t) = \frac{1}{2} (C_{\alpha, \beta}(\lambda) \Phi_{\lambda}^{\alpha, \beta}(t) + C_{\alpha, \beta}(-\lambda) \Phi_{-\lambda}^{\alpha, \beta}(t)), \quad (2.4)$$

where  $C_{\alpha, \beta}(\lambda)$  is the  $C$ -function given by

$$C_{\alpha, \beta}(\lambda) = \frac{2^{\rho} \Gamma(\frac{i\lambda}{2}) \Gamma(\frac{1+i\lambda}{2})}{\Gamma(\frac{\rho+i\lambda}{2}) \Gamma(\frac{\rho-2\beta+i\lambda}{2})} \quad (2.5)$$

(see [4, (2.5), (2.6)]). For convenience we assume  $\alpha > -1$  and  $\beta \in \mathbb{R}$  in the following. Then  $C_{\alpha, \beta}(-\lambda)^{-1}$  has only simple poles for  $\Im \lambda \geq 0$  which lie in the finite set  $D_{\alpha, \beta} = \{i(|\beta| - \alpha - 1 - 2m); m = 0, 1, 2, \dots, |\beta| - \alpha - 1 - 2m > 0\}$ . We denote the residue of  $(C_{\alpha, \beta}(\lambda) C_{\alpha, \beta}(-\lambda))^{-1}$  at  $\gamma \in D_{\alpha, \beta}$  by

$$d_{\alpha, \beta}(\gamma) = -i \operatorname{Res}_{\lambda=\gamma} (C_{\alpha, \beta}(\lambda) C_{\alpha, \beta}(-\lambda))^{-1}.$$

Let  $f$  be a compactly supported  $C^{\infty}$  even function on  $\mathbb{R}$ . We define the Jacobi transform  $\widehat{f}_{\alpha, \beta}(\lambda)$  by

$$\widehat{f}_{\alpha, \beta}(\lambda) = \frac{\sqrt{2}}{\Gamma(\alpha + 1)} \int_0^{\infty} f(x) \phi_{\lambda}^{\alpha, \beta}(x) \Delta_{\alpha, \beta}(x) dx, \quad (2.6)$$

where  $\Delta_{\alpha, \beta}(x) = (2 \sinh x)^{2\alpha+1} (2 \cosh x)^{2\beta+1}$  (see [4, (3.2)] and [5, (2.12)]). Then the inversion formula and the Plancherel formula are respectively given as follows:

$$f(x) = \frac{\sqrt{2}}{\Gamma(\alpha + 1)} \left( \int_0^{\infty} \widehat{f}_{\alpha, \beta}(\lambda) \phi_{\lambda}^{\alpha, \beta}(x) |C_{\alpha, \beta}(\lambda)|^{-2} d\lambda + \sum_{\gamma \in D_{\alpha, \beta}} \widehat{f}_{\alpha, \beta}(\gamma) \phi_{\gamma}^{\alpha, \beta}(x) d_{\alpha, \beta}(\gamma) \right), \quad (2.7)$$

$$\int_0^{\infty} |f(x)|^2 \Delta_{\alpha, \beta}(x) dx = \int_0^{\infty} |\widehat{f}_{\alpha, \beta}(\lambda)|^2 |C_{\alpha, \beta}(\lambda)|^{-2} d\lambda + \sum_{\gamma \in D_{\alpha, \beta}} |\widehat{f}_{\alpha, \beta}(\gamma)|^2 d_{\alpha, \beta}(\gamma) \quad (2.8)$$

(see [4, Theorem 4.2, (5.1)] and [5, Theorems 2.3 and 2.4]).

Let  $h_t^{\alpha, \beta}$  denote the heat kernel for the Jacobi transform, that is, an even function on  $\mathbb{R}$  satisfying

$$(h_t^{\alpha, \beta})_{\alpha, \beta}^{\wedge}(\lambda) = e^{-t(\lambda^2 + \rho^2)}, \quad t, \lambda \in \mathbb{R}. \quad (2.9)$$

We return to harmonic analysis on  $SU(1, 1)$ . Let  $n, m \in \mathbb{Z}$  and  $\psi_\lambda^{n,m}(g)$  ( $\lambda \in \mathbb{R}$ ,  $g \in G$ ) denote the matrix coefficient of  $\pi_{0,\lambda}(g)$  with  $K$ -type  $(n, m)$ . Let  $f$  be a compactly supported  $C^\infty$  function on  $G$ . We define the scalar-valued spherical Fourier transform  $\tilde{f}_{n,m}(\lambda)$  by

$$\tilde{f}_{n,m}(\lambda) = \int_0^\infty f(x) \psi_\lambda^{(n,m)}(x) \Delta_{0,0}(x) dx. \quad (2.10)$$

We recall that the explicit form of  $\psi_\lambda^{n,m}(g)$  is given by using the Jacobi function (2.2) (cf. [5, (4.17)] and [6, (3.4.10)]): for  $g = k_\phi a_x k_\psi \in G$ ,

$$\psi_\lambda^{n,m}(g) = (\cosh x)^{n+m} (\sinh x)^{|n-m|} Q_{n,m}(\lambda) \phi_\lambda^{|n-m|, n+m}(x) e^{in\phi} e^{im\psi},$$

where  $Q_{n,m}(\lambda)$  can be expressed by binomial coefficient as

$$Q_{n,m}(\lambda) = \binom{-\frac{1}{2} - \frac{i\lambda}{2} \mp m}{|n-m|}, \quad (2.11)$$

and  $\mp m$  is equal to  $-m$  if  $m \geq n$  and  $m$  if  $m \leq n$ . Hence from (2.6) and (2.9) it follows that

$$\begin{aligned} \tilde{f}_{n,m}(\lambda) &= 2^{-2(n+m)-2|n-m|} Q_{n,m}(\lambda) \\ &\times \left( f(x) (\sinh x)^{-|n-m|} (\cosh x)^{-(n+m)} \right)_{|n-m|, n+m}^\wedge(\lambda). \end{aligned} \quad (2.12)$$

We shall consider the case of  $n = m$ . Let  $F$  be a compactly supported  $C^\infty$  even function on  $\mathbb{R}$ . We put

$$f(g) = F(x) (\cosh x)^{2n} e^{in(\phi+\psi)}, \quad g = k_\phi a_x k_\psi \in G. \quad (2.13)$$

Then letting  $\alpha = 0, \beta = 2n$  in (2.8) and (2.11), we see that

$$\int_0^\infty |f(x)|^2 \Delta_{0,0}(x) dx = \int_0^\infty |\tilde{f}_{n,n}(\lambda)|^2 |C_{0,0}(\lambda)|^{-2} d\lambda + \sum_{k=0}^{|n|-1} \left(k + \frac{1}{2}\right) |\tilde{f}_{n,n}((2k+1)i)|^2 \quad (2.14)$$

(see [6, (4.21)] and [8, Theorem 8.2]). This is nothing but the Plancherel formula for central compactly supported  $C^\infty$  functions on  $G$ . We denote by  $f_P$  and  ${}^\circ f$  respectively the principal part and discrete part of  $f$  on  $G$ ;

$$f = f_P + {}^\circ f.$$

Then (2.14) corresponds to the relation  $\|f\|_{L^2(G)}^2 = \|f_P\|_{L^2(G)}^2 + \|{}^\circ f\|_{L^2(G)}^2$ .

### 3 Asymptotic Behavior of Heat Kernels

When  $\alpha \geq \beta \geq -\frac{1}{2}$ , the asymptotic behavior of  $h_t^{\alpha,\beta}(x)$  is well-known (see [1] and [3, Theorem 3.1]). In particular,

$$h_t^{0,0}(x) \sim t^{-1} e^{-\rho^2 t} e^{-\rho x} e^{-x^2/(4t)} (1+t+x)^{-1/2} (1+x). \quad (3.1)$$

In this section we shall treat the case of  $\alpha, \beta = 0, 1, 2, \dots$ , and we shall investigate a leading term of  $h_t^{\alpha,\beta}(x)$  when  $x \rightarrow \infty$ . In the following, we fix  $t > 0$  and we denote  $a = \frac{1}{4t}$  for simplicity.

For an even function  $f$  on  $\mathbb{R}$  let  $W_\mu^\sigma(f)$ ,  $\mu \in \mathbb{C}$ ,  $\sigma > 0$ , denote the Weyl type fractional integral of  $f$ , which is defined by

$$W_\mu^\sigma(f)(y) = \Gamma(\mu)^{-1} \int_y^\infty f(x) (\cosh \sigma x - \cosh \sigma y)^{\mu-1} d(\cosh \sigma x) \quad (3.2)$$

for  $\Re \mu > 0$  and is extended to an entire function in  $\mu$  (see [4, (3.10), (3.11)]). Then it is known that

$$\widehat{f}_{\alpha, \beta}(\lambda) = \mathcal{F}(2^{3\alpha+3/2} W_{\alpha-\beta}^1 \circ W_{\beta+1/2}^2(f)),$$

where  $\mathcal{F}$  denotes the Euclidean Fourier transform (see [4, (3.7), (3.12)]). Therefore, letting  $\alpha = \beta = 0$ , it follows from (2.9) that

$$W_{1/2}^2(h_t^{0,0})(x) = \frac{1}{2^2 \sqrt{t}} e^{-t} e^{-ax^2} \quad (3.3)$$

and moreover, letting  $\alpha = m$ ,  $\beta = n$ ,  $2^{3m+3/2} e^{t(m+n+1)^2} W_{m-n}^1 \circ W_{n+1/2}^2(h_t^{m,n})$  does not depend on  $m, n$ . Hence, it follows that

$$h_t^{m,n} = 2^{-3m} e^{-t((m+n+1)^2-1)} W_{-1/2-n}^2 \circ W_{n-m}^1 \circ W_{1/2}^2(h_t^{0,0}). \quad (3.4)$$

**Lemma 3.1** For  $n = 0, 1, 2, \dots$ ,

$$W_{-n}^2 \circ W_n^1(f)(x) = \sum_{l=0}^{n-1} c_l^n (\cosh x)^{-(n+l)} W_l^1(f)(x), \quad (3.5)$$

where  $4c_l^n = c_l^{n-1} - (n+l-2)c_{l-1}^{n-1}$ . In particular,  $c_0^n = 2^{-2n}$ ,  $|c_{n-1}^n| = 2^{-2n}(2n-3)!!$ ,  $c_l^n > 0$  if  $l$  is even and  $c_l^n < 0$  if  $l$  is odd, and

$$|c_l^n| \leq \frac{(2n-3)!!}{2^{2n}(n-1-l)!}, \quad 0 < l \leq n-1.$$

**Proof** Since

$$W_{-1}^2 = \frac{1}{2 \sinh 2x} \frac{d}{dx} = \frac{1}{4 \cosh x} W_{-1}^1,$$

(3.5) and the recursive relation  $4c_l^n = c_l^{n-1} - (n+l-2)c_{l-1}^{n-1}$  follows from the induction on  $n$ . In particular,  $4c_0^n = c_0^{n-1}$  and  $4c_{n-1}^n = -(2n-3)c_{n-2}^{n-1}$ , and thus,  $c_0^n = 2^{-2n} > 0$  and  $|c_{n-1}^n| = 2^{-2n}(2n-3)!!$ . The sign of general  $c_l^n$  follows from the recursive relation. Since  $4^{n-1}(2n-3)|c_{l-1}^{n-1}| \leq 4^n |c_l^n|$ , it follows that  $4|c_l^n| \leq |c_l^{n-1}| + \frac{4(n+l-2)}{2n-3} |c_l^n|$  and thus,  $4|c_l^n| \leq \frac{2n-3}{n-l-1} |c_l^{n-1}|$ . This means that

$$|c_l^n| \leq \frac{(2n-3)!!}{2^{2(n-l-1)}(n-l-1)!(2l-1)!!} |c_l^{l+1}| = \frac{(2n-3)!!}{2^{2n}(n-1-l)!}.$$

**Lemma 3.2.** Let  $p, q \geq 0$  and suppose  $q = 0$  if  $p = 0$ . Then there exists a positive constant  $c$  such that for all  $l = 0, 1, 2, \dots$ , and  $x \geq \max(1, \frac{1}{a})$ ,

$$W_l^1(e^{-ax^2-px}x^q) \leq c 2^{-l} (2a)^l e^{l^2/(4a)} e^{-ax^2-(p-l)x} x^{q-l}.$$

Here, if  $l \geq 1$ , then  $(2a)^l x^{-l}$  can be replaced by  $\Gamma(l)^{-1}(2a)^{-1}x^{-1}$ . Moreover, if  $2^{-l}e^{l^2/(4a)}$  is replaced by  $2^{-2l}$ , then the lower bound follows.

**Proof** The case of  $l = 0$  is obvious, so we may suppose that  $l \geq 1$ . Since  $2 \sinh x \leq e^x$ ,  $s + x \leq (1 + s)(1 + x) \leq 2x(1 + s)$  for  $x \geq 1$ ,  $s \geq 0$ , and  $e^{-px}(1 + s)^q \leq c$ , it follows from (3.2) that

$$\begin{aligned}
 & \Gamma(l)W_l^1(e^{-ax^2-px}x^q) \\
 &= \int_x^\infty e^{-as^2-ps}s^q(\cosh s - \cosh x)^{l-1} \sinh s ds \\
 &= e^{-ax^2-px} \int_0^\infty e^{-as^2-ps-2axs}(x+s)^q \left(2 \sinh\left(\frac{s}{2} + x\right) \sinh\left(\frac{s}{2}\right)\right)^{l-1} \sinh(s+x) ds \\
 &\leq ce^{-ax^2-(p-l)x}x^q \int_0^\infty e^{-as^2+(l/2+1/2)s-2axs} \left(\sinh\left(\frac{s}{2}\right)\right)^{l-1} ds \\
 &\leq c2^{-l}e^{-ax^2-(p-l)x}x^q \int_0^\infty e^{-as^2+ls-2axs} ds \\
 &= c2^{-l}e^{l^2/(4a)}e^{-ax^2-(p-l)x}x^q(2ax)^{-1}, \tag{3.6}
 \end{aligned}$$

where we used the fact that  $-as^2 + ls = -a(s - \frac{l}{2a})^2 + \frac{l^2}{4a}$ . Since  $\sinh x \leq xe^x$ , the integral in (3.6) is also estimated as

$$\leq c2^{-l}e^{l^2/(4a)}c \int_0^\infty e^{-2axs}s^{l-1}ds = c2^{-l}ce^{l^2/(4a)}\Gamma(l)(2ax)^{-l}.$$

Therefore, we can deduce the first estimate.

We note that  $\sinh x \geq \frac{xe^x}{2}$  for  $0 \leq x \leq \frac{1}{2}$ . Since  $0 \leq \frac{1}{2ax} \leq \frac{1}{2}$  and  $s + x \geq x$  for  $s, x \geq 0$ , it follows that

$$\begin{aligned}
 \Gamma(l)W_l^1(e^{-ax^2-px}x^q) &\geq c2^{-2l}e^{-ax^2-(p-l)x}x^q \int_0^{1/(2bx)} e^{-as^2-(p-l)s-2axs}s^{l-1}ds \\
 &\geq c2^{-2l}e^{-ax^2-(p-l)x}(2ax)^{-l} \int_0^1 e^{-s}s^{l-1}ds.
 \end{aligned}$$

Since  $\Gamma(l)^{-1} \int_0^1 e^{-s}s^{l-1}ds$  is bounded below, the lower estimate follows.

**Lemma 3.3** Let  $p, q \geq 0$  and suppose  $q = 0$  if  $p = 0$ . Then there exist a positive constants  $c$  such that for all  $l, n = 0, 1, 2, \dots$ , and  $x \geq \max(1, \frac{1}{a})$ ,

$$\begin{aligned}
 & |W_{-1/2}^2((\cosh x)^{n+l}W_l^1(e^{-ax^2-px}x^q))| \\
 &\leq c2^n(2a)^{-l}e^{l^2/(4a)}e^{-ax^2-(n+1+p)x}x^{q-l}((n+l)x^{-1/2} + x^{1/2}).
 \end{aligned}$$

Here, if  $l \geq 1$ , then  $(2a)^{-l}x^{-l}$  can be replaced by  $\Gamma(l)^{-1}(2a)^{-1}x^{-1}$ . Moreover, if  $2^n e^{l^2/(4a)}$  is replaced by  $2^{n-l}$ , then the lower bound follows.

**Proof** Since  $W_{-1/2}^2 = W_{1/2}^2 \circ W_{-1}^2$ , it follows that

$$\begin{aligned}
 & W_{-1/2}^2 \circ ((\cosh x)^{-(n+l)}W_l^1) \\
 &= \frac{1}{4}W_{1/2}^2 \circ (-(n+l)(\cosh x)^{-(n+l+2)}W_l^1 + (\cosh x)^{-(n+l+1)}W_{l-1}^1). \tag{3.7}
 \end{aligned}$$

Therefore, we need to estimate

$$W_{1/2}^2((\cosh x)^{-(n+l+2)}W_l^1(e^{-ax^2-px}x^q)), \quad l = -1, 0, 1, 2, \dots$$

Substituting the estimate obtained in Lemma 3.2, we see that for  $l \geq 0$ ,

$$\begin{aligned} & W_{1/2}^2((\cosh x)^{-(n+l+2)}W_l^1(e^{-ax^2-px}x^q)) \\ & \leq cc_0 \int_x^\infty e^{-as^2-(n+2+p)s} s^{q-l} (\cosh 2s - \cosh 2x)^{-1/2} 2 \sinh 2s ds \\ & \leq cc_0 e^{-ax^2-(n+p)x} \int_0^\infty e^{-as^2-(n+p+2ax)s} (s+x)^{q-l} (\cosh 2(s+x) - \cosh 2x)^{-1/2} ds, \end{aligned}$$

where  $c_0 = c2^n(2a)^l e^{l^2/(4a)}$ . We note that, if  $l \geq q$ , then  $(s+x)^{q-l} \leq x^{q-l}$ , and if  $l \leq q$ , then  $(s+x)^{q-l} \leq (2x(1+s))^{q-l}$  and  $e^{-ps}(1+s)^{q-l} \leq e^{-ps}(1+s)^q \leq c$ . Therefore, applying [3, (3.1)] to  $(\cosh 2(s+x) - \cosh 2x)^{-1/2}$ , we have that the above formula could be estimated as

$$\begin{aligned} & \leq cc_0 e^{-ax^2-(n+1+p)x} x^{q-l} \int_0^\infty e^{-ax^2-(n+1+2ax)s} \left( \frac{1+2(x+s)}{s(x+s)} \right)^{1/2} ds \\ & \leq cc_0 e^{-ax^2-(n+1+p)x} x^{q-l} \left( \frac{1}{x} + 1 \right)^{1/2} \int_0^\infty e^{-2axs} \frac{1}{\sqrt{s}} ds \\ & \leq cc_0 e^{-ax^2-(n+1+p)x} x^{q-l-1/2}. \end{aligned} \tag{3.8}$$

When  $l = -1$ , we note that  $|W_{-1}^1(e^{-ax^2-px}x^q)| \leq c(1+x)^{q+1}e^{-ax^2-px}(\sinh x)^{-1}$ . Hence, (3.8) is replaced by

$$\begin{aligned} & \leq c2^n e^{-ax^2-(n+p)x} \int_0^\infty e^{-as^2-(n+p+2ax)s} (s+x)^{q+1} (\cosh(s+x) - \cosh x)^{-1/2} ds \\ & \leq c2^n e^{-ax^2-(n+1+p)x} x^q \int_0^\infty e^{-2axs} \frac{1}{\sqrt{s}} ((x+s)(1+2(x+s)))^{1/2} ds. \end{aligned}$$

The last integral is dominated by  $x^{1/2}$ . Substituting these estimates to (3.7), we can deduce the desired upper estimate. Other desired estimates follow from Lemma 3.2 and the arguments used in [3, Theorem 3.1].

When  $p = q = 0$  and  $l = 0$ , we have the following refinement.

**Lemma 3.4** For all  $n = 0, 1, 2, \dots$ ,

$$W_{-1/2}^2((\cosh x)^{-n}e^{-ax^2}) = c_0(\cosh x)^{-n}h_t^{0,0}(x) + O(2^n n e^{-ax^2-(n+1)x}x^{-1/2}),$$

where  $c_0 = 2^2\sqrt{t}e^t$ . Here  $f = O(g)$  means that  $|\frac{f(x)}{g(x)}| \leq C$  when  $x \rightarrow \infty$ . If  $C$  depends on some parameters  $\gamma$ , then we use the symbol  $f = O_{(\gamma)}(g)$ .

**Proof** Since  $e^{-ax^2} = c_0 W_{1/2}^2(h_t^{0,0})$  (see (3.3)), the case of  $n = 0$  is obvious and moreover, for  $n \geq 1$ , it follows that

$$\begin{aligned} & c_0^{-1} W_{-1/2}^2((\cosh x)^{-n}e^{-ax^2}) \\ & = - \int_x^\infty \frac{d}{d \cosh 2s} (((\cosh x)^{-n} - (\cosh s)^{-n}) W_{1/2}^2(h_t^{0,0})(s)) (\cosh 2s - \cosh 2x)^{-1/2} d \cosh 2s \\ & \quad + (\cosh x)^{-n} h_t^{0,0}(x) \\ & = - \int_x^\infty ((\cosh x)^{-n} - (\cosh s)^{-n}) e^{-as^2} (\cosh 2s - \cosh 2x)^{-3/2} 2 \sinh 2s ds + (\cosh x)^{-n} h_t^{0,0}(x). \end{aligned}$$

We note that for  $0 \leq x \leq s$ ,

$$(\cosh x)^{-n} - (\cosh s)^{-n} \leq \frac{n(\cosh s - \cosh x)}{\cosh x (\cosh x)^n}.$$

Therefore, the similar argument in the proof of Lemma 3.3 (or [3, Theorem 3.1]) yields that the last integral is dominated by  $2^n n e^{-ax^2 - (n+1)x} x^{-1/2}$ .

Now we shall obtain the asymptotic behavior of  $h_t^{m,n}(x)$  as  $x \rightarrow \infty$ . It follows from (3.3), (3.4) and (3.5) that

$$h_t^{m,n} = c_0^{-1} 2^{-3m} e^{-t((m+n+1)^2-1)} \sum_{l=0}^{n-1} c_l^n W_{-1/2}^2((\cosh x)^{-(n+l)} W_{l-m}^1(e^{-ax^2})). \quad (3.9)$$

Since

$$W_{-m}^1(e^{-ax^2}) \sim_{(m)} (2ax)^m e^{-ax^2 - mx},$$

here  $f \sim g$  means that there exist positive constants  $c_1, c_2$  such that  $c_1 f(x) \leq g(x) \leq c_2 f(x)$ . If  $c_1, c_2$  depend on some parameters  $\gamma$ , then we use the symbol  $f \sim_{(\gamma)} g$ . Lemma 3.4 implies that, when  $x \rightarrow \infty$ , the term corresponding to  $l = 0$  contributes to the asymptotic behavior of  $h_t^{m,n}(x)$ :

**Proposition 3.1** *We fix  $t > 0$  and  $m, n = 0, 1, 2, \dots$ . Then for  $x, ax \geq 1$ ,*

$$h_t^{m,n}(x) \sim_{(t,m,n)} e^{-\rho^2 t} e^{-\rho x} e^{-x^2/(4t)} (1+x)^{m+1/2}. \quad (3.10)$$

Next we shall consider the behavior of  $(\cosh x)^n h_t^{0,n}(x)$ . Let  $\epsilon > 0$  and we suppose that

$$x \geq \frac{1}{2} \log \left( \frac{1}{2^\epsilon - 1} \right) = x(\epsilon),$$

that is,  $\cosh x \leq 2^{-1+\epsilon} e^x$  if  $x \geq x(\epsilon)$  and  $x \rightarrow \infty$  if  $\epsilon \rightarrow 0$ . Then it follows from (3.9), Lemmas 3.3 and 3.4 that

$$(\cosh x)^n h_t^{0,n}(x) = 2^{n\epsilon} e^{-t((n+1)^2-1)} \left( c_0^n h_t^{0,0}(x) + O \left( \sum_{l=1}^{n-1} |c_l^n| e^{l^2/(4a)} \Gamma(l)^{-1} e^{-ax^2 - x} n x^{-1/2} \right) \right).$$

We note that  $e^{l^2/(4a)} \leq e^{(n-1)^2/(4a)}$  and

$$\sum_{l=1}^{n-1} |c_l^n| \Gamma(l)^{-1} \leq \sum_{l=1}^{n-1} \frac{(n-1)!}{2^{n+2}(n-l-1)! \Gamma(l)} = 2^{-4}(n-1).$$

Hence, it follows that

$$\begin{aligned} (\cosh x)^n h_t^{0,n}(x) &= 2^{n\epsilon} e^{-t((n+1)^2-1)} (c_0^n h_t^{0,0}(x) + O(e^{(n-1)^2/(4a)} e^{-ax^2 - x} n^2 x^{-1/2})) \\ &= 2^{n\epsilon} 2^{-2n} e^{-t((n+1)^2-1)} h_t^{0,0}(x) (1 + O(n^2 2^{2n} e^{(n-1)^2/(4a)} x^{-1})). \end{aligned} \quad (3.11)$$

Letting  $x \rightarrow \infty$ , we have the following

**Proposition 3.2** *We fix  $t > 0$  and  $n = 0, 1, 2, \dots$ . Then*

$$\lim_{x \rightarrow \infty} \frac{(\cosh x)^n h_t^{0,n}(x)}{h_t^{0,0}(x)} = 2^{-2n} e^{-t((n+1)^2-1)}. \quad (3.12)$$



## 4 Hardy's Theorem

We keep the notations in the previous section. We recall the proof of Hardy's theorem for the Jacobi transform of  $(\alpha, \beta)$ ,  $\alpha \geq \beta \geq -\frac{1}{2}$  (see [1, 3]). Then it is easy to see that Hardy's theorem for the Jacobi transform of  $(m, n)$ ,  $m, n = 0, 1, 2, \dots$ , also holds:

**Theorem 4.1** *Let  $m, n = 0, 1, 2, \dots$ , and  $f$  be a measurable function on  $\mathbb{R}_+$  satisfying*

$$(i) \quad g(x) = O_{(m,n)}(h_{1/(4a)}^{m,n}(x)),$$

$$(ii) \quad \hat{g}_{m,n}(\lambda) = O_{(m,n)}(e^{-b\lambda^2}).$$

*If  $ab > \frac{1}{4}$ , then  $g = 0$ , and if  $ab = \frac{1}{4}$ , then  $g$  is a constant multiple of  $h_b^{m,n}(x)$ .*

Applying this theorem, we have Hardy's theorem on  $SU(1, 1)$  for a fixed  $K$ -type (see the example in [7]).

**Theorem 4.2** *Let  $f$  be a measurable function on  $G$  of  $K$ -type  $(n, m)$ ,  $n, m = 0, 1, 2, \dots$ , satisfying*

$$(i) \quad f(x) = O_{(n,m)}(h_{1/(4a)}^{|n-m|,n+m}(x)(\sinh x)^{|n-m|}(\cosh x)^{n+m}),$$

$$(ii) \quad \tilde{f}_{n,m}(\lambda) = O_{(n,m)}(Q_{n,m}(\lambda)e^{-b\lambda^2}).$$

*If  $ab > \frac{1}{4}$ , then  $f = 0$ , and if  $ab = \frac{1}{4}$ , then  $f$  is a constant multiple of  $h_b^{|n-m|,n+m}(x) \cdot (\sinh x)^{|n-m|}(\cosh x)^{n+m}$ .*

**Proof** Let  $g(x) = f(x)(\sinh x)^{-|n-m|}(\cosh x)^{-(n+m)}$ . Then

$$g(x) = O_{(n,m)}(h_{1/(4a)}^{|n-m|,n+m}(x)),$$

$$\hat{g}_{|n-m|,n+m}(\lambda) = 2^{2(n+m)+2|n-m|}\tilde{f}_{n,m}(\lambda)Q_{n,m}^{-1}(\lambda) = O_{(n,m)}(e^{-b\lambda^2})$$

(see (2.12)). Then Theorem 4.1 implies that, if  $ab > \frac{1}{4}$ , then  $g = 0$ , and if  $ab = \frac{1}{4}$ , then  $g$  is a constant multiple of  $h_b^{|n-m|,n+m}(x)$  and thus,  $f$  is the desired form.

Let  $L_{0+}^2(G)$  denote the subspace of  $L^2(G)$  consisting of all  $f$  of the form

$$f = \sum_{n,m=0}^{\infty} f_{n,m},$$

where  $f_{n,m}$  is of  $K$ -type  $(n, m)$ .

**Corollary 4.1** *Let  $f$  be in  $L_{0+}^2(G)$  and satisfy for all  $n, m = 0, 1, 2, \dots$ ,*

$$(i) \quad f_{n,m}(x) = O_{(n,m)}(h_{1/(4a)}^{0,0}(x)),$$

$$(ii) \quad \tilde{f}_{n,m}(\lambda) = O_{(n,m)}(e^{-b\lambda^2}).$$

*If  $ab > \frac{1}{4}$ , then  $f = 0$ , and if  $ab = \frac{1}{4}$ , then  $f$  is of the form*

$$f(g) = \sum_{n=0}^{\infty} a_n h_b^{0,2n}(x)(\cosh x)^{2n} e^{in(\phi+\psi)}, \quad (4.1)$$

where  $g = k_\phi a_x k_\psi$  and  $a_n \in \mathbb{C}$ .

**Proof** Proposition 3.1 implies that

$$h_t^{|n-m|,n+m}(x)(\sinh x)^{|m-n|}(\cosh x)^{m+n} \sim_{(m,n)} h_t^{0,0}(x)(1+x)^{|m-n|} \quad \text{for } x, ax \geq 1. \quad (4.2)$$

Hence  $f_{n,m}(x) = O_{(n,m)}(h_{1/(4a)}^{|n-m|,n+m}(x)(\sinh x)^{|n-m|}(\cosh x)^{n+m}(1+x)^{-|n-m|})$  and  $\tilde{f}_{n,m}(\lambda) = O_{(n,m)}(e^{-b\lambda^2}) = O_{(n,m)}(Q_{n,m}(\lambda)e^{-b\lambda^2})$  (see (2.11)). Hence Theorem 4.2 implies that, if  $ab > \frac{1}{4}$ , then  $f_{n,m} = 0$ , for all  $n, m \in \mathbb{Z}$ , and thus  $f = 0$ . If  $ab = \frac{1}{4}$ , then  $f_{n,m}$  is a constant multiple of  $h_b^{|n-m|,n+m}(x)(\sinh x)^{|n-m|}(\cosh x)^{n+m}$ . Since  $f_{n,m}(x) = O_{(n,m)}(h_{1/(4a)}^{0,0}(x))$ , it follows that  $|n-m| = 0$  (see (4.2)). Thereby,  $f$  must be of the desired form.

## 5 Main Theorem

We retain the notations and suppose that

$$f(g) = \sum_{n=0}^{\infty} a_n h_b^{0,2n}(x)(\cosh x)^{2n} e^{in(\phi+\psi)} \in L_{0+}^2(G).$$

We recall that

$$\tilde{f}_{n,n}(\lambda) = a_n 2^{-4n} e^{-b((2n+1)^2 + \lambda^2)} \quad (5.1)$$

(see (2.9) and (2.12)). Then letting  $t = b = \frac{1}{4a}$  in (2.14), we obtain the  $L^2$ -norm of  $f$  on  $G$  as follows.

$$\begin{aligned} \int_G |f(g)|^2 dg &= \sum_{n=0}^{\infty} |a_n|^2 \int_0^{\infty} |h_b^{0,2n}(x)(\cosh x)^{2n}|^2 \Delta_{0,0}(x) dx \\ &= \sum_{n=0}^{\infty} |a_n|^2 2^{-8n} e^{-2b(2n+1)^2} \\ &\quad \times \left( \int_0^{\infty} e^{-2b\lambda^2} |C_{0,0}(\lambda)|^{-2} d\lambda + \sum_{k=0}^{n-1} \left(k + \frac{1}{2}\right) e^{2b(2k+1)^2} \right). \end{aligned} \quad (5.2)$$

We define the partial sum  $f_N$ ,  $N = 0, 1, 2, \dots$ , of  $f$  as

$$f_N(g) = \sum_{n=0}^N a_n h_b^{0,2n}(x)(\cosh x)^{2n} e^{in(\phi+\psi)}.$$

Then Proposition 3.2 implies that

$$\lim_{N \rightarrow \infty} \lim_{x \rightarrow \infty} h_b^{0,0}(x)^{-1} f_N(k_\phi a_x) = \sum_{n=0}^{\infty} a_n 2^{-4n} e^{-b((2n+1)^2 - 1)} e^{in\phi} = \sum_{n=0}^{\infty} d_n e^{in\phi} = F(\phi), \quad (5.3)$$

where  $d_n = 2^{-4n} e^{-b((2n+1)^2 - 1)} a_n$ . Obviously, (5.2) implies that

$$\|F\|_{L^2(\mathbb{T})} = c \|f_P\|_{L^2(G)} \quad \text{and} \quad \sum_{n=0}^{\infty} |d_n|^2 \left(1 + \sum_{k=0}^{n-1} k e^{2b(2k+1)^2}\right) \sim \|f\|_{L^2(G)}^2. \quad (5.4)$$

Since  $\sum_{n=0}^{\infty} |d_n|^2 < \infty$  and  $\sum_{n=1}^{\infty} |d_n|^2 (n-1) e^{2b(2n-1)^2} < \infty$ , there exists a positive constant  $C$  such that  $|d_n|(1+n)^{1/2} e^{b(2n-1)^2} \leq C$  for all  $n = 0, 1, 2, \dots$ . This means that  $F$  is real analytic and

$$|a_n| 2^{-4n} (1+n)^{1/2} e^{-b(2n+1)^2} e^{b(2n-1)^2} \leq C \quad \text{for all } n = 0, 1, 2, \dots$$

Hence (5.1) implies

$$\tilde{f}_{n,n}(\lambda) = O((1+n)^{-1/2} e^{-b(2n-1)^2} e^{-b\lambda^2}).$$

We introduce a subspace  $A_b^2(\mathbb{T})$  of  $H^2(\mathbb{T})$  as follows:

$$A_b^2(\mathbb{T}) = \left\{ F(\phi) = \sum_{n=0}^{\infty} d_n e^{in\phi} \in H^2(\mathbb{T}); \|F\|_{A_b^2(\mathbb{T})}^2 = \sum_{n=0}^{\infty} |d_n|^2 \left( 1 + \sum_{k=0}^{n-1} k e^{2b(2k+1)^2} \right) < \infty \right\}.$$

For  $F(\phi) = \sum_{n=0}^{\infty} d_n e^{in\phi} \in A_b^2(\mathbb{T})$ , we define a function  $f$  on  $G$  as (4.1) with  $a_n = 2^{4n} e^{b((2n+1)^2-1)} d_n$ . Then (5.2) and (5.4) imply  $\|f\|_{L^2(G)} \leq c \|F\|_{A_b^2(\mathbb{T})}^2$ . Clearly,  $|f_{n,n}(x)| = |a_n| h_b^{0,2n}(x) (\cosh x)^{2n} = O_{(n)}(h_b^{0,0}(x))$  and (5.1) implies  $\tilde{f}_{n,n}(\lambda) = O_{(n)}(e^{-b\lambda^2})$ . Moreover,

$$\lim_{N \rightarrow \infty} \lim_{x \rightarrow \infty} h_b^{0,0}(x)^{-1} f_N(k_\phi a_x) = F(\phi).$$

Finally, we have the following theorem.

**Theorem 5.1** *Let  $ab = \frac{1}{4}$ . Let  $f$  be in  $L_{0+}^2(G)$  and satisfy for all  $n, m = 0, 1, 2, \dots$ ,*

$$(i) \quad f_{n,m}(x) = O_{(n,m)}(h_{1/(4a)}^{0,0}(x)),$$

$$(ii) \quad \tilde{f}_{n,m}(\lambda) = O_{(n,m)}(e^{-b\lambda^2}).$$

*Then, as an  $L^2$ -function on  $\mathbb{T}$ ,*

$$(iii) \quad \lim_{N \rightarrow \infty} \lim_{x \rightarrow \infty} h_b^{0,0}(x)^{-1} f_N(k_\phi a_x) = F(\phi)$$

*exists and  $F \in A_b^2(\mathbb{T})$ . Here  $\|F\|_{L^2(\mathbb{T})} = c \|f_P\|_{L^2(G)}$  and  $\|F\|_{A_b^2(\mathbb{T})} \sim \|f\|_{L^2(G)}$ . Let  $F(\phi) = \sum_{n=0}^{\infty} d_n e^{in\phi}$  denote the Fourier series of  $F$ . Then  $f$  is uniquely determined as a central function*

$$f(g) = \sum_{n=0}^{\infty} d_n 2^{4n} e^{b((2n+1)^2-1)} h_b^{0,2n}(x) (\cosh x)^{2n} e^{in(\phi+\psi)},$$

*where  $g = k_\phi a_x k_\psi$ , and each  $\tilde{f}_{n,n}(\lambda)$  satisfies*

$$\tilde{f}_{n,n}(\lambda) = O((1+n)^{-1/2} e^{-b(2n-1)^2} e^{-b\lambda^2}).$$

*Conversely, if  $F \in A_b^2(\mathbb{T})$ , then there exists a function  $f \in L_{0+}^2(G)$  such that  $f$  satisfies (i), (ii) and (iii).*

**Remark 5.1** (1) We note that, if  $f \in L_{0+}^2(G)$  is of the form in (4.1) and  $F$  is given by Theorem 5.1(iii), then  $|f(g) - h_b^{0,0}(x)F(\phi)|$ ,  $g = k_\phi a_x k_\psi$ , is dominated as

$$\left( \sum_{n=0}^{\infty} |a_n| 2^{n\epsilon} e^{-t((n+1)^2-1)} n^2 e^{(n-1)^2/(4a)} \right) h_b^{0,0}(x) x^{-1}$$

(see (3.11)). Therefore, if this sum is finite, then we can replace Theorem 5.1(iii) by

$$(iii)' \quad \lim_{x \rightarrow \infty} h_b^{0,0}(x)^{-1} f(k_\phi a_x) = F(\phi),$$

and deduce that  $f \in L^1(G)$  and  $f(x) = O(h_b^{0,0}(x))$ .

(2) In Corollary 4.1 we can replace the condition (i) by

$$(i)' \quad f(x) = O(h_{1/(4a)}^{0,0}(x)),$$

because (i)' implies (i). Also, in Theorem 5.1, it is true if we ignore the last statement of the existence of  $f$  for  $F \in A_b^2(\mathbb{T})$ . As remarked in (1), in order to construct  $f \in L_{0+}^2(G)$  from  $F \in A_b^2(\mathbb{T})$ , which satisfies (i)', (ii) and (iii), it is necessary to control the series in (1).

(3) In Corollary 4.1 and Theorem 5.1, if we replace the condition (i) by

$$(i)'_P \quad (f_P)_{m,n}(x) = O_{(n,m)}(h_{1/(4a)}^{0,0}(x)),$$

then  $a_n = 0$  for  $n \neq 0$ , that is,  $f$  is  $K$ -biinvariant. Actually, since  $\tilde{f}_{m,n} = (f_P)_{m,n}$ , (i)'<sub>P</sub> and (ii) imply that  $f_P$  is of the form in (4.1). Since  $f_P$  has no discrete part, (5.2) implies that  $a_n$  must be 0 if  $n \neq 0$ .

## References

- [1] Andersen, N., Hardy's theorem for the Jacobi transform, *Hiroshima Math. J.*, **33**, 2003, 229–251.
- [2] Hardy, G. H., A theorem concerning Fourier transform, *J. London Math. Soc.*, **8**, 1933, 227–231.
- [3] Kawazoe, T. and Liu, J., Heat kernel and Hardy's theorem for Jacobi transform, *Chin. Ann. Math.*, **24B**(3), 2003, 359–366.
- [4] Koornwinder, T. H., A new proof of a Paley-Wiener type theorem for the Jacobi transform, *Ark. Mat.*, **13**, 1975, 145–159.
- [5] Koornwinder, T. H., Jacobi functions and analysis on noncompact semisimple Lie groups, *Special Functions: Group Theoretical Aspects and Applications*, R. A. Askey et al (eds.), D. Reidel Publishing Company, Dordrecht, 1984, 1–85.
- [6] Sally, P., Analytic Continuation of the Irreducible Unitary Representations of the Universal Covering Group of  $SL(2, \mathbb{R})$ , *Memoirs of the Amer. Math. Soc.*, Num., **69**, Amer. Math. Soc., Providence, Rhode Island, 1967.
- [7] Sarkar, R. P., A complete analogue of Hardy's theorem on semisimple Lie groups, *Colloq. Math.*, **93**, 2002, 27–40.
- [8] Sugiura, M., *Unitary Representations and Harmonic Analysis*, Second Edition, North-Holland, Amsterdam, 1990.
- [9] Thangavelu, S., *An Introduction to the Uncertainty Principle: Hardy's Theorem on Lie Groups*, Progress in Mathematics, Birkhäuser, Boston, 2003.