

# Notes on Homoclinic Solutions of the Steady Swift-Hohenberg Equation\*

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**Abstract** This paper considers the steady Swift-Hohenberg equation

$$u'''' + \beta^2 u'' + u^3 - u = 0.$$

Using the dynamic approach, the authors prove that it has a homoclinic solution for each  $\beta \in [\sqrt[4]{8} - \epsilon_0, \sqrt[4]{8}]$ , where  $\epsilon_0$  is a small positive constant. This slightly complements Santra and Wei's result [Santra, S. and Wei, J., Homoclinic solutions for fourth order traveling wave equations, *SIAM J. Math. Anal.*, 41, 2009, 2038–2056], which stated that it admits a homoclinic solution for each  $\beta \in (0, \beta_0)$  where  $\beta_0 = 0.9342 \dots$ .

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## 1 Introduction

The well-known Swift-Hohenberg (SH) equation

$$u_t = \mu u - (1 + \partial_{xx})^2 u - u^3 \quad (1.1)$$

was originally derived in the context of Rayleigh-Bénard convection with thermal fluctuations in the limit of infinite Prandtl numbers (see [11]). It captures much of the observed physical behavior and has now become a general tool used to investigate not only the Rayleigh-Bénard convection, but also other pattern-forming systems (see [2]). This equation has been studied intensively by many authors. Mielke and Schneider [7] proved the existence of the global attractor in a weighted Sobolev space on the whole real line. Hsieh et al [4–5] remarked that the elemental instability mechanism is the negative diffusion term  $-u_{xx}$ . Yari [13] discussed the bifurcation and asymptotic behavior of its solutions with the Dirichlet boundary condition and Peletier and Rottschäfer [8] gave some numerical results.

If we consider only the steady solutions of (1.1) for  $\mu > 1$ , it can be written as

$$u'''' + \beta^2 u'' + u^3 - u = 0 \quad (1.2)$$

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after scaling, where the prime ' means taking the derivative with respect to  $x$ . Van den Berg, Peletier and Troy [12] investigated the existence of periodic solutions. Bonheure [1] established the existence of multitransition kinks and pulses obtained as local minima of the associated functional. Smets and Van den Berg [10] used the mountain-pass lemma and Struwe's monotonicity trick to prove that the equation (1.2) has a homoclinic solution for almost all  $\beta \in (0, \sqrt[4]{8})$  while Santra and Wei [9] employed the energy and the Morse index to show that it admits a homoclinic solution for each  $\beta \in (0, \beta_0)$ , where  $\beta_0 = \sqrt{\frac{\sqrt{2}}{k_0}} \approx 0.9342 \dots$ , and  $k_0$  is a solution of  $4k_0^2 - 2k_0 - 3 = 0$ .

For each fixed  $\beta \in [\beta_0, \sqrt[4]{8})$ , whether (1.2) has a homoclinic solution is still an open question. In this paper, we use the dynamic approach and the norm form analysis in particular to obtain that the equation (1.2) has a homoclinic solution for each  $\beta$  less than but close to  $\sqrt[4]{8}$ . This partially complements the result in [9].

## 2 Main Results

Note that (1.2) has three stationary solutions: 0, 1 and  $-1$ . We are here interested in the last two stationary solutions 1 and  $-1$  which correspond to the homoclinic solutions. By symmetry, it is sufficient to consider the solutions homoclinic to 1. Let  $v = u - 1$  which changes (1.2) into

$$v'''' + \beta^2 v'' + v^3 + 3v^2 + 2v = 0. \quad (2.1)$$

The linear equation of (2.1) is

$$v'''' + \beta^2 v'' + 2v = 0, \quad (2.2)$$

whose eigenvalues satisfy

$$\lambda_{\pm}^2 = \frac{-\beta^2 \pm \sqrt{\beta^4 - 8}}{2}. \quad (2.3)$$

Thus, the threshold  $\beta = \sqrt[4]{8}$  corresponds to the upper limit for saddle-focus equilibria. If  $|\beta| < \sqrt[4]{8}$ , then four eigenvalues  $\pm \lambda_{\pm}$  are complex and have nonzero real parts. We have the following theorem.

**Theorem 2.1** *There exists a positive constant  $\epsilon_0$  such that for each  $\beta \in [\sqrt[4]{8} - \epsilon_0, \sqrt[4]{8})$ , the equation (1.2) has a homoclinic solution exponentially approaching 1.*

**Remark 2.1** In the same way, we can prove that (1.2) has a homoclinic solution exponentially approaching  $-1$ .

The proof of Theorem 2.1 is based on the normal form analysis and the results given by Iooss and Pérouème [6]. In the following, we will give the proof.

Firstly we focus on the normal form of (2.1) for  $\beta$  near  $\sqrt[4]{8}$  and assume that for simplicity

$$\beta^2 = \sqrt{8} - a\mu, \quad (2.4)$$

where  $a$  is a positive constant and  $\mu > 0$  is a small parameter. Let  $U = (v, u_1, u_2, u_3)^T$ , where  $u_1 = v_x$ ,  $u_2 = v_{xx}$  and  $u_3 = v_{xxx}$ , which changes (2.1) into

$$U' = AU + \mu A_{\mu}U + N(U), \quad (2.5)$$

where

$$A = \begin{pmatrix} 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \\ -2 & 0 & -\sqrt{8} & 0 \end{pmatrix}, \quad A_{\mu} = \begin{pmatrix} 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & a & 0 \end{pmatrix}, \quad N(U) = \begin{pmatrix} 0 \\ 0 \\ 0 \\ -v^3 - 3v^2 \end{pmatrix}. \quad (2.6)$$

The system (2.5) is reversible with a reverser  $S$  defined by

$$S(v, u_1, u_2, u_3) = (v, -u_1, u_2, -u_3),$$

that is,  $S(v, u_1, u_2, u_3)(-x)$  is also a solution whenever  $(v, u_1, u_2, u_3)(x)$  is. A solution  $(v, u_1, u_2, u_3)$  is reversible if  $S(v, u_1, u_2, u_3)(-x) = (v, u_1, u_2, u_3)(x)$ .

$\mathcal{A}$  has two eigenvalues  $\pm is_1$  which are double, where  $s_1 = \sqrt[4]{2}$ . The eigenvector  $U_1$  and the generalized eigenvector  $U_2$  of  $\mathcal{A}$  corresponding to  $is_1$  are given by

$$U_1 = \left( -\frac{1}{\sqrt[4]{2^3}}, -\frac{i}{\sqrt{2}}, \frac{1}{\sqrt[4]{2}}, i \right)^T, \quad U_2 = \left( -\frac{3}{2}i, \sqrt[4]{2}, \frac{i}{\sqrt{2}}, 0 \right)^T, \quad (2.7)$$

respectively, and they satisfy

$$SU_1 = \overline{U}_1, \quad SU_2 = -\overline{U}_2.$$

The eigenvector  $V_2$  and the generalized eigenvector  $V_1$  of the adjoint operator  $\mathcal{A}^* = \mathcal{A}^T$  corresponding to  $-is_1$  are given by

$$V_1 = \left( \frac{1}{2\sqrt[4]{2}}, 0, \frac{3}{2\sqrt[4]{2^3}}, \frac{i}{2} \right)^T, \quad V_2 = \left( -\frac{i}{2}, \frac{1}{2\sqrt[4]{2}}, -\frac{i}{2\sqrt{2}}, \frac{1}{2\sqrt[4]{2^3}} \right)^T, \quad (2.8)$$

respectively. They satisfy  $(U_j, V_j) = 1$  and  $(U_j, V_k) = 0$  for  $j \neq k$ ,  $j, k = 1, 2$ , where  $(\cdot, \cdot)$  denotes the scalar product in  $\mathbf{C}^4$ .

We can write

$$U = AU_1 + BU_2 + \overline{AU}_1 + \overline{BU}_2 \quad (2.9)$$

such that (2.5) is changed into

$$X' = LX + F(\mu, X), \quad (2.10)$$

where  $X = (A, B, \overline{A}, \overline{B})^T$ ,  $L$  is given by

$$L = \begin{pmatrix} is_1 & 1 & 0 & 0 \\ 0 & is_1 & 0 & 0 \\ 0 & 0 & -is_1 & 1 \\ 0 & 0 & 0 & -is_1 \end{pmatrix}, \quad (2.11)$$

$F(\mu, 0) = 0$ ,  $D_X F(0, 0) = 0$  and  $F(0, X) = O(|X|^2)$ .

Note that the reverser  $S$  is given by

$$S(A, B) = (\overline{A}, -\overline{B}) \quad (2.12)$$

and

$$SF = -FS, \quad SL = -LS. \quad (2.13)$$

From the general theory of normal forms (see [3]), (2.10) can be written as (for the sake of convenience, we still use  $X$ )

$$\begin{aligned} A' &= is_1 A + B + iAP_0(\mu, A\overline{A}, i(A\overline{B} - \overline{A}B)) + R_1(\mu, X), \\ B' &= is_1 B + iBP_0(\mu, A\overline{A}, i(A\overline{B} - \overline{A}B)) + AQ_0(\mu, A\overline{A}, i(A\overline{B} - \overline{A}B)) + R_2(\mu, X), \end{aligned} \quad (2.14)$$

and their complex conjugates, where  $P_0$  and  $Q_0$  are real polynomials of their arguments with degree  $n$  (arbitrary but fixed), and

$$\begin{aligned} P_0(\mu, A\bar{A}, i(A\bar{B} - \bar{A}B)) &= O(|(\mu, A, B)|^{n-1}), \\ Q_0(\mu, A\bar{A}, i(A\bar{B} - \bar{A}B)) &= O(|(\mu, A, B)|^{n-1}), \\ R_1(\mu, X) &= O(|(A, B)| |(\mu, A, B)|^n), \\ R_2(\mu, X) &= O(|(A, B)| |(\mu, A, B)|^n). \end{aligned}$$

Moreover,

$$S \begin{pmatrix} R_1(\mu, X) \\ R_2(\mu, X) \end{pmatrix} = - \begin{pmatrix} R_1(\mu, SX) \\ R_2(\mu, SX) \end{pmatrix}$$

and  $Q_0$  has the form

$$Q_0(\mu, A\bar{A}, i(A\bar{B} - \bar{A}B)) = q_0\mu + q_1A\bar{A} + \cdots,$$

where

$$q_0 = \frac{a}{4} > 0, \quad q_1 = -1.$$

The calculations of  $q_0$  and  $q_1$  are similar to the ones in [3] by (2.7) and (2.8). Note that the normal form of the system (2.14) is exactly the system (4) in [6]. If we replace  $\mu$  by  $-\mu$  and use the results in IV. 3 of [6], we have the following theorem.

**Theorem 2.2** *There exists a positive constant  $\mu_0$  such that for  $0 < \mu \leq \mu_0$  the system (2.14) has a homoclinic solution which is reversible and exponentially tends to 0 as  $x \rightarrow \pm\infty$ .*

From this theorem, we easily obtain Theorem 2.1.

## References

- [1] Bonheure, D., Multitransition kinks and pulses for fourth order equations with a bistable nonlinearity, *Ann. Inst. H. Poincaré Anal. Non Linéaire*, **21**, 2004, 319–340.
- [2] Cross, M. C. and Hohenberg, P. C., Pattern formation outside of equilibrium, *Rev. Mod. Phys.*, **65**, 1993, 851–1112.
- [3] Iooss, G. and Adelmeyer, M., Topics in bifurcation theory and applications, World Scientific, Singapore, 1992.
- [4] Hsieh, D. Y., Elemental mechanisms of hydrodynamic instabilities, *Acta Mech. Sinica*, **10**, 1994, 193–202.
- [5] Hsieh, D. Y., Tang, S. Q. and Wang, X. P., On hydrodynamic instabilities, chaos and phase transition, *Acta Mech. Sinica*, **12**, 1996, 1–14.
- [6] Iooss, G. and Pérouème, M. C., Perturbed homoclinic solutions in reversible 1:1 resonance vector fields, *J. Diff. Eqs.*, **102**, 1993, 62–88.
- [7] Mielke, A. and Schneider, G., Attractors for modulation equations on unbounded domains-existence and comparison, *Nonlinearity*, **8**, 1995, 734–768.
- [8] Peletier, L. A. and Rottschäfer, V., Pattern selection of solutions of the Swift-Hohenberg equation, *Phys. D*, **194**, 2004, 95–126.
- [9] Santra, S. and Wei, J., Homoclinic solutions for fourth order traveling wave equations, *SIAM J. Math. Anal.*, **41**, 2009, 2038–2056.
- [10] Smets, D. and Van den Berg, J. B., Homoclinic solutions for Swift-Hohenberg and suspension bridge type equations, *J. Diff. Eqs.*, **184**, 2002, 78–96.
- [11] Swift, J. and Hohenberg, P. C., Hydrodynamic fluctuations at the convective instability, *Phys. Rev. A*, **15**, 1977, 319–328.
- [12] Van den Berg, J. B., Peletier, L. A. and Troy, W. C., Global branches of multi-bump periodic solutions of the Swift-Hohenberg equation, *Arch. Rational Mech. Anal.*, **158**, 2001, 91–153.
- [13] Yari, M., Attractor bifurcation and final patterns of the  $n$ -dimensional and generalized Swift-Hohenberg equations, *Discrete Contin. Dyn. Syst. Ser. B*, **7**, 2007, 441–456.