# Two-Level Additive Schwarz Methods Using Rough Polyharmonic Splines-Based Coarse Spaces<sup>\*</sup>

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(In Honor of the Scientific Contributions of Professor Luc Tartar)

**Abstract** This paper introduces a domain decomposition preconditioner for elliptic equations with rough coefficients. The coarse space of the domain decomposition method is constructed via the so-called rough polyharmonic splines (RPS for short). As an approximation space of the elliptic problem, RPS is known to recover the quasi-optimal convergence rate and attain the quasi-optimal localization property. The authors lay out the formulation of the RPS based domain decomposition preconditioner, and numerically verify the performance boost of this method through several examples.

Keywords Numerical homogenization, Domain decomposition, Two-level Schwarz additive preconditioner, Rough polyharmonic splines
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## 1 Introduction

Problems with many scales are ubiquitous in nature. Among all the multi-scale problems, the following divergence-form scalar elliptic equation (1.1) with highly heterogeneous coefficients  $\alpha(x) \in L^{\infty}(\Omega)$  is perhaps the most intensively studied one, with a wide range of applications in reservoir modeling, composite materials, etc. Our main objective in this paper is to develop and test a class of overlapping domain decomposition methods using the so-called rough polyharmonic splines (RPS for short) based multiscale coarse grid solvers, which is a multiscale basis with the optimal convergence rate and the optimal localization property (see [23]).

$$\begin{cases} -\nabla \cdot \alpha(x) \nabla u(x) = f(x) & \text{in } \Omega, \\ u = 0 & \text{on } \partial \Omega. \end{cases}$$
(1.1)

The contrast of  $\alpha(x)$  is expressed as the ratio between high and low conductivity values and defined by

$$\frac{\max_{x \in \Omega} \alpha(x)}{\min_{x \in \Omega} \alpha(x)} \tag{1.2}$$

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which brings a small scale into the problem. This type of problems commonly arises in several fields that include subsurface flows and oil reservoir modeling as the representative examples.

The domain decomposition method is a powerful tool to construct efficient parallel solvers for large-scale linear systems arising from the fine scale discretization of partial differential equations (see [26, 28]). To construct preconditioners for the fine-scale system, domain decomposition algorithms include small local problems in the subregions and a coarse problem. Note that without a coarse space component, the algorithms cannot be scalable, i.e., they have a rate of convergence independent of the number of subregions.

The number of iterations required by domain decomposition based methods is determined by the condition number of the preconditioned system. Therefore, the design of preconditioners must address the effects of the following two issues on the condition number, the size of the problem, and the contrast in the media properties. Commonly used domain decomposition methods (see [19, 26, 28]) can make the condition number of the preconditioned system independent of the system size for certain problems including the scalar elliptic equation (1.1).

The design of robust preconditioners with respect to the contrast turns out to be a more challenging problem. When the variation of coefficients is mild inside the coarse-grid blocks, classical domain decomposition preconditioners with a linear coarse space result in a system with the condition number independent of the contrast (see, e.g., [19, 28]). If the variation of coefficients is large inside coarse grid blocks, the classical method fails to be robust. In this case, a judicious choice of the coarse space is the key to constructing the robust domain decomposition preconditioner which has a condition number independent of the coefficient contrast.

In recent works [1, 11, 14], robust preconditioners with respect to the contrast were constructed with the help of coarse space adaptive to small scale features. The connection between multiscale finite elements and robust preconditioners was first explored in [1]. In [14–15], robust preconditioners were constructed for a variety of binary (i.e., two-scale) media model problems using multiscale finite element based coarse space. In a series of works [9–12], the coarse space was further enriched by local spectral basis functions to be suitable for high contrast problems. It is now well understood that the construction of coarse space is closely associated with the development of the multiscale finite-element methods, or in the general sense, numerical homogenization methods.

To this end, we would like to give a very short introduction about numerical homogenization. The field of numerical homogenization concerns the numerical approximation of the solution space of (1.1) with a finite-dimensional space. This problem is motivated by the fact that standard methods, such as the finite-element method with piecewise linear elements (see [4]), can perform arbitrarily badly for PDEs with rough coefficients such as (1.1). Although some numerical homogenization methods are developed from classical homogenization concepts such as periodic homogenization and scale separation (see [5]), as well as localized cell problems (see [24]), one of the main objectives of numerical homogenization is to achieve a numerical approximation of the solution space of (1.1) with arbitrary rough coefficients (i.e., in particular, without the assumptions found in classical homogenization, such as scale separation and ergodicity at fine scales). In this direction, oscillating test functions, G or H-convergence and compensated compactness (see [13, 20, 27]) have always been a great source of inspiration. Professor Luc Tartar has made profound contribution to the development of the above theories. The multiscale finite-element method (MsFEM for short) (see [8, 16–17]) can be seen as a numerical generalization of the idea of oscillating test functions found in H-convergence (see [20]) and was

justified for problems with scale separation. For problems with nonseparable scales, the authors and collaborators proposed the method of harmonic coordinates for scalar elliptic equations in 2D (see [21]). In [6], the transfer property of the flux-norm was introduced to identify the global basis, and then in [22] the computation of the basis was localized to sub-domains of size  $O(\sqrt{H}\ln(\frac{1}{H}))$ . In [22], we also concluded the strong compactness of the solution space, which guarantees the existence of an accurate finite-dimensional approximation space as long as the right-hand side is not too singular. Then the name of the game becomes how to achieve such a finite-dimensional space with the least cost, namely, a space with the best localized basis.

In [23], we introduced an approximation space generated by an interpolation basis over scattered points with resolution H which minimizes the  $L^2$  norm of the source terms; its (pre-)computation involves minimizing  $O(H^d)$  quadratic (cell) problems on (super-)localized sub-domains of size  $O(H \ln(\frac{1}{H}))$ . The resulting interpolation basis functions are biharmonic for  $d \leq 3$ , and polyharmonic for  $d \geq 4$  and can be seen as a generalization of polyharmonic splines to differential operators with arbitrary rough coefficients. Therefore, the basis is called rough polyharmonic splines (RPS for short). The accuracy of the method O(H) in the energy norm is optimal and independent of aspect ratios of the mesh formed by the scattered points. For development in this direction, please also see [3, 18].

In this paper, we use RPS to construct the coarse space of the domain decomposition method, and evaluate the performance of RPS-based preconditioners with some typical domain decomposition preconditioners (linear and MsFEM-based). We apply our method to some benchmark problems and obtain promising numerical results. Theoretical justification of the method is in progress.

#### 2 Mathematical Formulations

In this section, we present the continuous and discrete formulations of the problems that will be considered for preconditioning.

Consider the variational formulation of (1.1) as follows: Find  $u \in \mathcal{H}_0^1(\Omega)$  such that

$$a(u, v) = (f, v) \quad \text{for all } v \in \mathcal{H}_0^1(\Omega),$$

$$(2.1)$$

where

$$a(u,v) := \int_{\Omega} \alpha(x) \nabla u \cdot \nabla v \mathrm{d}x, \quad (f,v) := \int_{\Omega} f v \mathrm{d}x,$$

where  $\Omega \subset \mathbb{R}^2$  is a bounded polygonal domain, and  $f(x) \in L^2(\Omega)$ . The coefficients  $\alpha(x) \geq \alpha_0 > 0$ , and  $\alpha(x) \in L^{\infty}(\Omega)$  are allowed to be highly variable in an unstructured way and have a high contrast in  $\Omega$ . Without loss of generality, we assume that  $\alpha_0 \geq 1$  which can always be achieved by scaling the problem with  $\frac{1}{\min \alpha(x)}$ .

Let the domain  $\overline{\Omega} = \bigcup_{i=1}^{N} \overline{\Omega}_i$  and  $\Omega_i$  be disjoint-shape regular polygonal subdomains of diameters  $H_i$ . Denote the subdomain boundaries by  $\partial \Omega_i$ . For each  $\Omega_i$ , we introduce a quasi-uniform triangulation  $\mathcal{T}_i$  by triangles with the mesh size  $h_i$ . The resulting triangulation  $\mathcal{T}_h$  on  $\Omega$  is assumed to be conforming, i.e., the subdomain meshes match across  $\partial \Omega_i$ .

Without loss of generality, we assume that the coefficient  $\alpha(x)$  restricted to any fine triangle  $\tau \in \mathcal{T}_h$  is a constant, denoted by  $\alpha_{\tau}$ . The analysis will depend on the coefficient on a boundary

layer near subdomain boundaries. For each subdomain  $\Omega_i$ , we define the boundary layer  $\Omega_i^h$  as the union of fine triangles in  $\Omega_i$  that touch the boundaries  $\partial \Omega_i$ . Then we set

$$\overline{\alpha}_i := \max_{x \in \Omega_i^h} \alpha(x). \tag{2.2}$$

Denote by  $V_h(\Omega)$  the standard finite-element space of continuous functions on the domain  $\Omega$ , which are piecewise linear on the fine mesh  $\mathcal{T}_h$  and vanishing on  $\partial\Omega$ . The discrete problem of (2.1) is of the following form: Find  $u_h \in V_h(\Omega)$  such that

$$a(u_h, v_h) = (f, v_h) \quad \text{for all } v_h \in V_h(\Omega).$$
(2.3)

This weak form results in a linear system as

$$Au_h = \mathbf{f},\tag{2.4}$$

where the global stiffness matrix A is large, sparse, and symmetric positive definite, and  $\mathbf{f}$  is the load vector corresponding to the right-hand side of (2.3).

#### 3 Additive Schwarz Methods with RPS-Based Coarse Spaces

The two-level additive Schwarz methods use the solutions of local problems and a coarse problem to construct preconditioners for the fine scale system (2.4). In this section, we introduce an additive Schwarz method with the RPS approximation space as a new coarse space. The construction of RPS generalizes the well-known concepts of polyharmonic splines (see [2, 7, 25]). Unlike standard polyharmonic splines, RPS incorporates information about the coefficients  $\alpha(x)$ . In [23], the properties of RPS, necessary for construction of an optimal localized computational basis, were established and rigorously justified.

We denote by  $\{\Omega'_i\}_{i=1}^N$  the overlapping partition from the original nonoverlapping partition  $\{\Omega_i\}_{i=1}^N$  by extending each subregion  $\Omega_i$  into

$$\Omega_i' = \Omega_i \cup \{ x \in \Omega \,|\, \operatorname{dist}(x, \Omega_i) < \delta_i \}, \quad i = 1, \cdots, N,$$
(3.1)

where the dist is some distance function. Here we consider the case of minimal overlap with  $\delta_i = h_i$ , that is,  $\Omega_i$  is extended to  $\Omega'_i$  by adding one layer of fine triangles in  $\mathcal{T}_i$  that touch the outside of  $\partial \Omega_i$  by edges and/or vertices. The space  $V_h(\Omega)$  can be decomposed as follows:

$$V_h(\Omega) = V_0(\Omega) + V_1(\Omega) + \dots + V_N(\Omega)$$

where

$$V_i(\Omega) = \{ v \in V_h(\Omega) \mid v(x) = 0, \ x \in \Omega \setminus \Omega'_i \} \text{ for all } i = 1, \cdots, N.$$

The local matrix is defined by  $A_i = R_i A R_i^{\mathrm{T}}$ , where  $R_i^{\mathrm{T}} : V_i(\Omega) \to V_h(\Omega)$  is the extension by zero operator.

The coarse space  $V_0(\Omega)$  will be defined in a special way with the coarse basis functions  $\{\Phi_k\}_{k=1}^{N_c}$ , where  $N_c$  is equal to the number of coarse nodes excluding those on the global boundary  $\partial\Omega$ . We associate a specific function  $\Phi_k$  with each coarse node  $x_k$ , i.e., a rough polyharmonic spline centered at  $x_k$ . The definition of RPS is provided below (see also [23]).

#### ASM with RPS Coarse Spaces

Let V be the set of functions  $u \in \mathcal{H}^1_0(\Omega)$  such that  $\operatorname{div}(\alpha \nabla u) \in L^2(\Omega)$ . Let  $\|.\|_V$  be the norm on V defined by

$$||u||_{V} := ||\operatorname{div}(\alpha \nabla u)||_{L^{2}(\Omega)}.$$
(3.2)

For each coarse node  $x_k$  we define

$$V_k := \{ \Phi \in V \mid \Phi(x_k) = 1 \text{ and } \Phi(x_j) = 0 \text{ for } j \in \{1, \cdots, N_c\} \text{ such that } j \neq k \}$$

and consider the following optimization problem over  $V_k$ :

$$\begin{cases} \text{Minimize } \int_{\Omega} |\operatorname{div}(\alpha \nabla \Phi)|^2 \mathrm{d}x, \\ \text{Subject to } \Phi \in V_k. \end{cases}$$

It was shown in [23] that the minimizer of the above problem exists and is unique, which can be finally identified as the coarse basis function  $\Phi_k$ .

Then we can define

$$V_0(\Omega) = \operatorname{span}\{\Phi_k\}_{k=1}^{N_c},\tag{3.3}$$

and the coarse matrix  $A_0 = R_0 A R_0^{\mathrm{T}}$  with  $R_0^{\mathrm{T}} = [\Phi_1, \dots, \Phi_{N_c}]$ , i.e.,  $(R_0)_{kj} = \Phi_k(x_j^h)$  where  $x_j^h$  is the fine nodal point. The corresponding two-level additive Schwarz preconditioner is of the form

$$M_{\rm AS}^{-1} = \sum_{i=0}^{N} R_i^{\rm T} A_i^{-1} R_i.$$
(3.4)

**Theorem 3.1** The condition numbers of the preconditioned stiffness matrices using the proposed two-level additive Schwarz preconditioners satisfy

$$\kappa(M_{\rm AS}^{-1}A) \le C\beta \frac{H}{h},\tag{3.5}$$

where  $\beta = \max_{i} \overline{\alpha}_{i}$ ,  $\frac{H}{h} = \max_{i} \frac{H_{i}}{h_{i}}$ , and C is independent of  $H_{i}$ ,  $h_{i}$  and  $\alpha$ .

**Remark 3.1** For the proof of Theorem 3.1, we follow the general abstract theory for the additive Schwarz methods developed in [26, 28]. Using the same abstract arguments, the proof follows by checking three assumptions. We start with Assumption II (strengthened Cauchy-Schwarz inequalities) and Assumption III (local stability) which are quite easy to verify. In our case, this is straightforward to prove that  $\rho(\mathcal{E}) = 1$  and  $\omega = 1$  in these two assumptions, respectively. Finally, we only need to check the constant in Assumption I (stable decomposition).

The sharpness of the dependence on the contrast and the mesh ratio in (3.5) will be verified numerically in the next section, while the rigorous proof remains the work in progress.

# 4 Numerical Results

Let the domain  $\Omega$  be a unit square  $(0,1)^2$ . We firstly triangulate  $\Omega$  into  $N \times N$  equal coarse squares, then decompose each coarse square into  $M \times M$  equal fine squares and divide each fine square into two sub-triangles along its diagonal with a positive slope. The distribution of coefficients in each example is presented by figures, where  $\alpha(x) = \hat{\alpha}$  in the red shaded region, and  $\alpha(x) = 1$  else where. We use the additive Schwarz method with RPS-based coarse spaces (see Section 3), and iterate with the preconditioned conjugate gradient (PCG for short) method. The iteration in each test stops whenever the  $l_2$  norm of the residual is reduced by a factor of  $10^{-6}$ .

## 4.1 Dependence on the contrast

We implement the two-level additive Schwarz methods with RPS-based coarse spaces for Examples 4.1–4.2 (see Figure 1) to investigate how the condition numbers of preconditioned systems depend on the contrast in the coefficients. Here we take N = 8 and M = 8, i.e.,  $H = \frac{1}{8}$  and  $h = \frac{1}{64}$ .

Figure 1 Left: Example 4.1. Right: Example 4.2.

In Example 4.1, the high-conductivity subregions are completely in the interior of subdomains, while the high-conductivity subregions are only in the boundary layers in Example 4.2. The results in Table 1 indicate that the condition numbers are independent of the contrast in the interior of subdomains, while they may (linearly) depend on the contrast in the boundary layers.

 Table 1 Examples 4.1–4.2. Iteration numbers (condition numbers)

$\widehat{\alpha}$	Example 4.1	Example 4.2
1e0	16(7.392e0)	16(7.392e0)
1e1	15(5.154e0)	25(2.001e1)
1e2	14(4.469e0)	40(1.435e2)
1e3	14(4.438e0)	56(9.536e2)
1e4	15(4.444e0)	58(2.670e3)
1e5	15(4.440e0)	99(1.240e5)
1e6	15(4.443e0)	136(1.234e6)

for different values of contrast, coarse dof 49.

# 4.2 Dependence on the mesh ratio

We test the proposed method with Example 4.1, where the distribution of coefficients is kept as shown in Figure 1 with  $H = \frac{1}{8}$  and  $\hat{\alpha} = 10^6$ . We change M = 8, 16, 32, i.e.,  $h = \frac{1}{64}, \frac{1}{128}, \frac{1}{256}$ . The log-log plot of the condition numbers with respect to the mesh ratio  $\frac{H}{h}$  is reported in Figure 2. The slope of the least square line in Figure 2 is approximated by 1.01, which indicates the linear dependence of the condition number on the mesh ratio.



Figure 2 The log-log plot of the condition number vs.  $\frac{H}{h}$  for Example 4.1.

#### 4.3 Comparisons

We test domain decomposition preconditioners based on the following three coarse spaces: The classical linear coarse space (linear for short), the MsFEM-based coarse space (MsFEM for short), and the global RPS-based coarse space (RPS for short). We consider three different distributions of the coefficients (see Examples 4.3–4.5 in Figure 3, and test different values of the contrast in each example). In all the three examples, we choose N = 8 and M = 10, i.e.,  $H = \frac{1}{8}$  and  $h = \frac{1}{80}$ .



Figure 3 Top Left: Example 4.3. Top Right: Example 4.4. Bottom: Example 4.5.

It follows from the numerical results in Tables 2–4 that for unstructured permeabilities with high contrast, domain decomposition preconditioners constructed by coarse spaces with small scale features (MsFEM and RPS) have a better performance than the classical linear coarse

spaces. In particular, RPS outperforms MsFEM by 3–4 orders of magnitude in the condition number and has a much less iteration counts for the high contrast case. This difference is not so profound for the simpler case, Example 4.3, but for the more complicated permeability fields in Examples 4.4–4.5, the performance gain with RPS is also more significant. We believe it is due to the fact that RPS is a basis with the (quasi-)optimal convergence rate and the (quasi-)optimal localization property as an approximation space of the original elliptic problem (1.1).

Contrast	Linear	MsFEM	RPS
1e0	18(7.906e0)	20(8.843e0)	17(8.686e0)
1e1	22(9.385e0)	21(9.269e0)	18(9.240e0)
1e2	38(2.624e1)	34(2.400e1)	19(9.324e0)
1e3	54(6.594e1)	42(4.021e1)	19(9.323e0)
1e4	61(9.244e1)	47(4.788e1)	19(9.323e0)
1e5	65(9.884e1)	52(5.681e1)	19(9.323e0)
1e6	63(9.957e1)	56(5.832e1)	19(9.301e0)

Table 2 Example 4.3. Iteration numbers (condition numbers) of different methods, coarse dof 49.

Table 3 Example 4.4. Iteration numbers (condition numbers) of different methods, coarse dof 49.

Contrast	Linear	MsFEM	RPS
1 0	10(7,000,0)		17(0,000,0)
1e0	18(7.906e0)	20(8.843e0)	17(8.686e0)
1e1	28(1.632e1)	26(1.449e1)	23(1.292e1)
1e2	59(9.852e1)	48(5.273e1)	30(2.409e1)
1e3	100(8.840e2)	70(3.414e2)	31(2.991e1)
1e4	134(8.725e3)	89(2.919e3)	34(3.079e1)
1e5	164(8.713e4)	127(2.837e4)	36(3.089e1)
1e6	199(8.718e5)	177(2.828e5)	38(3.081e1)

Table 4 Example 4.5. Iteration numbers (condition numbers) of different methods, coarse dof 49.

Contrast	Linear	MsFEM	RPS
1e0	18(7.906e0)	20(8.843e0)	17(8.686e0)
1e1	30(1.744e1)	28(1.648e1)	24(1.587e1)
1e2	62(1.003e2)	48(5.759e1)	35(3.903e1)
1e3	92(8.927e2)	72(3.876e2)	45(8.771e1)
1e4	115(8.808e3)	90(3.764e3)	47(1.051e2)
1e5	141(8.796e4)	116(3.753e4)	44(1.071e2)
1e6	168(8.795e5)	133(3.752e5)	44(1.073e2)

# 5 Conclusion

In this paper, we formulate a domain decomposition method which uses rough polyharmonic splines (RPS for short) to construct its coarse space. As RPS has quasi-optimal properties as an approximation space of the original elliptic equation (1.1), we expect that the performance of the domain decomposition preconditioner will be greatly improved, especially for problems with highly variable coefficients. This is verified numerically through several examples. The theoretical analysis is on progress, which lies in checking the assumption of stable decomposition. A promising future work is to investigate the performance of localized RPS coarse spaces.

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