EXISTENCE AND NON-EXISTENCE OF GLOBAL SMOOTH SOLUTIONS FOR QUASILINEAR HYPERBOLIC SYSTEMS*

and size in our view yet

LIN LONGWEI(林龙威)** ZHENG YONGSHU(郑永树)***

Abstract

Consider initial value probiom $v_t - u_x = 0$, $u_t + p(v)_x = 0$, (E), $v(x, 0) = v_0(x)$, $u(x, 0) = u_0(x)$, (I), where $A \ge 0$, $p(v) = K^2 v^{-\gamma}$, K > 0, $0 < \gamma < 3$. As $0 < \gamma \le 1$, the authors give a sufficient condition for that (E), (I) to have a unique global smooth solution. As $1 \le \gamma < 3$, a necessary condition is given for that.

§ 1. Introduction

The first order quasilinear system

$$v_t - v_x = 0, \ u_t + p(v)_x = -2Au, \quad (A \geqslant 0)$$
 (E)

arises in a variety of ways in several areas of applied mathematics. There has been much investigation to the problem of establishing global existence and non-existence theorems of smooth solutions for initial-boundary value problems.

Nishida^[1] has considered the initial value problem for the damped quasilinear hyperbolic system (E) (p'(v) < 0) with smooth initial data

$$v(x, 0) = v_0(x), \ u(x, 0) = v_0(x),$$
 (I)

which is small with respect to A>0. He has shown that initial value problem for (E), (I) will possess a unique global smooth solution. In other words, under some conditions, there exists $\varepsilon>0$, such that if

$$\sup_{x} |r(x, 0)| + \sup_{x} |s(x, 0)| + \sup_{x} |r_{x}(x, 0)| + \sup_{x} |s_{x}(x, 0)| < s,$$

where r, s are Riemann invariants, then the initial value problem (E), (I) has a unique smooth solution in the large time. Here s>0 depends on A, and $s(A)\to 0$, as $A\to 0$.

Slemrod^[2,8], on the other hand, has shown that if |r(x, 0)| |s(x, 0)| are

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^{**} Department of Mathematics, Zhong Shan University, Guangzhou, Guangdong, China.

^{***} Department of Mathematics, Hua Chiao University, Quanzhou, Fujian, China.

 (λ, Ω)

sufficiently small and $r_{x}(x, 0)$ or $s_{x}(x, 0)$ is negative and sufficiently large in the absolute value at any point $x \in R$, then the initial value problem for (E), (I) has a C^1 -solution only for a finite time.

It is more interesting to investigate the problem with "big" initial value.

In the case $A = 0^{16-111}$ the necessary and sufficient condition of the existence for global smooth solution is

$$r_0'(x) \geqslant 0, \quad s_0'(x) \geqslant 0.$$

The condition is available for "big" initial data.

In section 2 of this paper, we give some conditions for existence or nonexistence of global smooth solution for certain p(v). The conditions are available for $A \geqslant 0$ ica Nec Inch<mark>a</mark>g s wa and "big" initial data.

In detail, suppose that $p(v) = K^2 v^{-\gamma}$, K > 0, $0 < \gamma < 3$, v > 0. Introduce Riemann invariants as $r = u + \Phi(v)$, $s = u - \Phi(v)$, where $\Phi(v) = \int_{1}^{v} \sqrt{-p'(s)} ds$. The main result in section 2 can be stated as follows:

Under conditions (A), (C),

$$\sup_{x} |r_0(x)| + \sup_{x} |s_0(x)| < \min\{-2\Phi(0), 2\Phi(\infty)\}. \tag{A}$$

$$r_0(x), s_0(x) \in C^1(R), |r'_0(x)| + |s'_0(x)| \leq M,$$
 (C)

(1) As $0 < \gamma \le 1$, if condition (B)

$$r'_0(x) \geqslant -\frac{4A}{3-\gamma} v_0(x), s'_0(x) \geqslant -\frac{4A}{3-\gamma} v_0(x)$$
 (B)

holds for all $x \in R$, then the initial value problem (E), (I) has a unique global smooth solution.

(2) As $1 \le \gamma < 3$, if condition (B) fails at any $x \in R$, then the initial value problem (E), (I) has a C¹-solution for only a finite time.

Combining (1), (2), we obtain a necessary and sufficient condition of the existence of global smooth solutions for $\gamma=1$. This is the result of Zheng in [5].

The condition (A) ensures that the C^1 -solution is strictly away from $v=\infty$ (as $1 < \gamma < 3$) or v = 0 (as $0 < \gamma < 1$). In view of the case A = 0 (see [10, 11]) we believe that the right side of the inequality (A) can be replaced by ∞ .

§ 2. The Existence and Nonexistence of Global Smooth Solutions godfaret (8.1) i (6.8)

Consider the initial value problem (E), (I), where $p(v) = K^2 v^{-\gamma}$, K > 0, $0 < \gamma$ <3, v>0. The characteristics are $\lambda = -\sqrt{-p'(v)} = -b^{-2}v^{-\frac{\gamma+1}{2}}, \ \mu = \sqrt{-p'(v)} = b^{-2}v^{-\frac{\gamma+1}{2}},$

$$\lambda = -\sqrt{-p'(v)} = -b^{-2}v^{-\frac{\gamma+1}{2}}, \ \mu = \sqrt{-p'(v)} = b^{-2}v^{-\frac{\gamma+1}{2}},$$

where $b = K^{-\frac{1}{2}} \gamma^{-\frac{1}{4}}$, and the Riemann invariants are taken as

$$r = u + \Phi(v), \quad \mathbf{s} = u - \Phi(v),$$

$$\Phi(v) = \int_{-1}^{v} \mu(s) ds.$$

The Rimann invariants give a one to one smooth mapping from $\Omega = \{(u, v) | u \in R, v \in R_+\}$ onto $\Omega_1 = \{(r, s) | 2\Phi(0) < r - s < 2\Phi(\infty)\}$. The Riemann invariants diagonalize the principal parts of the system (E) as

$$r_t + \lambda r_x = -A (r+s), \quad s_t + \mu s_x = -A (r+s)$$
 (E)₁

and transform the initial data (I) into

$$r(x, 0) = r(u_0(x), v_0(x)) = r_0(x),$$

$$s(x, 0) = s(u_0(x), v_0(x)) = s_0(x),$$

(I)₁

We borrow a priori estimate of Nishida^[1].

Lemma 2.1. Under the conditions (A) (C), the initial value problem (E)₁, (I)₁ has the a priori estimate for C^1 -solution:

$$\sup_{x} |r(x, t)| + \sup_{x} |s(x, t)| \leq \sup_{x} |r_0(x)| + \sup_{x} |s_0(x)| \leq \min\{-2\Phi(0), 2\Phi(\infty)\},\$$

for $t \ge 0$ as long as the C^1 -solution exists. Therefore the solution remains in the region Ω_1 .

We are going to discuss the a priori estimate of $r_x(x, t)$, $s_x(x, t)$. Following Slemrod^[2,3], we have

$$\theta' = -f\theta^2 - A\theta - Ag', \tag{2.1}_1$$

$$\varphi' = -f\varphi^2 - A\varphi - Ag', \qquad (2.1)_2$$

where

$$'=\frac{\partial}{\partial t}+\lambda\frac{\partial}{\partial x}, \qquad =\frac{\partial}{\partial t}+\mu\frac{\partial}{\partial x},$$

$$\theta = \mu^{\frac{1}{2}} \frac{\partial r}{\partial x} = b^{-1} v^{-\frac{r+1}{4}} \frac{\partial r}{\partial x}, \quad \varphi = \mu^{\frac{1}{2}} \frac{\partial s}{\partial x} = b^{-1} v^{-\frac{r+1}{4}} \frac{\partial s}{\partial x}, \tag{2.2}$$

$$f = -\mu_r \mu^{-\frac{1}{2}} = \frac{r+1}{4} b v^{-\frac{3-r}{4}}, \quad g = \int_{0}^{r-s} \frac{1}{2} \mu^{-\frac{1}{2}}(\xi) d\xi. \tag{2.3}$$

Then

$$g' = \frac{1}{2} \mu^{-\frac{1}{2}} (r - s)' = \mu^{-\frac{1}{2}} \overline{\Phi}' = \mu^{\frac{1}{2}} v' = b^{-1} v^{-\frac{r+1}{4}} v', \qquad (2.4)$$

$$\left(\frac{1}{f}\right)' = \left(\frac{4}{r+1}b^{-1}v^{-\frac{3-r}{4}}\right) = C^{-1}b^{-1}v^{-\frac{r+1}{4}}v', \tag{2.5}$$

where

$$C = \frac{r+1}{3-r} \tag{2.6}$$

(2.4), (2.5) implies

$$g' = \left(\frac{C}{f}\right)'. \tag{2.7}$$

Moreover, we set

$$w = \theta + \frac{AO}{f}, \qquad (2.8)$$

so that we can rewrite (2.1)₁ as

$$w' = -f\left(W - \frac{AC}{f}\right)^2 - A\left(w - \frac{AC}{f}\right), \tag{2.9}$$

or

$$w' = -fw^2 + A(2C-1)w + \frac{A^2C}{f}(1-C). \tag{2.9}_2$$

Similarly, put $y = \varphi + \frac{AC}{f}$, (2.1)₂ implies

$$y' = -fy^2 + A(2O-1)y + \frac{A^2O}{f}(1-O).$$

Lemma 2.2. Consider the initial value problem

$$z'(t) = -F(t)z^{2}(t) + A(2C-1)z(t) + \frac{A^{2}C}{F(t)}(1-C), \qquad (2.10)$$

$$z(0) = z_0, (2.11)$$

where A>0, $F(t) \in C^1(R_+)$, $F_* = \inf_{t>0} F(t) > 0$.

(1) If $0 < C \le 1$, $z_0 \ge 0$, then the initial value problem (2.10), (2.11) has the a priori estimate for C^1 -solution z(t):

$$0 \leqslant z(t) \leqslant \max \left\{ \frac{AC}{F}, \ z_0 \right\}. \tag{2.12}$$

(2) If $1 \le C < \infty$, $z_0 < 0$, then the initial value problem (2.10), (2.11) has a C^1 -solution for only a finite time.

Proof (1) The proof is done by contradiction. If the a priori estimate (2.12) fails, then there exists $\tau > 0$, such that

$$z(\tau) < 0, \tag{2.13}$$

or exists $\tau > 0$, such that

$$z(\tau) > \max\left\{\frac{AC}{F}, z_0\right\}.$$
 (2.14)

However, from (2.13), since $z_0 \ge 0$, there exists $0 < \delta \le \tau$, such that

$$z(t) < 0, \ \tau - \delta < t \leqslant \tau, \tag{2.15}$$

$$z(\tau - \delta) = 0. \tag{2.16}$$

(2.16), (2.10) and $F_*>0$ imply

$$z'(\tau - \delta) > 0$$
, as $0 < C < 1$. (2.17)

(2.16), (2.17) contradict (2.15). As C=1,

$$z(t) = 0, \ t \geqslant \tau - \delta, \tag{2.18}$$

is the unique solution of the initial value problem (2.10), (2.16). But (2.18) contradicts (2.15) also.

From (2.14), there exists $0 < \delta \le \tau$, such that

$$z(t) > \max\left\{\frac{AC}{F_x}, z_0\right\}, \quad \tau - \delta < t < \tau,$$
 (2.19)

$$z(\tau-\delta) = \max\left\{\frac{AC}{F_*}, z_0\right\}. \tag{2.20}$$

Rewrite (2.10) as

$$z'(t) = -F\left(z - \frac{AC}{F}\right)^2 - A\left(z - \frac{AC}{F}\right). \tag{2.21}$$

(2.19), (2.21) imply

$$z'(t) < 0, \ \tau - \delta < t \leq \tau. \tag{2.22}$$

(2.20), (2.22) contradict (2.19).

(2) Since $1 \le C < \infty$, $z_0 < 0$, it is easy to prove that z(t) < 0 as long as C^1 —solution z(t) exists, and

$$z'(t) \le -F_*z^2(t) + A(2C-1)z(t)$$
. (2.23)

Consider the following initial value problem

$$\begin{cases} y'(t) = -F_* y^2(t) + A(2C - 1)y(t), \\ y(0) = z_0. \end{cases}$$
 (2.24)

Solving (2.24) by elementary method, we get

$$y(t) = \frac{A(2C-1)z_0}{F_*z_0(1-e^{-A(2C-1)t}) + A(2C-1)e^{-A(2C-1)t}}.$$

Obviously, y(t) goes to $-\infty$ in a finite time for $z_0 < 0$. Comparing (2.23), (2.24), by comparison theorem, we get $z(t) \le y(t) < 0$ as long as z(t) exists. Therefore z(t) exists for only a finite time.

Lemma 2.3. Suppose $6 < \gamma \le 1$. Under the condition (A), (C), any C¹-solution to the initial value problem $(E)_1$, $(I)_1$ has the a priori estimate:

$$|r_{\alpha}(x, t)| \leqslant M, |s_{\alpha}(x, t)| \leqslant M,$$

where M only depends on K, r, A and the bounds of $r_0(x)$, $s_0(x)$, $r_0'(x)$, $s_0'(x)$.

Proof Following the proof of the part (1) for Lemma 2.2, we can prove that C^1 -solution to the equation (2.9) has the a priori estimate

$$0 \leq w(x, t) \leq \max\left\{\frac{AC}{f_*}, w_0\right\}, \tag{2.25}$$

if

$$w(x,0) \geqslant 0, x \in R. \tag{2.26}$$

From conditions (A), (C), $f_* = \inf f(x, t) > 0$, $w_0 = \sup w(x, 0) < \infty$. Since $\mu_r = \mu'(v)v_r$, $\Phi = \frac{1}{2}(r-s)$, $\mu v_r = \frac{1}{2}$, we have $\mu_r = \frac{1}{2}\mu^{-1}\mu' = -\frac{\gamma+1}{4v}$. From (2.7), (2.2), (2.3), (2.6),

$$w = \mu^{\frac{1}{2}} (r_x - AC\mu_r^{-1}) = \mu^{\frac{1}{2}} (r_x + \frac{4Av}{3-\gamma}).$$

Consequently, condition (2.26) is equivalent to the condition (B). r_x is bounded. Its bound depends on K, γ , A and the bounds of $r_0(x)$, $s_0(x)$, $r'_0(x)$, $s'_0(x)$. Similarly, s_x has an analogues a priori estimate.

The a priori estimates in Lemmas 2.1, 2.3, and the local existence theorem for $(E)_1$, $(I)_1$ give the following global existence theorem.

Theorem 2.4. Suppose $0 < \gamma \le 1$. Under conditions (A), (C), if condition (B) holds for all $x \in R$, then the initial value problem (E), (I) has a unique global smooth

solution.

According to the part (2) of Lemma 2.2, the following theorem can be proved by contradition.

Theorem 2.5. Suppose $1 \le \gamma < 3$. Under conditions (A), (C), if condition (B) fails at any point $x \in R$, then the initial value problem (E), (I) has a C^1 -solution for only a finite time.

References

- [1] Nishida, T., Nonlinear hyperbolic equations and related topics in fluid dynamics, Publications Mathematiques D'osay. 78.02, Department de Mathematique, Paris-Sud, 1978.
- [2] Slemrod, M., Instability of steady shearing flows in a nonlinear viscoelastic fluid, Arch. Rational Mech. Anal., 68(1978), 111—225.
- [3] Slemrod, M., Damped conservation laws in continuum mechanics, Nonlinear Analysis and Mechanics, Vol. III, pp. 135—175, Pitman, New York, 1978.
- [4] Bloom, F., On the damped nonlinear evolution equation $W_{tt} = \sigma(W)_{xx} \gamma W_t$, J. Math. Anal. and Appl., 96(1983), 551—583.
- [5] Zheng, Y. S., Global smooth solutions for systems of gas dynamics with the dissipation (to appear).
- [6] Lin, L. W., On the global existence of the continuous solutions of the reducible quasilinear hyperbolic systems, *Acta Scien. Nat. Jilin Univ.*, 4(1963), 83—96.
- [7] Lax, P. D., Development of singularities of solutions on nonlinear hyperbolic partial differential equations, J. Math. Phy., 5(1964), 611—613.
- [8] Johnson, J. L., Global continuous solutions of hyperbolic systems of quasilinear equations, *Bull. Amer. Math. Soc.*, **73**(1967), 639—641.
- [9] Yamaguti, M. and Nishida, T., On the global solution for quasilinear hyperbolic equations, *Funkcial Evac.*, 11(1968), 51—57.
- [10] Lin, L. W., On the vaccum state for the equations of isentropics gas dynamics, J. Math. Anal. & Appli. 121: 2(1982) 406—425.
- [11] Lin, L. W., Vaccum state and equidistribution of the random sequence for Glimm's scheme, J. Math & Appli., 124: 1(1987), 117—126.